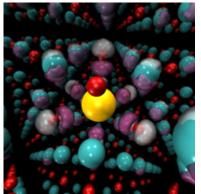
## Ab initio Quantum Monte Carlo

Paul Kent - kentpr@ornl.gov Oak Ridge National Laboratory, Tennessee, USA

https://web.ornl.gov/~kentpr



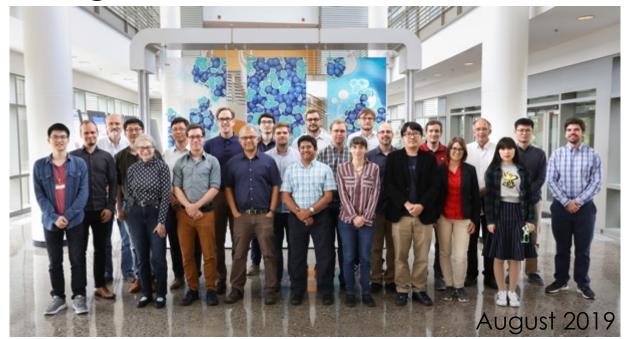






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## Acknowledgments



- Members of Center for Predictive Simulation of Functional Materials, from Livermore, Sandia, Argonne, Oak Ridge Labs, UC Berkeley, Brown, & North Carolina State U.
- While focused on QMC and open source QMCPACK, we also do DFT, DMFT, GW and experimental validation.
- INCITE supercomputer time.





## We are hiring – talk/email/apply ASAP if interested

- Staff scientist in computational many-body theory
  - Interested in all post-DFT methods: GW, DMFT, QMC,...
- Post-doc in Quantum Monte Carlo (Reboredo)
- Distinguished Fellowships
  - "Named postdocs" often leading to a staff position
- Post-doc in first-principles calculations for MXenes
  - Strong experimental collaborations





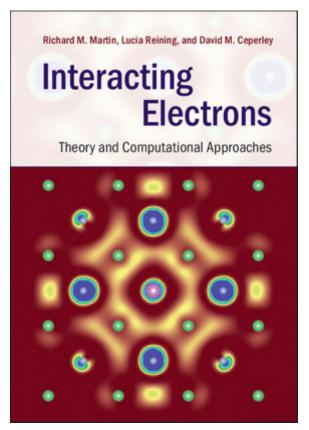


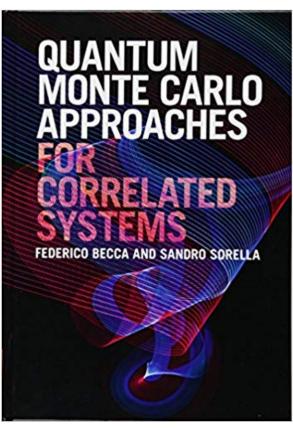


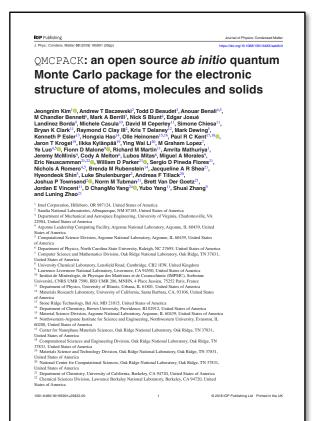
#### Outline

- Motivation and Introduction
  - Why take on an NP hard problem?
  - What is achieved today?
  - Why is the field growing?
- Quantum Monte Carlo Methods
  - There are many; Details matter
  - Real-space VMC, DMC, orbital space AFQMC
- The Future
  - New capabilities
  - New (& old) connections between QMC, GW, DMFT, DFT.

### To find out more







Classic Review: W. Foulkes et al, Rev. Mod. Phys. **73** 33 (2001) [real space only, older] QMCPACK citation: J. Kim et al, JPCM **30** 195901 (2018) [newer references, methods]

## Our Challenge

Solve the many-body Schrodinger equation for general systems and with only readily controllable approximations

#### Our Goals

Develop a practical and convergent method for real materials and chemistry

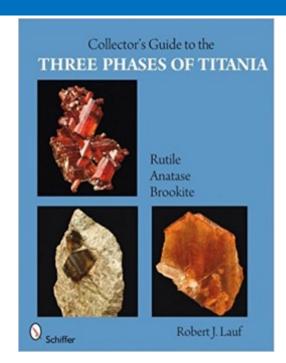
Understand many-body physics, chemistry, materials

Provide useful benchmarks and make connections to other methods, particularly for periodic systems, and eventual upscaling



# Collaboration wanted: Experiment can't provide all the answers we need or are asked to provide

#### No Definitive Result



TiO<sub>2</sub> shows multiple metastable phases. Different theory and experimental camps have argued which is the ground state.

#### **New materials**



MXenes (Prototype: Ti<sub>3</sub>C<sub>2</sub>), the largest family of 2D nanomaterials. Robust predictions needed to guide synthesis.

#### New phenomena



Neutrons show  $\alpha$ -RuCl<sub>3</sub> "is close" to Kitaev quantum spin liquid. Models for QSL depend strongly on Hamiltonian parameters.

## DFT challenges...



**Density functional theory is** straying from the path toward

to chod

Many DFT at We need methods to rationally choose and improve DFT approximations in physics, chemistry, materials

15 MAY 1981

lectron systems

electron densities for atomic species, as functionals. We found that these densitie theoretical advances, until the early 2000 functionals sacrificing physical rigor for

Medvedev et al. S + comm ~ Energies & densities shou Empirical fitting h

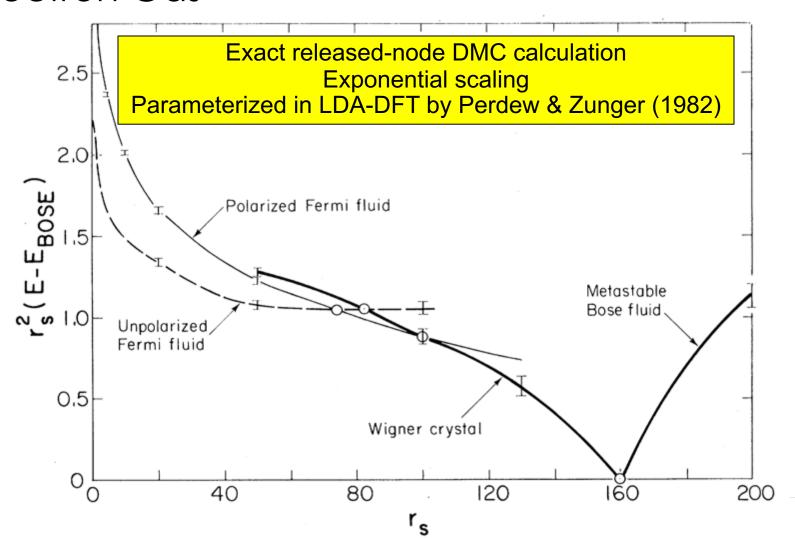
Solar Energy Research Institute, Golden, Colorado 80401 and Department of Physics, University of Colorado, Boulder, Colorado 80302 (Received 31 October 1980)

The exact density functional for the ground-state energy is strictly self-interaction-free (i.e., orbitals demonstrably do not self-interact), but many approximations to it, including the local-spin-density (LSD) approximation for exchange and correlation, are not. We present two related methods for the self-interaction correction (SIC) of any density functional for the energy; correction of the self-consistent one-electron potential follows naturally from the variational principle. Both methods are sanctioned by the Hohenberg-Kohn theorem. Although the first method introduces an orbital-dependent single-particle potential, the second involves a local potential as in the Kohn-Sham scheme. We apply the first method to LSD and show that it properly conserves the number content of the exchangecorrelation hole, while substantially improving the description of its shape. We apply this method to a number of physical problems, where the uncorrected LSD approach produces systematic errors. We find systematic improvements, qualitative as well as quantitative, from this simple correction. Benefits of SIC in atomic calculation include (i) improved values for the total energy and for the separate exchange and correlation pieces of it, (ii) accurate binding energies of negative ions, which are wrongly unstable in LSD, (iii) more accurate electron densities, (iv) orbital eigenvalues that closely approximate physical removal energies, including relaxation, and (v) correct longrange behavior of the potential and density. It appears that SIC can also remedy the LSD underestimate of the band gaps in insulators (as shown by numerical calculations for the rare-gas solids and CuCl), and the LSD overestimate of the cohesive energies of transition metals. The LSD spin splitting in atomic Ni and s-d interconfigurational energies of transition elements are almost unchanged by SIC. We also discuss the admissibility of fractional occupation numbers, and present a parametrization of the electron-gas correlation energy at any density, based on the recent results of Ceperley and Alder.

Perdew & Zunger PRB 23 5048 (1981) The LDA DFT paper is about an error in LDA DFT Most functionals still have self-interaction error (really, delocalization error is key)



## Exact QMC results for the Homogeneous Electron Gas



Ceperley & Alder PRL **45** 566 (1980) See N. Tubman et al. JCP **135** 184109 (2011) for discussion of modern prospects

## Mn Doped Phosphors with DMC

K. Saritas et al. JPCL 10 67 (2019)

Excited states can be run

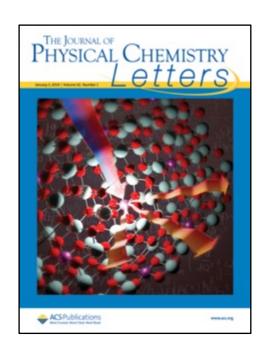
"Standard recipe" fixed-node single determinant DMC gives excellent results for emission energies vs experiment.

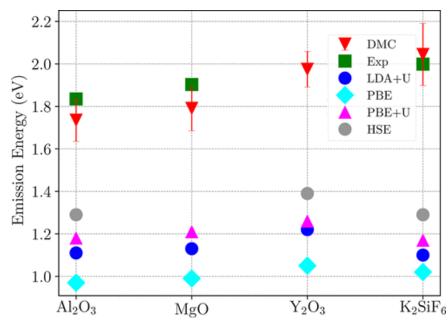
Predicts red Y<sub>2</sub>O<sub>3</sub>:Mn emission

Hybrid DFT is too low.

O(1M) CPU hours

No heroics.





## Scandium oxide polymorphs

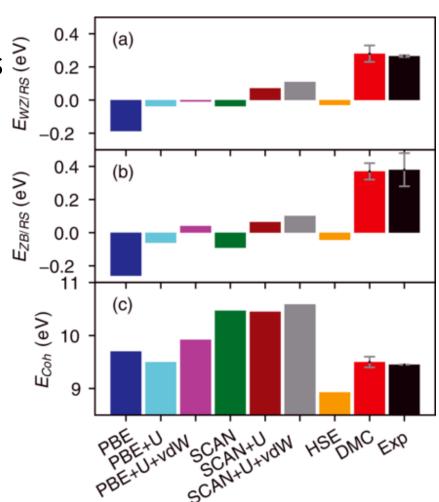
"Standard recipe" fixed-node single determinant DMC gives good thermodynamics, bulk moduli and correct energy ordering for rocksalt, wurtzite, zinc blende.

PBE, HSE, SCAN do not.

SCAN+U corrects order but overbinds.

Followed QMC workshop recipe.

Setup for a later series of perovskite calculations with defects.



K. Saritas et al. PRB 98 155130 (2018)



## Why the Field of QMC is Growing

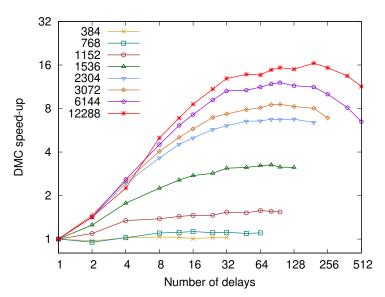


## Massive growth in computing increases accessible materials

- Accelerated computation (GPUs & more) driven by power efficiency
- Memory size & type
- Institutional level resources are very useable.

OAK RIDGE

"Delayed update" algorithm



## New ideas, methods, algorithms

- New methods including AFQMC, FCIQMC...
- Numerics: QMC is 2-10x faster for solids than ~2 years ago
- New related methods:
   Stochastic GWoCG master to edit

#### Outline

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- Quantum Monte Carlo Methods
  - There are many; Details matter
  - Real-space VMC, DMC, orbital space AFQMC
- The Future
  - New capabilities
  - New (& old) connections between QMC, GW, DMFT, DFT.

#### Monte Carlo Methods

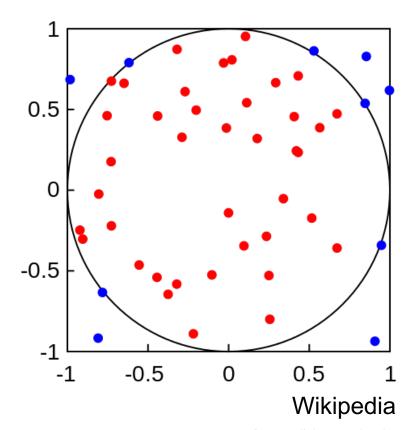
Use statistical methods to tackle the high dimensionality of the Schrodinger equation: Monte Carlo is more efficient than numerical integration in high dimensions.

Trade-off: all measurements have a **statistical error**.

Estimate  $\pi$  via random sampling and ratio of points inside circle to square.

Standard error of mean  $\propto \frac{1}{\sqrt{N}}$ 

100x increase in cost to reduce error 10x!
Sample accurately (importance sampling)
Plan calculations carefully!



## Flavors of Ab initio Quantum Monte Carlo All can treat strong electronic correlations, Van der Waals etc.

- Real space QMC
  - Sample electron positions in real space
  - Variational, diffusion, reptation...
  - Longest established, most results
- Auxiliary Field QMC Motta & Zhang Comp. Mol. Science 8 e1364 (2018)
  - Works in a basis. Strong basis set effects.
  - Easier access to observables than DMC
  - Appears more accurate by default than DMC, but larger cost prefactor. Fewer results.
- Full Configuration Interaction QMC
  - Works in a basis of determinants
  - Potentially exact, expensive. Booth et al. Nature 493 365 (2013)

### Variational Monte Carlo

Use Monte Carlo integration to obtain a variational energy bound from a trial wavefunction. Optimize parameters in trial to reduce variational energy and improve trial wavefunction.

$$E_T = \frac{\int \Psi_T^* \hat{H} \Psi_T}{\int |\Psi_T|^2} \ge E_0$$

$$\Psi_T = \sum_i \!\!\!\! c_i D_i e^J$$

 $\Psi_T = \sum_{\bf c} {\bf c}_{\bf i} D_i e^J \qquad \text{Use best available physics-motivated trial Great flexibility: only need to evaluate value &}$ derivatives in real space

$$E_T = \frac{\int |\Psi_T|^2 \frac{\hat{H}\Psi_T}{\Psi_T}}{\int |\Psi_T|^2}$$

Red-shaded part is a probability density: positive definite & normalized **Use textbook Metropolis Monte Carlo** For 1000 electrons = 3000 dimensional integral

### **Modern QMC trial wavefunctions**

$$\Psi_T = (D_0 + \sum_{ab} c_a^b D_a^b + \sum_{abcd} c_{ab}^{cd} D_{ab}^{cd} + \cdots) \exp(J(r_{ij}, r_{il}, r_{ijl}, \dots))$$



- (1) Start with best affordable density functional or quantum chemistry wavefunction.
- (2) Reoptimize some or all coefficients.

#### Molecules

State of art: O(10<sup>4-5</sup>) determinants + Jastrow and full reoptimization

#### **Solids**

Today mostly single determinant from DFT+U/EXX + Jastrow

Other options: Backflow, Geminals...

Parameterized **Jastrow factor**. **Build in physics**, e.g. wavefunction cusps. *Greatly* improves trial wavefunction. Does not change nodes

$$J(\{\mathbf{r}_i\}, \{\mathbf{r}_I\}) = \sum_{i=1}^{N-1} \sum_{j=i+1}^{N} u(r_{ij}) + \sum_{I=1}^{N_{\text{ions}}} \sum_{i=1}^{N} \chi_I(r_{iI})$$

$$+ \sum_{I=1}^{N_{\text{ions}}} \sum_{i=1}^{N-1} \sum_{j=i+1}^{N} f_I(r_{iI}, r_{jI}, r_{ij}) + \sum_{i=1}^{N-1} \sum_{j=i+1}^{N} p(\mathbf{r}_{ij})$$

$$+ \sum_{i=1}^{N} q(\mathbf{r}_i),$$

e.g. Drummond et al. PRB **70** 235119 (2004)

## Key features of VMC

#### **Advantages**

Explicit form of trial wavefunction

- Can use any trial wavefunction we can imagine
- Easy to compute any observable
- Simple Monte Carlo no timestep error or other discretization introduced

#### **Disadvantages**

Finite size scaling for periodic calculations

Explicit form of trial wavefunction

- Limited to forms of trial wavefunction we can imagine high accuracy difficult for solids, correlated physics.
- Potentially many parameters to optimize reliably. Not yet automated.

### Diffusion Quantum Monte Carlo

**Project out ground state** to minimize trial wavefunction dependence.

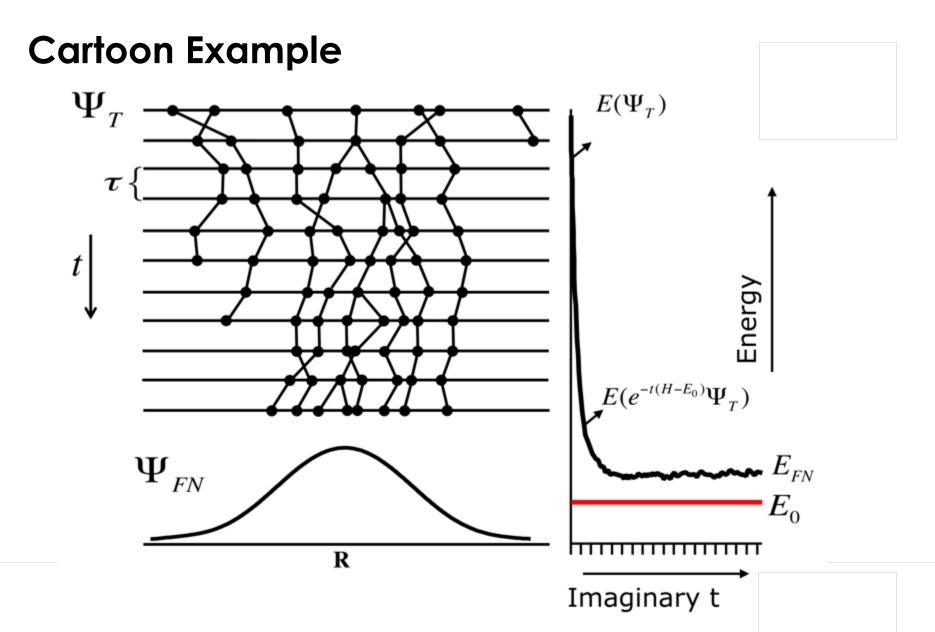
Write the time dependent Schrodinger equation in imaginary time

$$\frac{\partial |\psi\rangle}{\partial \tau} = -\hat{H}|\psi\rangle \qquad \qquad |\psi(\delta\tau)\rangle = \sum_{i=0}^{\infty} c_i e^{-\epsilon_i \delta\tau} |\phi_i\rangle$$

Maps to a branching importance sampled Monte Carlo

Enforce a fermionic solution via the "**fixed-node approximation**". Require nodes of projected wavefunction to be the same as trial. Variational error in energy. **Most significant approximation in DMC**.

Leads to a robust method with good properties: variational, consistently yields high-fraction of correlation energy, formally N<sup>2</sup>-N<sup>4</sup> scaling, readily parallelized...

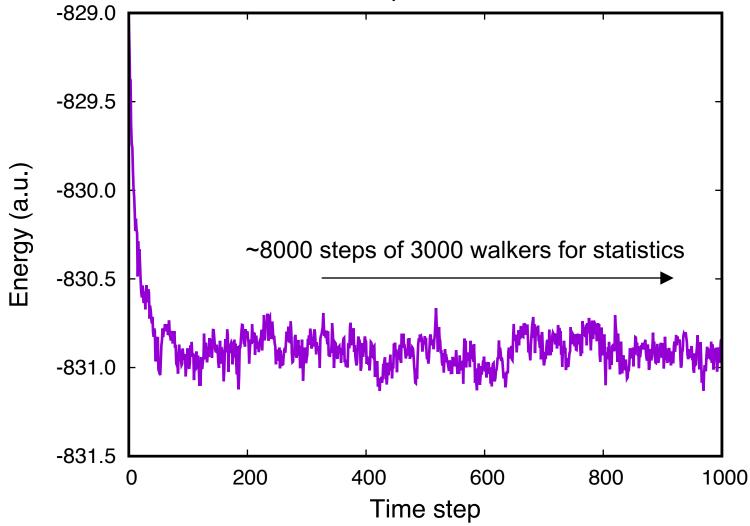


After initial projection to ground state, branching random walk has greatest density where wavefunction probability largest

## Real World DMC: Bulk VO<sub>2</sub>

Production run from Kylanpaa PRM 1 065408 (2017).

200 electrons, 0.01 a.u. timesteps



## Key features of DMC

#### **Advantages**

Gives very accurate and robust results, even with simple nodal surfaces/trial wavefunctions.

Easily take advantages of supercomputers.

Becomes faster, cheaper with improved trial wavefunctions.

#### **Disadvantages**

Finite size scaling & cost in solid-state (Supercells).

Small timestep needed for high Z.

Unlike VMC, no explicit wavefunction obtained. Mixed estimator problem for non-commuting observables.

#### **Errors, Approximations**

Potentially all controllable

"Fixed node"
Variational error in energy is key approximation.

### **Auxiliary Field QMC**

AFQMC is projection method like DMC but works in orbital space. Same or similar as AFQMC used for "model" Hamiltonians.

Far fewer published AFQMC than DMC results. First production open source implementation is in QMCPACK.

#### **AFQMC**

#### Strong basis set error

Flexible treatment of core electrons

Spin-orbit, operators are easy

Direct connection to models

K-point symmetries (NEW)

Well suited to GPUs

Simple workflow

Overall costs (time, memory)

#### **DMC**

Works at basis set limit

Additional approximations needed for non-local potentials

Spin-orbit is not straightforward

Trivial explicit correlation

Memory friendly

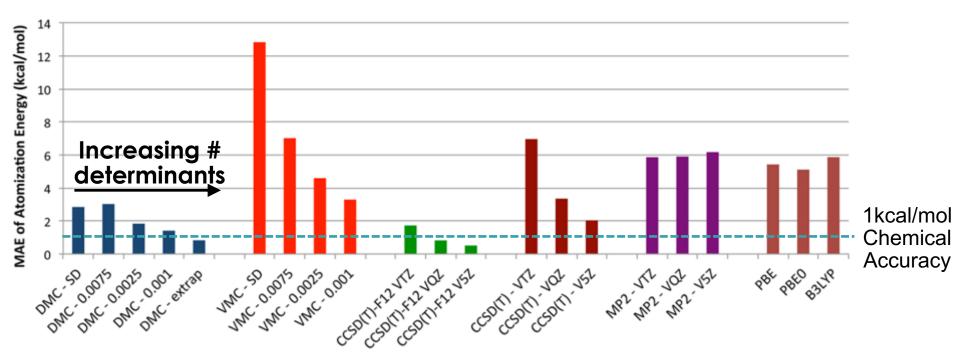
**Complex workflow** 

See <a href="https://github.com/QMCPACK/qmcpack\_workshop\_2019">https://github.com/QMCPACK/qmcpack\_workshop\_2019</a> + YouTube

## Molecular systems

Chemical accuracy <1kcal/mol achieved via the "linear method" for wf optimization of Umrigar et al. PRL (2007).

Trial wavefunctions use large multideterminant expansions DMC better than VMC.



G1 test set. Morales et al. JCTC **8** 2181 (2012) VdW Review: M. Dubecky Chem. Rev. **116** 5188 (2016)

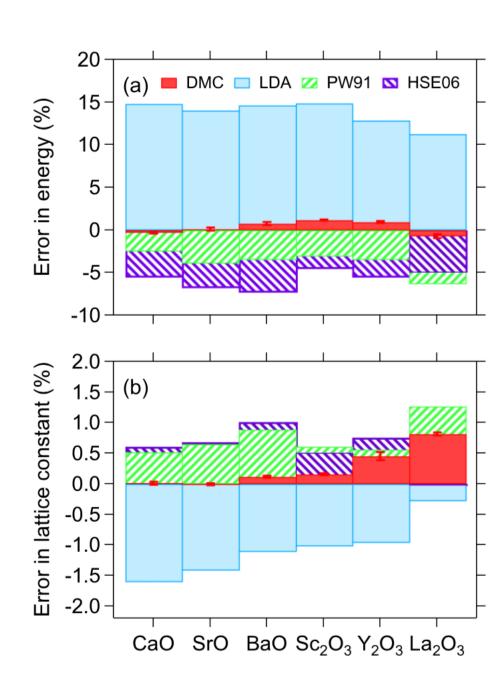
### Solids: Metal oxides

**Single determinant** DMC results are the most accurate in cohesive energy and lattice constant.

Error increases for heavier elements.

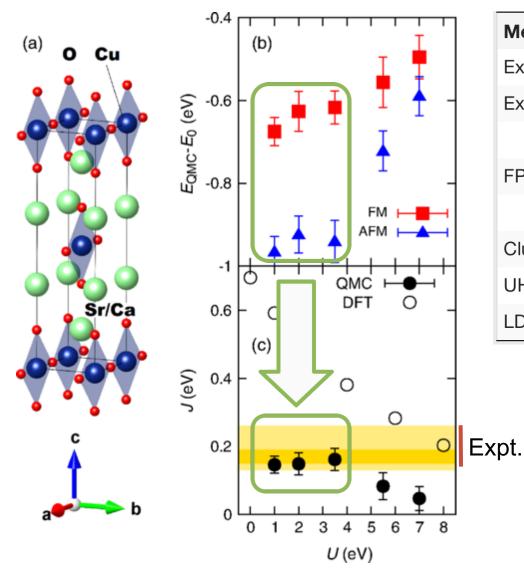
Recall: These are **not** exact calculations. Relative size of nodal and pseudopotential errors is not known.

J. Santana et al. JCP (2016,2017)



## Copper oxides

The best variational DMC results give the best agreement with experiment. Note: DFT+U is simply not predictive.



Method	J (eV)
Experiment (INS)	0.241(11) Ref. [12]
Experiment $[\chi(T)]$	0.146(13) Refs. [7,13]
	0.189(17) Ref. [8]
FP DMC	0.159(14) This work
	0.115(10) This work [14]
Cluster DDCl3	0.231 Ref. [10]
UHF	0.04 Ref. [10]
LDA	0.64 This work

K. Foyevtsova et al. PRX **4** 031003 (2014)

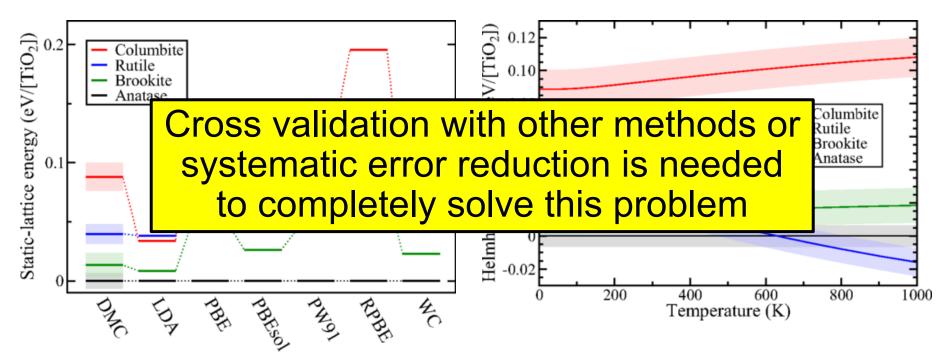
Other studies inc. La<sub>2</sub>CuO<sub>4</sub>: L. Wagner et al. PRB **90** 125129 (2014), PRB **92** 161116 (2015)

## The Titania (TiO<sub>2</sub>) Conundrum

Rutile generally considered to be stable phase at ambient.

Most DFTs at 0K **find anatase most stable**, famously claimed to be an error.

Two QMC results (different codes, choices) also **find anatase most stable** at 0K. Finite T vibrational contributions to stability are important (via DFT).

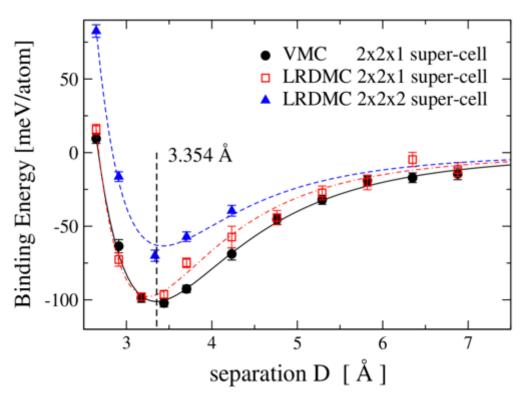


J. Trail et al. PRB **95** 121108 (2017), Y. Luo et al. NJP **18** 113049 (2016)

## Solids: Graphite & Van de Waals

Graphite (A-B stacked graphene sheets) is bound via weak Van de Waals forces. Long a challenge for DFT.

Accurate treatment of Van de Waals critical for the increasing number of 2D materials & heterostructures.



DMC calculations with up to 64 atoms, 256 electrons.

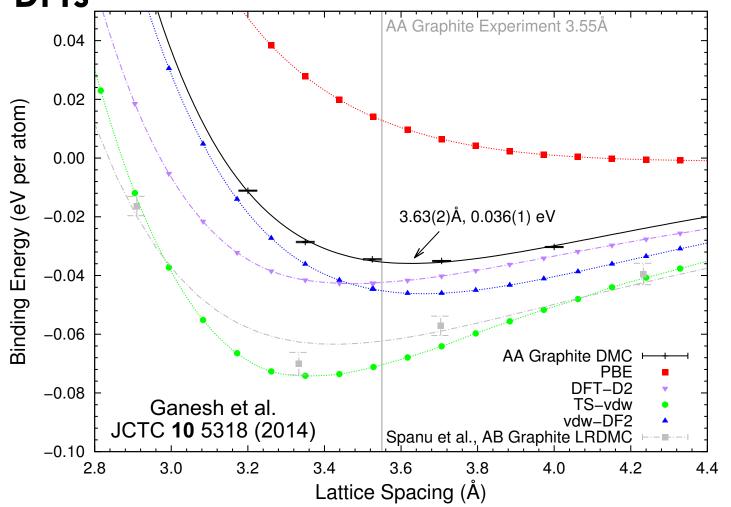
Simple, single determinant trial wavefunctions from DFT.

Predicted binding energy (56 meV/atom) close to experiment (62 meV/atom), lattice parameter within 2%.

Now well predicted by various Van de Waals DFT methods.

L. Spanu, S. Sorella, G. Galli PRL 103 196401 (2009)

QMC for A-A Graphite Helps Identify Preferred VdW DFTs



DMC is within 0.1A of experiment. A-A stacked graphite sensibly higher in energy than A-B stacked ground state.

Self-consistent VdW functionals perform best in this case.

#### The Future

"The future is already here – it's just not very evenly distributed"

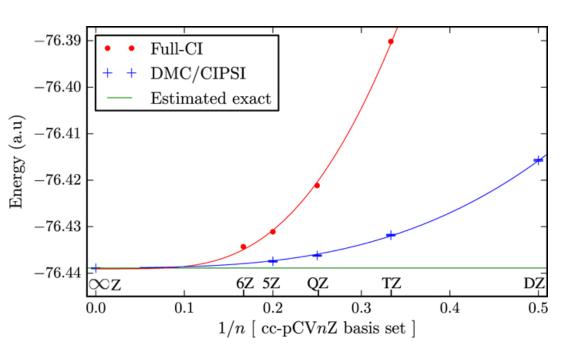
William Gibson

## Towards Systematic convergence

We can use selected-CI methods to build large multideterminant wavefunctions & reoptimize.

Pioneered for QMC by Toulouse group: Use CIPSI to perturbatively grow a wavefunction with a single threshold parameter.

Near push-button workflow developed for molecules.



H<sub>2</sub>O molecule -76.438 94(12) a.u. CIPSI-DMC vs -76.438 9 a.u. Experiment

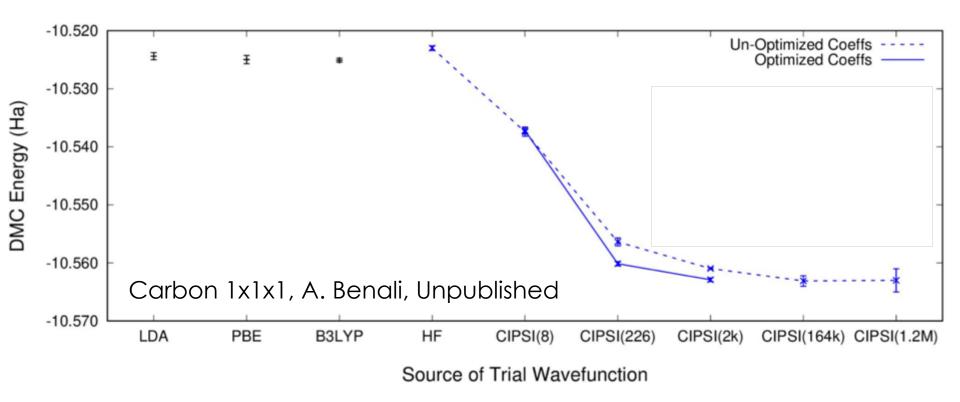
Most accurate theory calculation to-date

Caffarel et al. JCP **144** 151103 (2016)

## Systematic convergence in solids

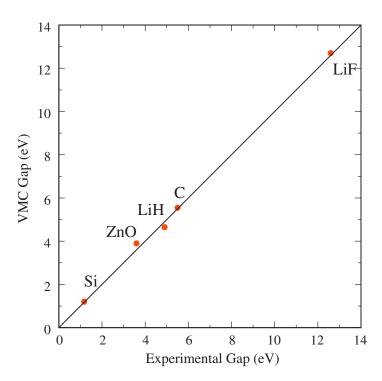
Similar techniques should work in solids – the question is how complex a solid can be treated?

Current practical bottleneck is LCAO interface, Gaussians



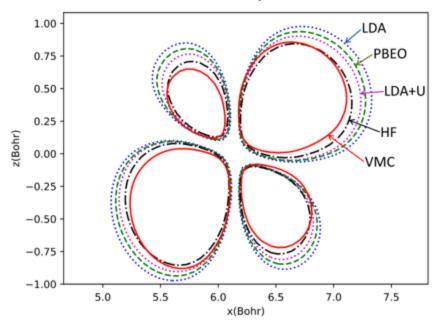
## **Convergent Excited states**

Build multideterminant wavefunctions, selectively target and optimize for ground & excited states, match error (variance) between states.



Simple VMC gives accurate results! No supercomputers used!

#### ZnO hole density isocontours



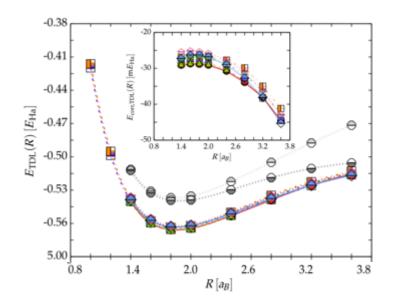
Analyze wavefunction to understand + guide G<sub>0</sub>W<sub>0</sub> accuracy in ZnO



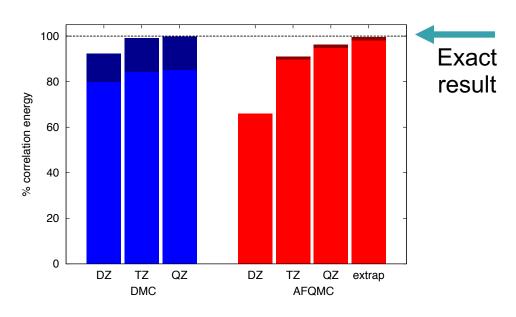
Zhao & Neuscamman PRL **123** 036402 (2019)

#### **Cross-validation**

Use of multiple, distinct methods will give stronger predictions and help drive methods improvement for strongly correlated materials.



H<sub>10</sub> chain
Many methods:
DMC, AFQMC, DMRG,...
Motta PRX 2017



Carbon diamond primitive cell
Multideterminants, large basis sets in DMC
Strong basis set dependence in AFQMC
Unpublished, M. Morales & A. Benali.

## Improving DFT functionals

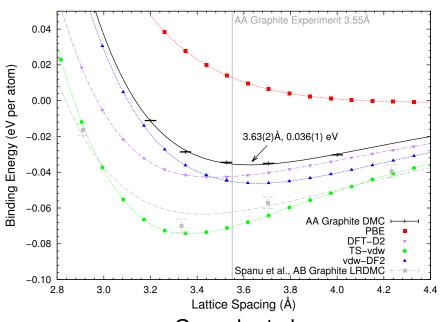
Clearly, this is a hard problem...

Use QMC data from real materials, models to:

- Choose best functional for given application
- Obtain improved understanding by analyzing manybody wavefunctions for solid-state

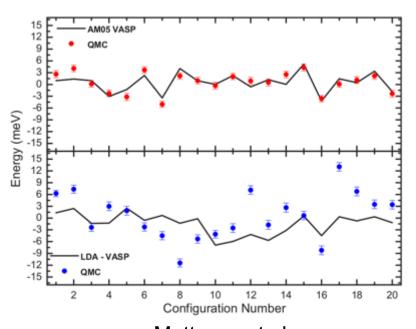
## QMC can already inform DFT choice

## Graphite, Lithium-Graphite Pick best vdw method for more extensive DFT studies



Ganesh et al. JCTC **10** 5318 (2014)

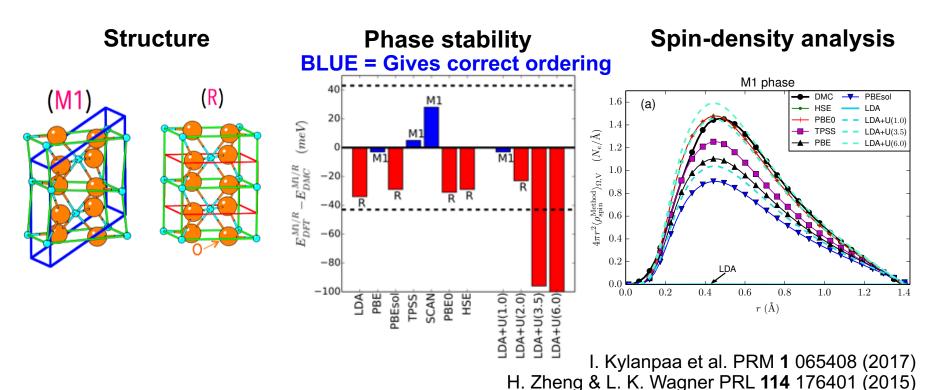
## Kr liquid at high pressure AM05 has best agreement with QMC. Use AM05 for large scale dynamics.



Mattsson et al. PRB **90** 184105 (2014)

## VO<sub>2</sub> metal-insulator transition and phase diagram

Long a challenge for "band theories" to obtain the correct phase ordering and physics. M1 transitions to rutile phase at ~340K, becoming metallic. Analysis of the QMC charge density -- with good statistics -- finds that functionals have difficulty with the vanadium d electrons (presumably self-interaction error.)



With access to energies and densities for many materials, we could empirically fit a functional (dangerous) or help design one (preferred).

## **Analyzing the wavefunction**

A better route to **understanding** successes & failures of DFT in both real materials and model systems could be analyzing the many-body wavefunction in more depth.

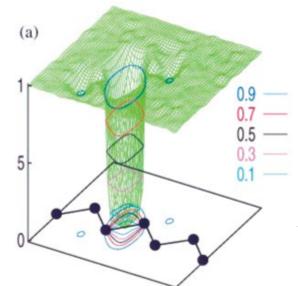
Accessible many-body quantities: density  $n(\mathbf{r})$ , pair correlation function  $g(\mathbf{r},\mathbf{r}')$ , density matrix  $n(\mathbf{r},\mathbf{r}')$ , exchange correlation energy density  $e_{xc}(\mathbf{r})$ ,...

Little explored: bulk Si, some model and atomic systems.

$$\overline{g}_{\alpha\beta}(\mathbf{r},\mathbf{r}') = \frac{N(N-1)}{n_{\alpha}(\mathbf{r})n_{\beta}(\mathbf{r}')} \int_{0}^{1} d\lambda \int d\mathbf{x}_{3} \cdots d\mathbf{x}_{N} \times |\Psi_{\lambda}(\mathbf{r}\alpha,\mathbf{r}'\beta,\mathbf{x}_{3},\ldots,\mathbf{x}_{N})|^{2}, \tag{4}$$

Monte Carlo integral over N-2 particle positions + Coupling constant integration.

R. Q. Hood et al. PRL 78 3350 (1997)



g(**r**,**r**') spin parallel, one electron on bond center, Si (110) plane

## **Bulk silicon analysis**

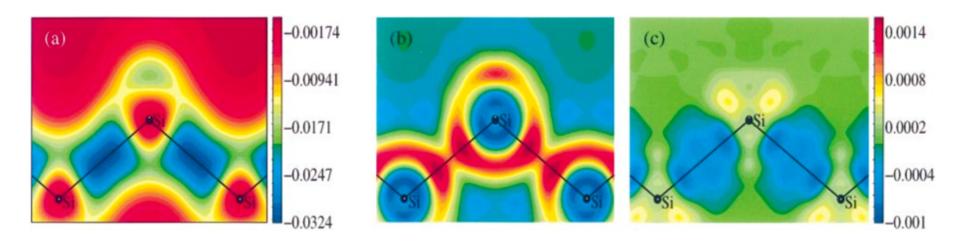


FIG. 4(color). Contour plots along the (110) plane for (a)  $e_{xc}^{VMC}(\mathbf{r})$ , (b)  $e_{xc}^{VMC}(\mathbf{r}) - e_{xc}^{LDA}(\mathbf{r})$ , and (c)  $e_{xc}^{VMC}(\mathbf{r}) - e_{xc}^{ADA}(\mathbf{r})$ . (b) and (c) have the same legend shown to the right of (c). The atoms and bonds are schematically represented for bond chains along the [111] direction.

Some evidence of real space cancelation of errors in LDA ADA (averaged density approximation) performs better than LDA in this case.

Datasets are very rich!

R. Q. Hood et al. PRL **78** 3350 (1997)

R. Q. Hood et al. PRB **57** 8972 (1998)

This analysis has not been repeated in other materials or with modern functionals.

## **Summary: Informing DFT**

- QMC can be used to help select the most accurate existing functional today. Can connect to DMFT the same way. Similar to use of quantum chemistry for molecular systems.
- Analysis of many-body wavefunctions may yield greater understanding.
- Important to heed the lessons of Medvedev et al. Science (2017).

#### Conclusions

Stochastic methods are a promising route to meeting the challenge of the full quantum many-body problem.

Today, QMC can to be applied to important materials where DFT approximations are questioned.

Accurate wavefunctions from QMC in solids can potentially inform DFT, GW, DMFT and other theories.

kentpr@ornl.gov



www.qmcpack.org
Fully open source, github.com/QMCPACK. 50 contributors!



# A New Generation of Effective Core Potentials Pseudopotentiallibrary.org

New many-body construction method H-Kr, L. Mitas and co-workers, JCP 2017, JCP 2018x2, arxiv 2019

Improved accuracy compared to previous potentials including for stretched+compressed bonds.

Open website, various quantum chemistry formats + UPF with KB projectors for plane wave codes

