

Theoretical Spectroscopy and Electronic Excitations

Fabio Caruso, Patrick Rinke and Heiko Appel

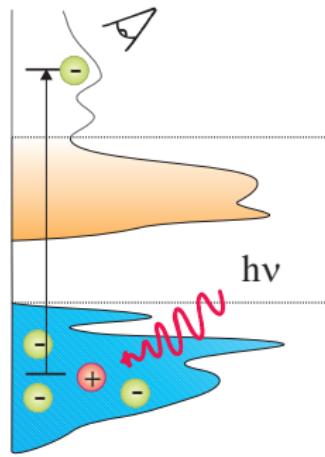
Fritz-Haber-Institut der Max-Planck-Gesellschaft, Berlin

FHI Hands-On Workshop – Tutorial 6

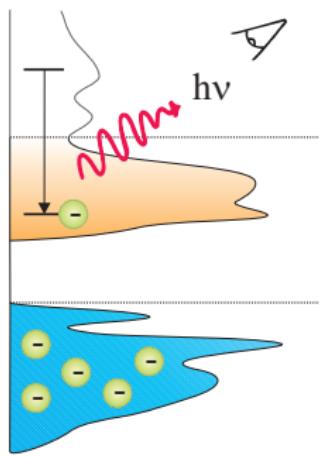


Band structures: photo-electron spectroscopy

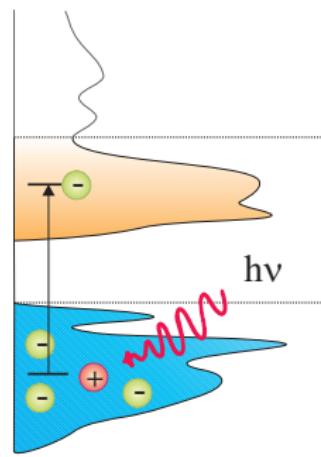
Photoemission



Inverse Photoemission



Absorption

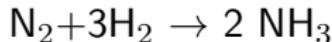


In this tutorial

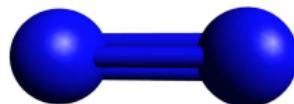
Theoretical spectroscopy of N₂ (and H₂O)

Nitrogen:

- 78% of the earth's atmosphere
- strong triple bond
- ammonia (NH₃) synthesis



- ▶ essential for production of fertilizers
- ▶ Nobel Prize for Fritz Haber in 1918



Haber, F.

Photo-Electron Excitation Energies

• Photoemission

- ▶ electron removal

$$\psi_s(\mathbf{r}) = \langle N - 1, s | \hat{\psi}(\mathbf{r}) | N \rangle$$

- ▶ removal energy

$$\epsilon_s = E(N) - E(N - 1, s)$$

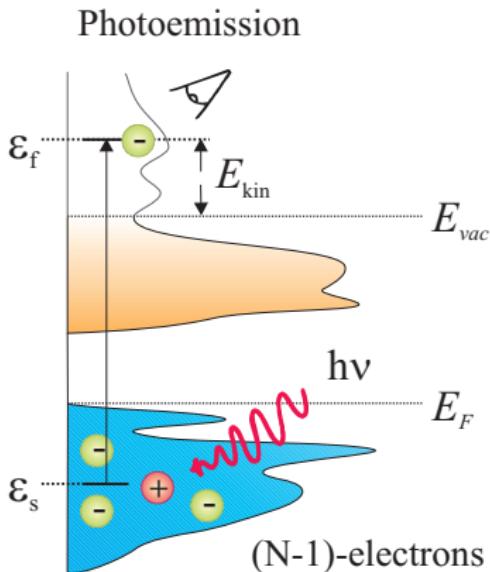
• Inverse Photoemission

- ▶ electron addition

$$\psi_s(\mathbf{r}) = \langle N | \hat{\psi}(\mathbf{r}) | N + 1, s \rangle$$

- ▶ addition energy

$$\epsilon_s = E(N + 1, s) - E(N)$$



$|N\rangle$: N -particle ground state

$|N - 1, s\rangle$: $(N - 1)$ -particle excited state s

$E(N, s)$: total energy in state s

Ionisation Potential, Electron Affinity and (Band) Gaps

Could use total energy method to compute

$$\epsilon_s = E(N \pm 1, s) - E(N)$$

Ionisation potential: *minimal energy to remove an electron*

$$I = E(N - 1) - E(N)$$

Electron affinity: *minimal energy to add an electron*

$$A = E(N) - E(N + 1)$$

(Band) gap: $E_{gap} = I - A$

Ionisation Potential, Electron Affinity and (Band) Gaps

Exercise:

- use HF and DFT-PBE total energies to compute I and A of N_2

This procedure is often called ΔSCF
(difference of two self-consistent field calculations)

Δ SCF versus eigenvalues

Hartree-Fock:

Koopmans' theorem: (Physica 1, 104 (1934))

$$E^*(N \pm 1, s) - E(N) = -\epsilon_s^{\text{HF}}$$

where the same orbitals are used in calculating E^* and E (frozen orbitals)

but:

- orbital relaxations can be important
- correlations absent from HF

Δ SCF versus eigenvalues

Density-functional theory:

- in exact KS-DFT: *ionization potential given by KS eigenvalue of highest occupied state*

$$I_{\text{KS}} = -\epsilon_N(N)$$

- otherwise Janak's theorem: (PRA 18, 7165 (1978)))

$$\frac{\partial E}{\partial n_s} = \epsilon_s$$

rearranging and making mid point approximation:

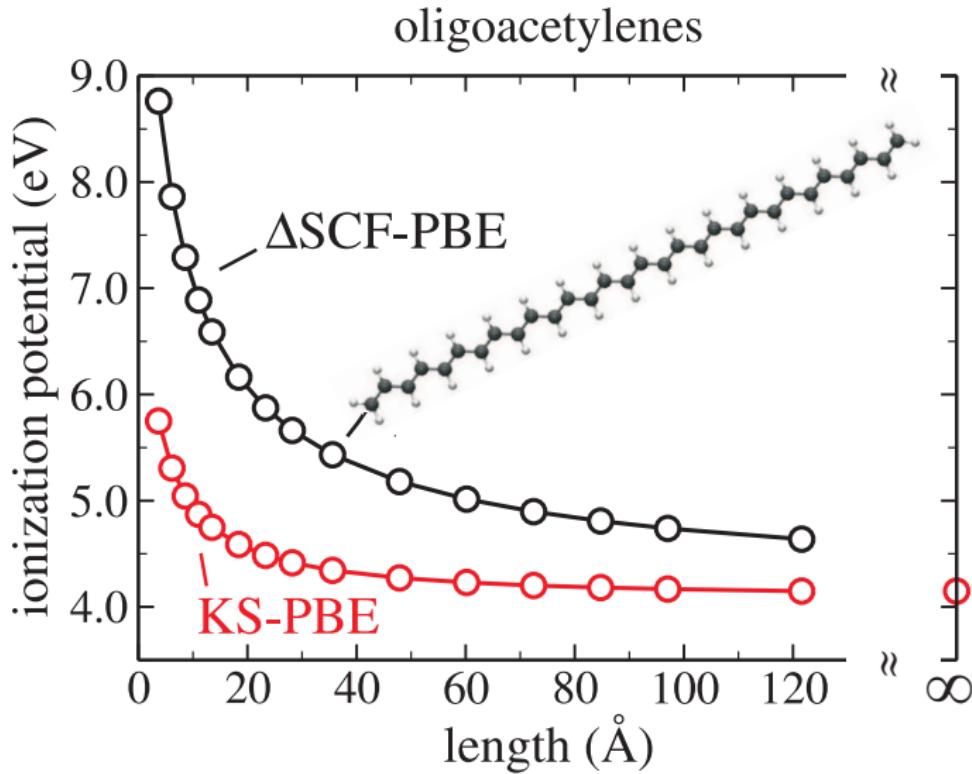
$$E(N+1, s) - E(N) = \int_0^1 dn \epsilon_s(n) \approx \epsilon_s(0.5)$$

Δ SCF versus eigenvalues

Exercise:

- calculate I and A of N_2 using HF and DFT-PBE eigenvalues

Δ SCF versus eigenvalues for finite systems



data courtesy of Max Pinheiro

Single Particle Green's Functions

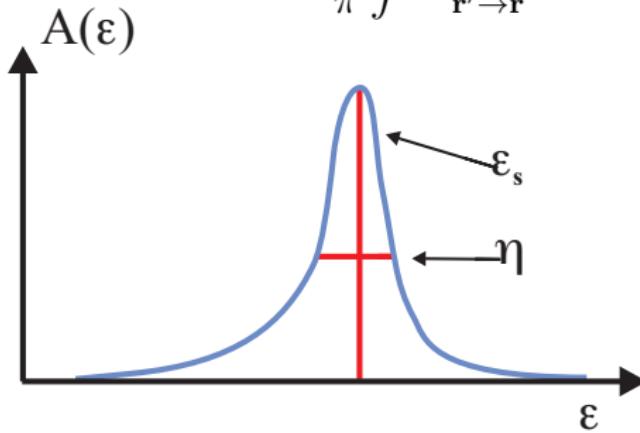
- Lehmann representation of single particle Green's function G :

$$G(\mathbf{r}, \mathbf{r}'; \epsilon) = \lim_{\eta \rightarrow 0^+} \sum_s \frac{\psi_s(\mathbf{r}) \psi_s^*(\mathbf{r}')}{\epsilon - (\epsilon_s + i\eta \operatorname{sgn}(E_f - \epsilon_s))}$$

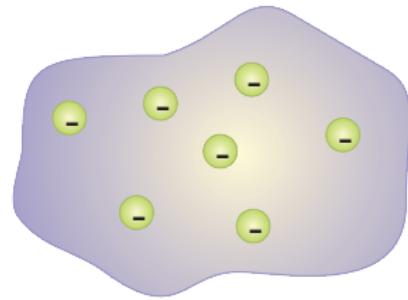
Excitation energies are poles of the Green's function!

- spectroscopically relevant quantity: spectral function A :

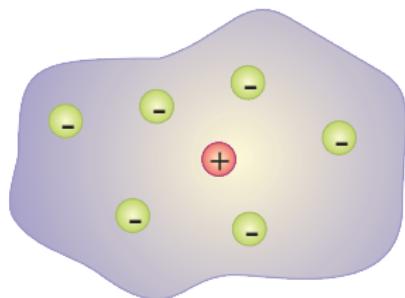
$$A(\epsilon) = -\frac{1}{\pi} \int d\mathbf{r} \lim_{\mathbf{r}' \rightarrow \mathbf{r}} \operatorname{Im} G(\mathbf{r}, \mathbf{r}'; \epsilon)$$



Δ SCF - the infinite time limit

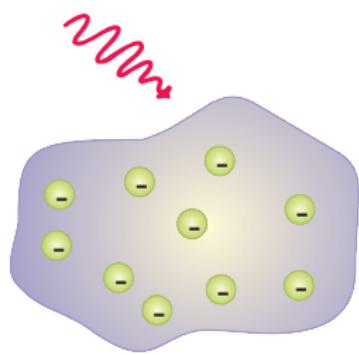


N electrons

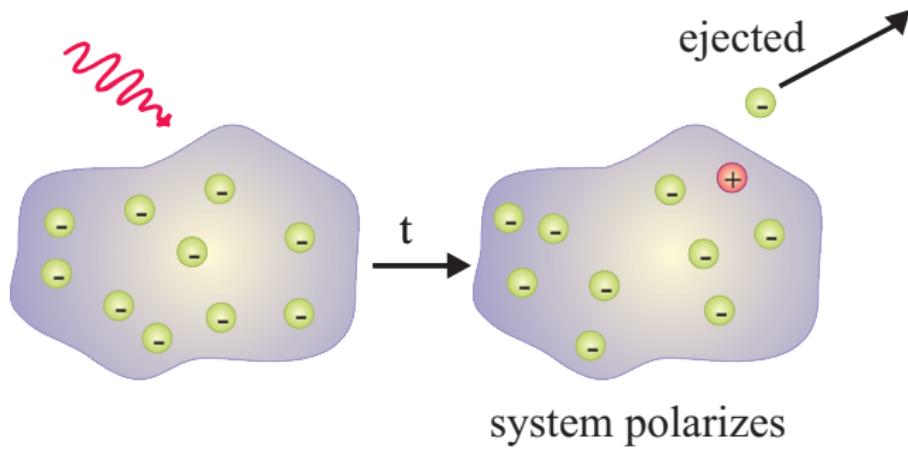


$N-1$ electrons

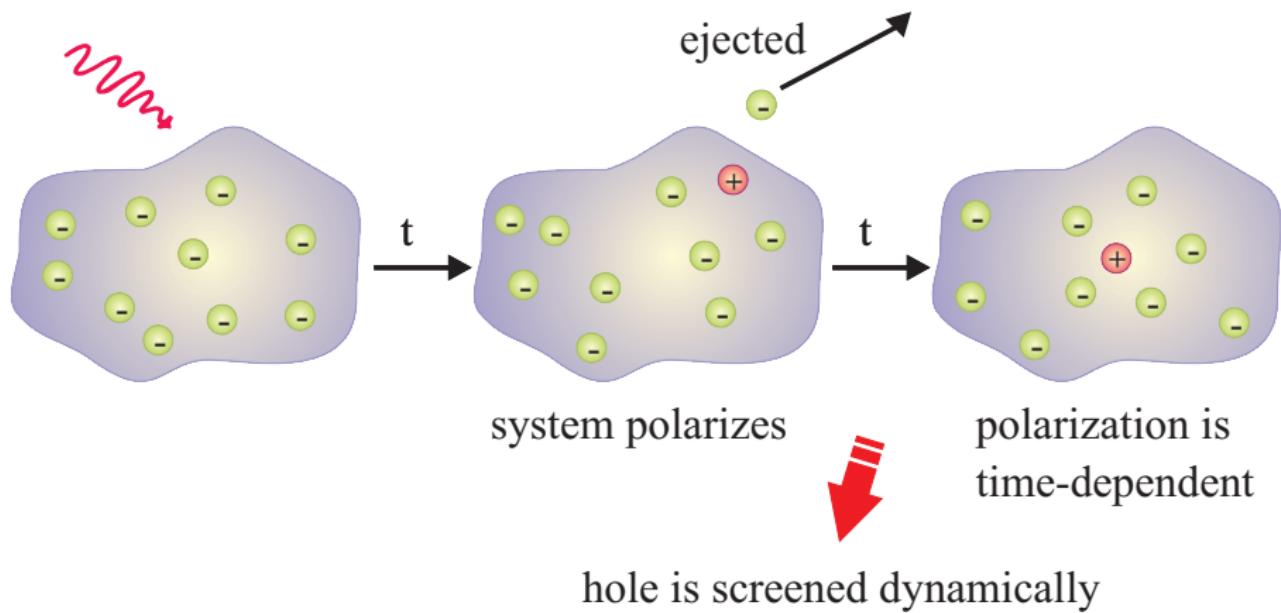
Green's function and screening



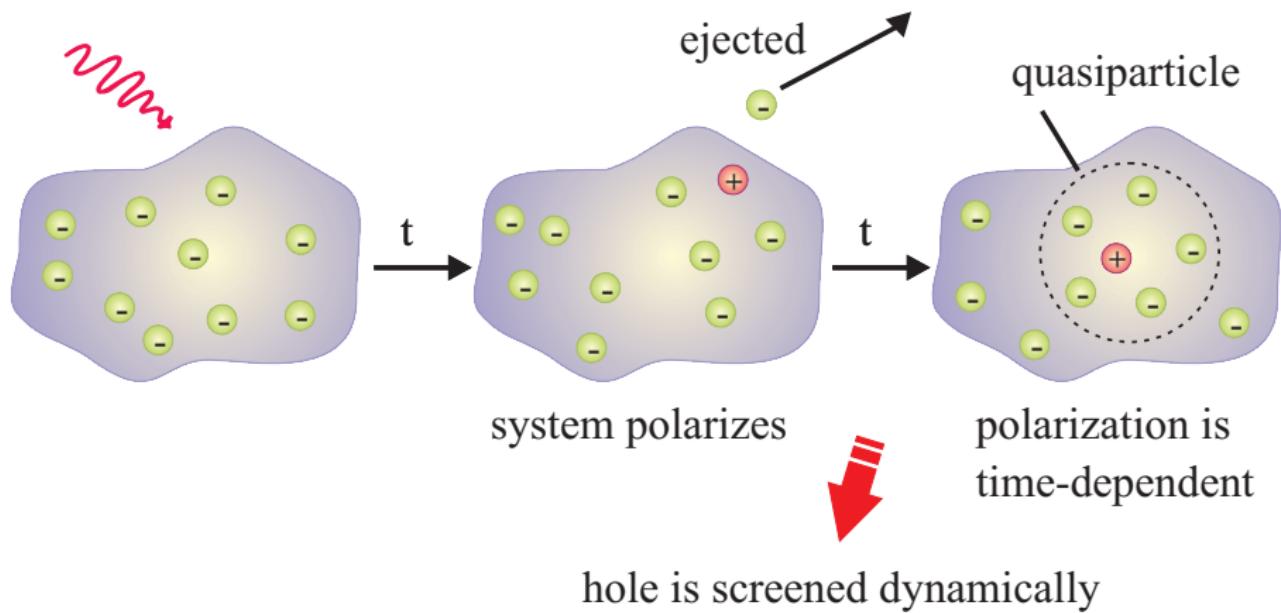
Green's function and screening



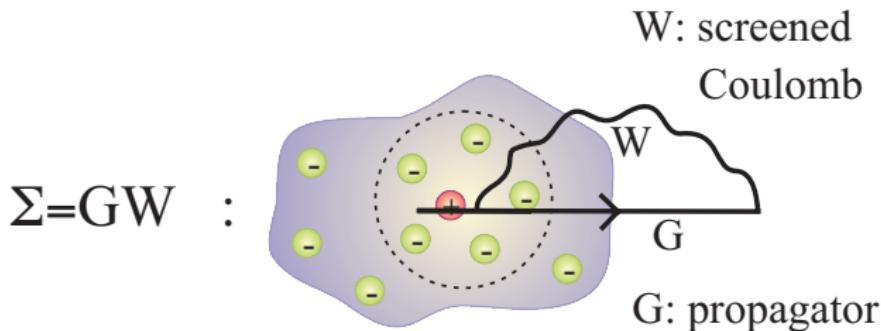
Green's function and screening



Green's function and screening



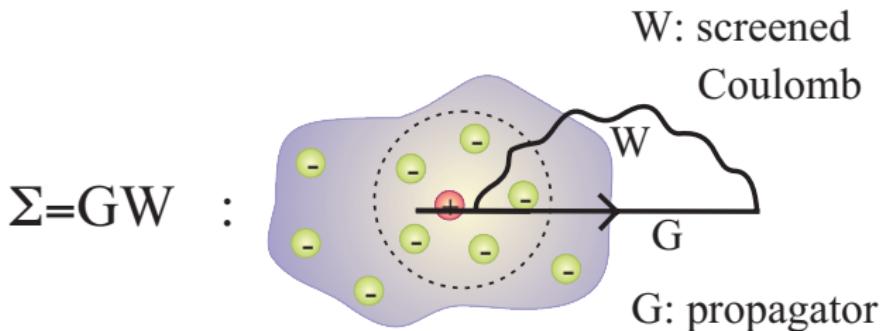
GW Approximation - Screened Electrons



Self-Energy

$$\Sigma^{GW}(\mathbf{r}, \mathbf{r}', \omega) = -\frac{i}{2\pi} \int d\omega e^{i\omega\eta} G(\mathbf{r}, \mathbf{r}', \omega + \omega') W(\mathbf{r}, \mathbf{r}', \omega')$$

GW Approximation - Screened Electrons



Self-Energy: $\Sigma = \Sigma_x + \Sigma_c$

- $\Sigma_x = iGv$:
 - ▶ exact (Hartree-Fock) exchange
- $\Sigma_c = iG(W - v)$:
 - ▶ correlation (screening due to other electrons)

GW – The Issue of Self-Consistency

Hedin's GW equations:

$$G(1, 2) = G_0(1, 2) \quad \text{notation: } 1 = (\mathbf{r}_1, \sigma_1, t_1)$$

$$\Gamma(1, 2, 3) = \delta(1, 2)\delta(1, 3)$$

$$P(1, 2) = -iG(1, 2)G(2, 1^+)$$

$$W(1, 2) = v(1, 2) + \int v(1, 3)P(3, 4)W(4, 2)d(3, 4)$$

$$\Sigma(1, 2) = iG(1, 2)W(2, 1)$$

G_0W_0 quasiparticle energy corrections

- Quasiparticle equation:

$$\hat{h}_0(\mathbf{r})\psi_s(\mathbf{r}) + \int d\mathbf{r}' \Sigma(\mathbf{r}, \mathbf{r}'; \epsilon_s^{qp})\psi_s(\mathbf{r}') = \epsilon_s^{qp}\psi_s(\mathbf{r})$$

- 1 solve for quasiparticle energies ϵ_s^{qp} and wavefunctions $\psi_s(\mathbf{r})$
- 2 perturbation theory: $\psi_s(\mathbf{r}) = \phi_s^{\text{KS}}(\mathbf{r})$

$$\Rightarrow \epsilon_s^{qp} = \epsilon_s^{\text{KS}} + \langle s | \Sigma(\epsilon_s^{qp}) | s \rangle - \langle s | v_{xc} | s \rangle$$

G_0W_0 exercise

Exercise:

- calculate the quasiparticle excitations of N₂ in G_0W_0 starting from
 - ▶ DFT-PBE
 - ▶ HF
 - ▶ DFT-PBE0

GW – The Issue of Self-Consistency

Hedin's *GW* equations:

$$G(1, 2) = G_0(1, 2) \quad \text{notation: } 1 = (\mathbf{r}_1, \sigma_1, t_1)$$

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$$W(1, 2) = v(1, 2) + \int v(1, 3)P(3, 4)W(4, 2)d(3, 4)$$

$$\Sigma(1, 2) = iG(1, 2)W(2, 1)$$

self-consistency

Dyson's equation:

$$G^{-1}(1, 2) = G_0^{-1}(1, 2) - \Sigma(1, 2)$$

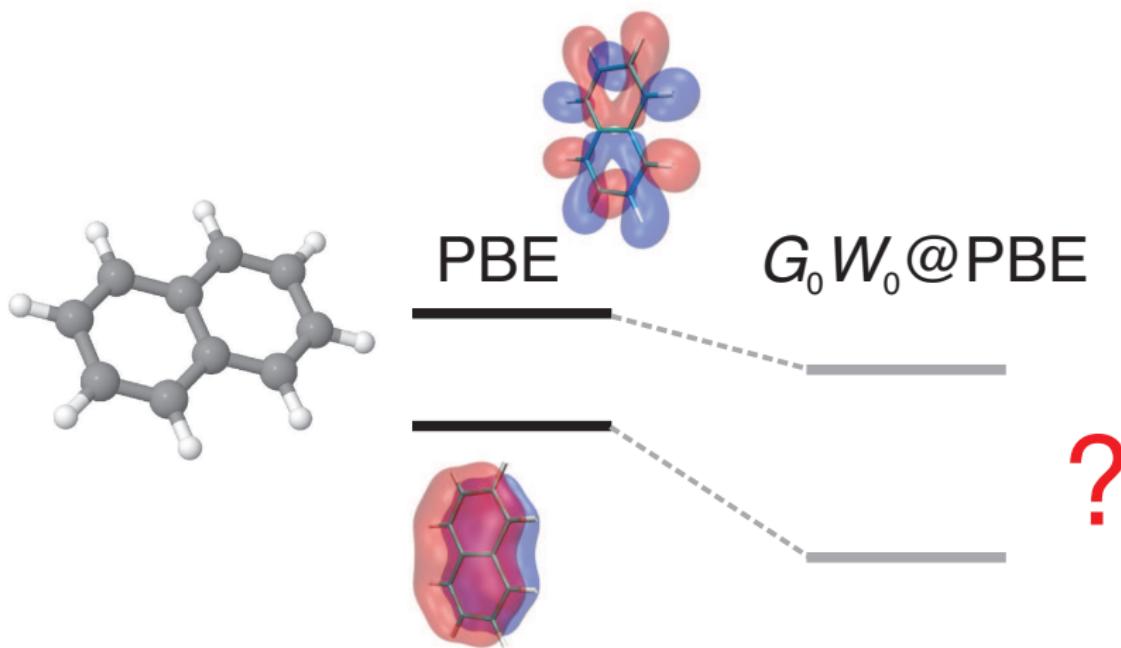
self-consistency

scGW exercise

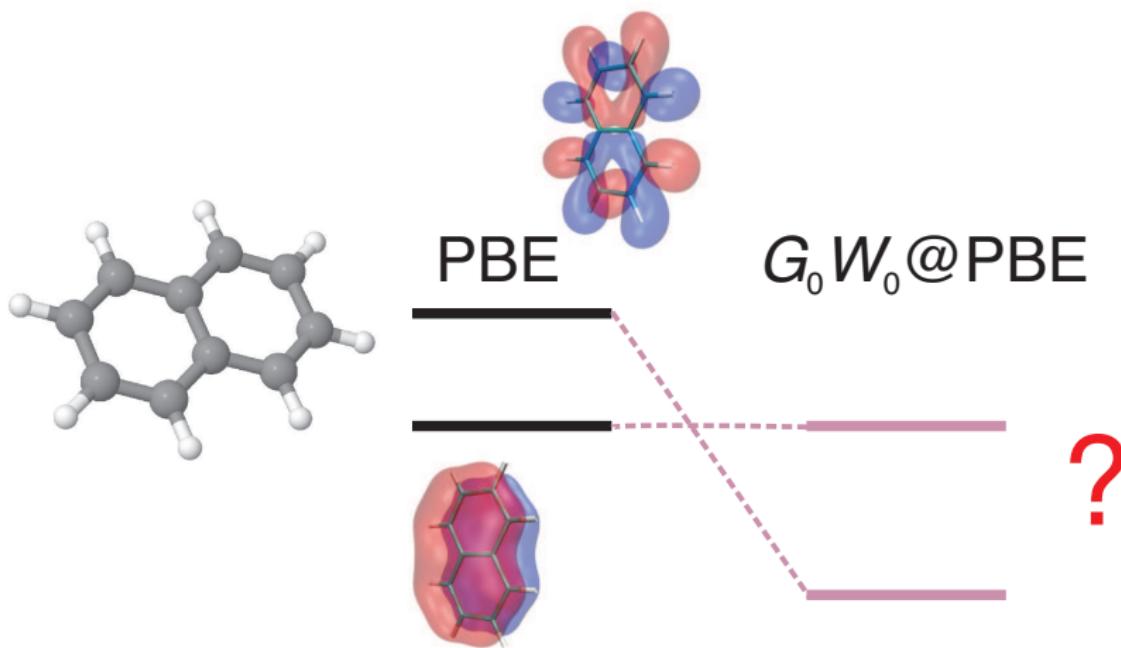
Exercise:

- calculate the quasiparticle spectrum of N_2 in scGW
- extract I and A from the spectrum

Naphthalene and the self-interaction error

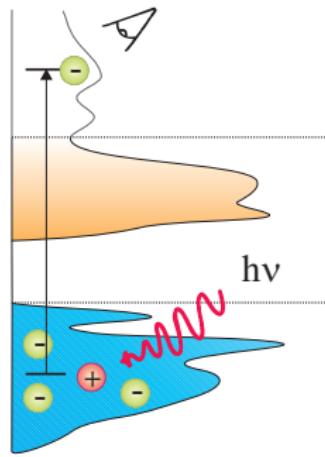


Naphthalene and the self-interaction error

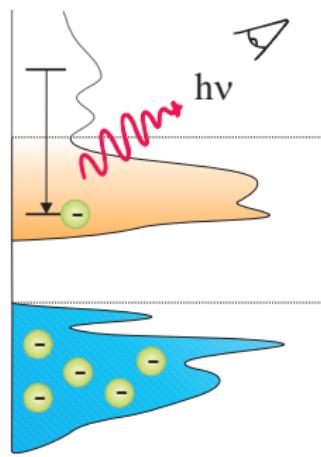


Absorption spectroscopy

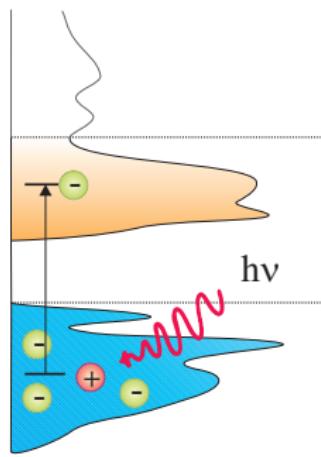
Photoemission



Inverse Photoemission



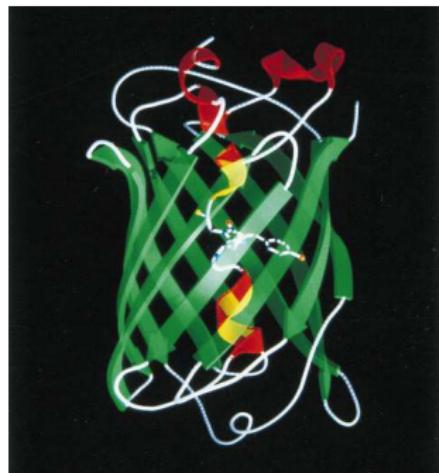
Absorption



Green fluorescent protein

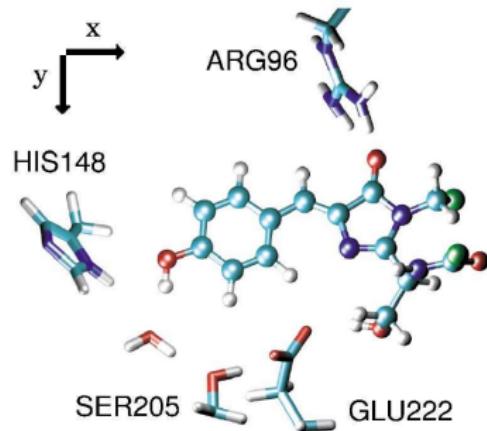
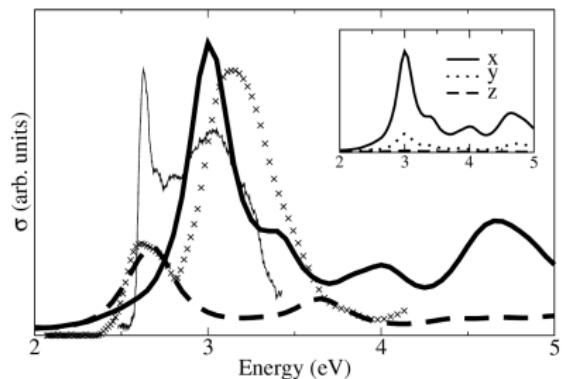


Aequorea victoria - crystal jelly



Ribbon diagram of the Green fluorescent protein (PNAS, 94, 2306 (1997)).

Green fluorescent protein



TDDFT photoabsorption spectrum of neutral (thick solid line) and anionic (thick dashed line) green fluorescent protein chromophores (Marques, et. al., PRL, 90, 258101 (2003).)

Photoabsorption cross section

- Frequency dependent polarizability

$$\alpha_{ij}(\omega) = -\frac{2}{E} \int \rho_1^{(i)}(\mathbf{r}, \omega) r_j d^3r \quad i, j = x, y, z$$

- Photoabsorption cross section

$$\sigma_{ij} = \frac{4\pi\omega}{c} \Im \alpha_{ij}(\omega)$$

- Van der Waals C_6 coefficients from Casimir Polder formula

$$C_6(A, B) = \frac{3}{\pi} \int_0^\infty \alpha_A(i\omega) \alpha_B(i\omega) d\omega$$

Standard linear response formalism

- Exact many-body eigenstates

$$\hat{H}(t = t_0)|m\rangle = E_m|m\rangle$$

- Lehmann representation of linear response function:

$$\chi(\mathbf{r}, \mathbf{r}'; \omega) = \lim_{\eta \rightarrow 0^+} \sum_m \left(\frac{\langle 0 | \hat{\rho}(\mathbf{r}) | m \rangle \langle m | \hat{\rho}(\mathbf{r}) | 0 \rangle}{\omega - (E_m - E_0) + i\eta} - \frac{\langle 0 | \hat{\rho}(\mathbf{r}') | m \rangle \langle m | \hat{\rho}(\mathbf{r}') | 0 \rangle}{\omega + (E_m - E_0) + i\eta} \right)$$

Neutral excitation energies are poles of the linear response function!

- Exact linear density response to perturbation $v_1(\omega)$

$$\rho_1(\omega) = \hat{\chi}(\omega)v_1(\omega)$$

Linear response with TDDFT

- Dyson-type equation for response functions

$$\chi(\mathbf{r}, \mathbf{r}'; \omega) = \chi_S(\mathbf{r}, \mathbf{r}'; \omega) + \int \int \chi_S(\mathbf{r}, \mathbf{r}'; \omega) \left(\frac{1}{|\mathbf{r} - \mathbf{r}'|} + f_{xc}(\mathbf{r}, \mathbf{r}'; \omega) \right) \chi(\mathbf{r}, \mathbf{r}'; \omega) d^3\mathbf{r} d^3\mathbf{r}'$$

- Kohn-Sham linear response function:

$$\chi_S(\mathbf{r}, \mathbf{r}'; \omega) = \lim_{\eta \rightarrow 0^+} \sum_{i,j} \frac{(f_j - f_k) \phi_j(\mathbf{r}) \phi_j^*(\mathbf{r}') \phi_k(\mathbf{r}') \phi_k^*(\mathbf{r})}{\omega - (\epsilon_j - \epsilon_k) + i\eta}$$

Density response from real-time propagation

- Many-body TDSE

$$i\partial_t |\Psi(t)\rangle = (\hat{H}_0 + \alpha H_1(t)) |\Psi(t)\rangle$$

- Density response

$$\begin{aligned}\rho(t) &= \langle \Psi(t) | \hat{\rho} | \Psi(t) \rangle \\ &= \rho_0 + \alpha \rho_1(t) + \alpha^2 \rho_2(t) + \dots\end{aligned}$$

- Advantage: Allows to compute also nonlinear response

Density response from TDDFT real-time propagation

- Time-dependent Kohn-Sham equation

$$i\partial_t |\Phi(t)\rangle = \hat{H}_{KS}(t) |\Phi(t)\rangle$$

- Kohn-Sham density response

$$\begin{aligned}\rho(t) &= \langle \Phi(t) | \hat{\rho} | \Phi(t) \rangle \\ &= \rho_0 + \alpha \rho_1(t) + \alpha^2 \rho_2(t) + \dots\end{aligned}$$

- For the exact functional the time-dependent Kohn-Sham density is equivalent to the exact time-evolving many-body density.
- Need only to consider occupied states!

Real-time evolution - Magnus expansion

- Time-ordered evolution operator

$$\hat{U}(t + \Delta t, t) = \hat{T} \exp \left(-i \int_t^{t+\Delta t} \hat{H}_{\text{KS}}(\tau) d\tau \right)$$

- Magnus expansion

$$\hat{U}(t + \Delta t, t) = \exp \left(\hat{\Omega}_1 + \hat{\Omega}_2 + \hat{\Omega}_3 + \dots \right)$$

- Magnus operators

$$\hat{\Omega}_1 = -i \int_t^{t+\Delta t} \hat{H}_{\text{KS}}(\tau) d\tau$$

$$\hat{\Omega}_2 = \int_t^{t+\Delta t} \int_t^{\tau_1} [\hat{H}_{\text{KS}}(\tau_1), \hat{H}_{\text{KS}}(\tau_2)] d\tau_2 d\tau_1$$

⋮

Real-time evolution - Magnus expansion

- Exponential midpoint rule

$$\begin{aligned}\hat{U}(t + \Delta t, t) &= \exp\left(\hat{\Omega}_1\right) + O(\Delta t^3) \\ \hat{\Omega}_1 &= -i\hat{H}(t + \Delta t/2) + O(\Delta t^3).\end{aligned}$$

- Time-evolved state

$$|\Psi(t + \Delta t)\rangle \approx \exp\left(\hat{H}(t + \Delta t/2)\right) |\Psi(t)\rangle$$

- Task: Compute exponential of operator/matrix

Approximate propagators

- Exact matrix exponential: Diagonalize matrix, exponentiate eigenvalues, transform with eigenvectors
- Taylor/Chebyshev expansion of exponential
- Pade approximation of exponential, e.g. lowest order (Crank-Nicholson)

$$\exp(-i\hat{H}\Delta t) \approx \frac{1 - i\hat{H}\Delta t/2}{1 + i\hat{H}\Delta t/2}$$

$$(1 + i\hat{H}\Delta t/2)|\Psi(t + \Delta t)\rangle = (1 - i\hat{H}\Delta t/2)|\Psi(t)\rangle$$

- Time-reversal symmetry? Symplectic (phase-space preserving) properties?

Properties of approximate real-time propagators

Exercise 6: Getting familiar with real-time propagators

- Crank-Nicholson/Caley, Runge-Kutta 4th order (efficiency, stability)
- Coherent and squeezed states as initial state
- Spectra of autocorrelation functions
- Movies of wavepacket propagation

Absorption spectra from TDDFT

Exercise 7:

- Calculation of photoabsorption cross section for N_2 , H_2O

Part 1: Hands-on charged excitations

What you should do:

N_2	Exp	PBE	HF	Δ -SCF	G_0W_0 @PBE	G_0W_0 @HF	scGW
HOMO							
HOMO-1				***			
HOMO-3				***			

How does DFT eigenvalues compare to experiments for charged excitation?

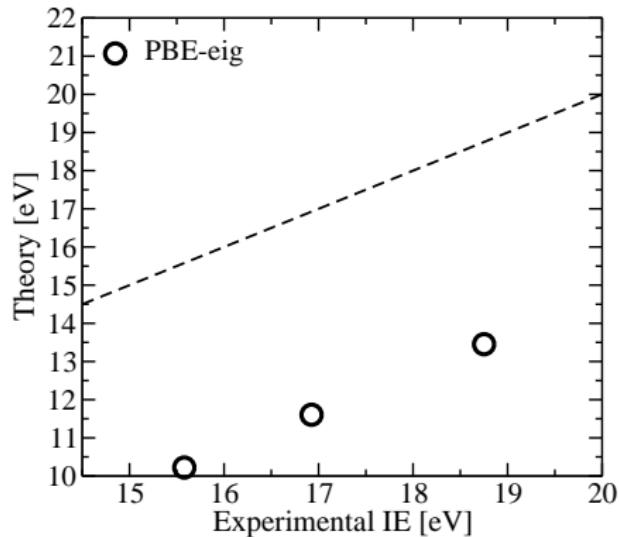
Exercise 1:

Perform a KS-DFT calculation, using the PBE exchange correlation functional for the N_2 molecule (or optionally for H_2O). Compare the first three non-degenerate DFT eigenvalue with the experimental ionization energies of N_2 (see documentation),

- What is the average difference between the Kohn-Sham eigenvalues and experimental values?

How does DFT eigenvalues compare to experiments for charged excitation?

Exercise 1: RESULTS



Hartree-Fock ionization energies: the Koopman theorem

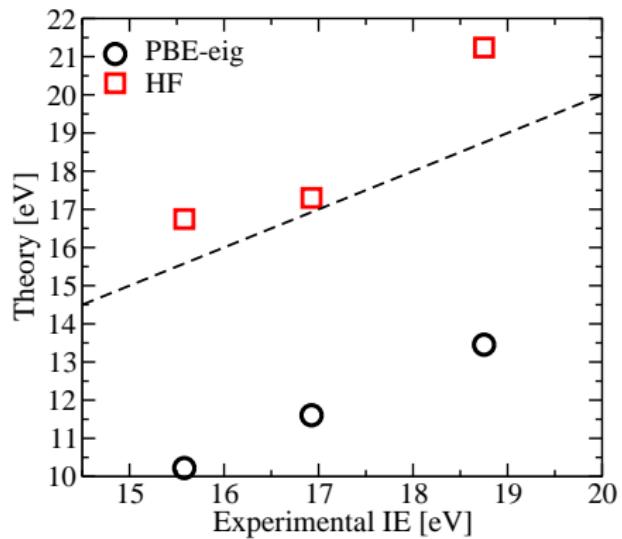
Exercise 2:

Evaluate the first 3 ionization energies of N_2 from an Hartree-Fock calculation. In Hartree-Fock, the Koopman theorem justifies the interpretation of the eigenvalues as addition and removal energies.

- What is the effect of neglecting the electron-electron correlation on ionization energies?

Hartree-Fock ionization energies: the Koopman theorem

Exercise 2: RESULTS



Electron removal energies from the Δ -SCF approach

Exercise 3:

Evaluate the first ionization energy I and the electron affinity A of N_2 from KS-DFT and Hartree-Fock total energy differences using the formula:

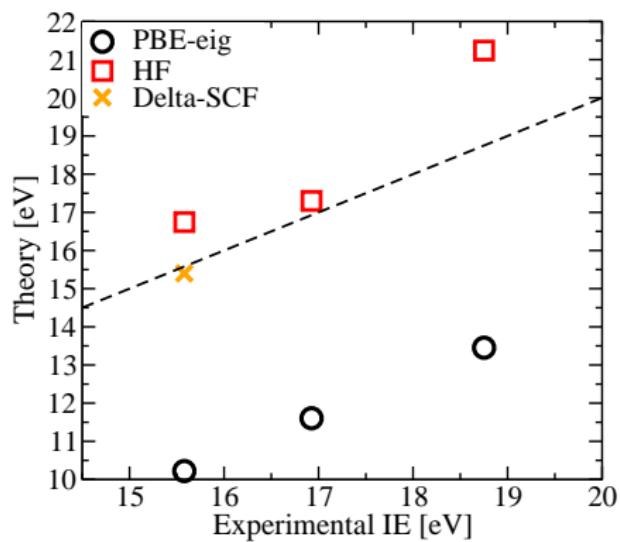
$$I = E(N - 1) - E(N) \quad (1)$$

$$A = E(N) - E(N + 1) \quad (2)$$

- Is the agreement with the first experimental ionization energy improved compared to those computed in Exercise 1?
- Is the Koopman theorem satisfied?

The GW approximation: self-consistent GW

Exercise 3: RESULTS



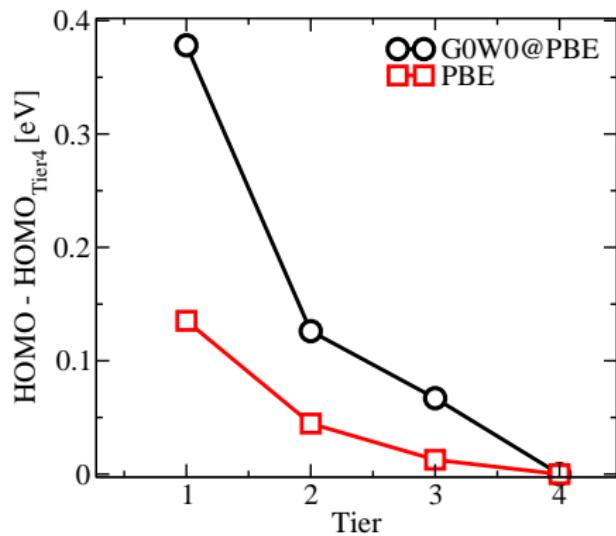
The GW approximation: one-shot G_0W_0

Exercise 4:

- Perform a $G_0W_0@\text{PBE}$ calculation for N_2 . How do the quasi-particle energies compare with experiments?
- Basis set convergence: plot the convergence of the quasi-particle HOMO level of N_2 for Tier 1, 2, 3 and 4.
- **Optional:** Dependence on the starting point. Perform a $G_0W_0@\text{HF}$ and a $G_0W_0@\text{PBE0}$ calculation of the ionization energies and compare with $G_0W_0@\text{PBE}$ the obtained values.

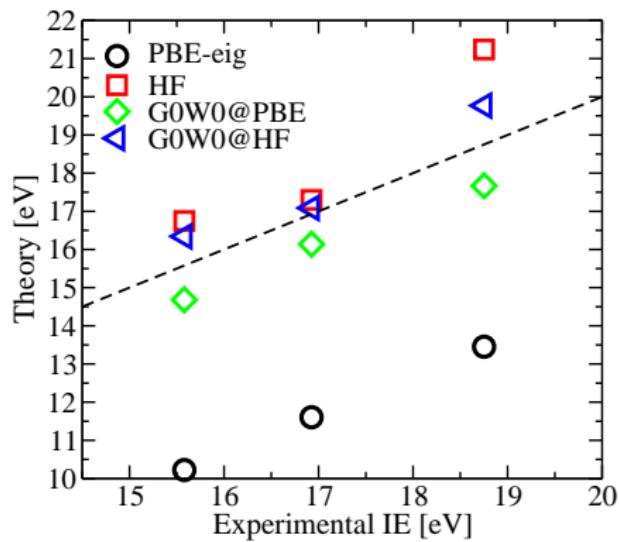
The GW approximation: one-shot G_0W_0

Exercise 4: RESULTS - Basis set convergence



The GW approximation: one-shot G_0W_0

Exercise 4: RESULTS - Ionization energies



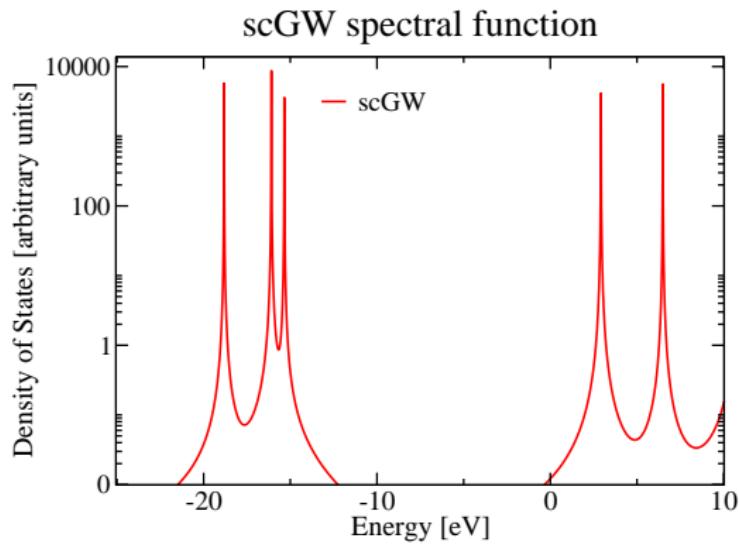
The GW approximation: self-consistent GW

Exercise 5:

- Perform a self-consistent GW calculation for N_2 (**Tier 1** basis set) and plot the GW spectral function.
- Extract the ionization energies from the spectral function (gnuplot, xmGrace, ...).
- Compare the density obtained from the self-consistent Green's function and PBE .
- **Optional:** Try sc- GW with a Hartree-Fock starting point.

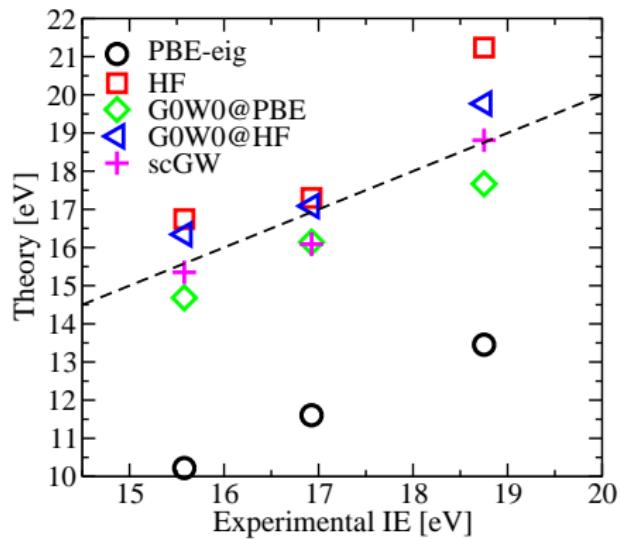
The GW approximation: self-consistent GW

Exercise 5: RESULTS

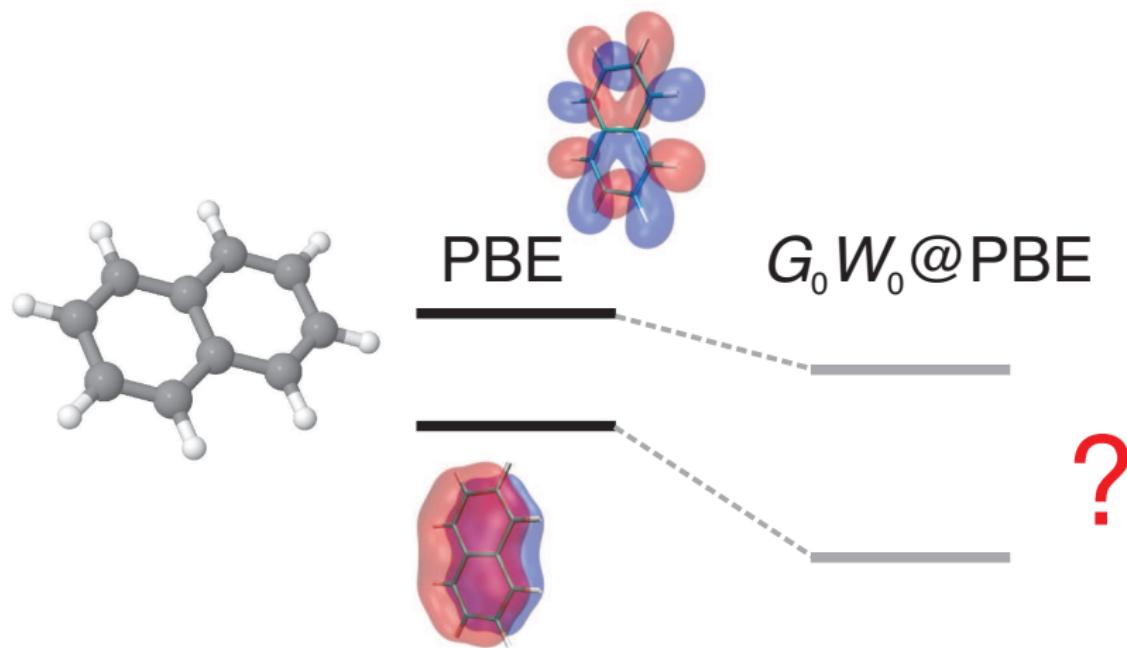


The GW approximation: self-consistent GW

Exercise 5: RESULTS



Naphthalene and the self-interaction error



Self-interaction error in Naphthalene

Bonus Exercise:

- Perform a G_0W_0 @PBE calculation of Naphthalene.
- Determine which orbitals are energetically swapped in G_0W_0 @PBE compared to PBE.
- Visualize the orbital 27 and 28. Is the shape of the orbitals consistent with the self-interaction error picture?

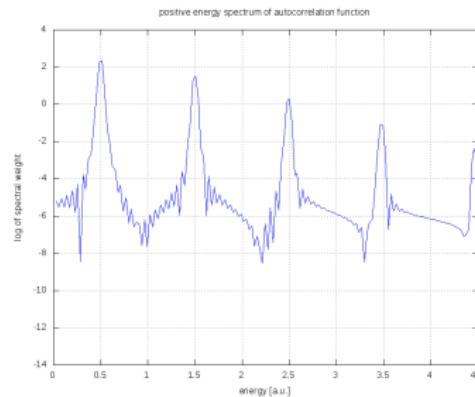
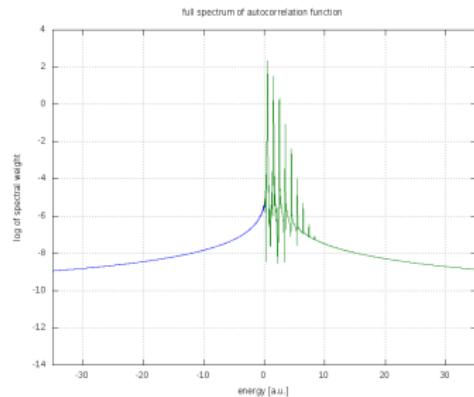
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Exercise 6: Getting familiar with real-time propagators

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Properties of approximate real-time propagators

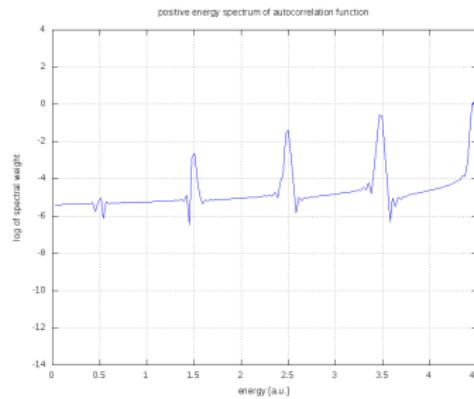
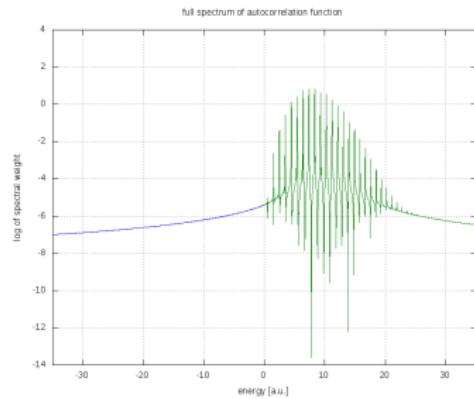
Exercise 6: RESULTS - autocorrelation function



xmin = 1.0, squeeze = 1.2, 9 spectral peaks

Properties of approximate real-time propagators

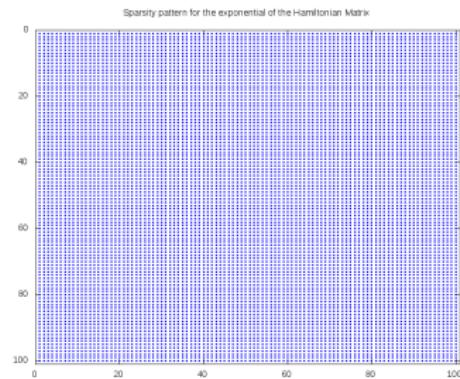
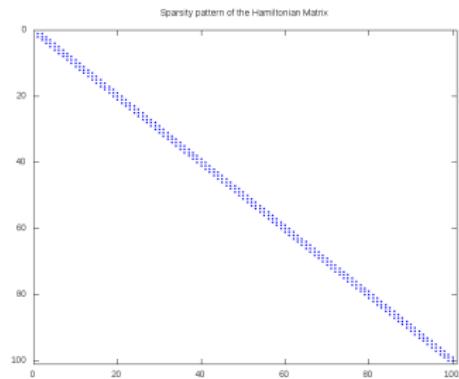
Exercise 6: RESULTS - autocorrelation function



`xmin = 4.0, squeeze = 1.2, 28 spectral peaks`

Properties of approximate real-time propagators

Exercise 6: RESULTS - sparsity of matrices



Properties of approximate real-time propagators

Exercise 6: RESULTS - Runge-Kutta 4th order

```
rk1 = -sqrt(-1) * dt * Hm_t * (psi);
rk2 = -sqrt(-1) * dt * Hm_t_dt2 * (psi + rk1/2);
rk3 = -sqrt(-1) * dt * Hm_t_dt2 * (psi + rk2/2);
rk4 = -sqrt(-1) * dt * Hm_t_dt * (psi + rk3);
psi = psi + (rk1 + 2*rk2 + 2*rk3 + rk4)/6;
```

Properties of approximate real-time propagators

Exercise 6: RESULTS - Quiz time

- Which method allows for the largest time step?
 - ▶ Crank-Nicholson/Caley
- Which method allows for the fastest time stepping?
 - ▶ Runge-Kutta 4th order, Taylor expansion
- How is the grid spacing related to the time step?
 - ▶ The spacings appear in the ratio $\Delta t / \Delta x^2$ in the exponential

Absorption spectra from TDDFT

Exercise 7:

- Calculation of photoabsorption cross section for N_2 , H_2O