

Algorithms: Total Energy Minimization

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Goal: Calculate total energy

$$E_{tot}^{BO}(\{\vec{R}_i\}) = \min_{\{\phi_i\}} E_{tot}[\{\vec{R}_I\}, \{\phi_i\}]$$

Questions:

- How to construct E_{tot} ?
- **How to find (efficiently) the electronic ground state?**
⇒ Solve Kohn-Sham Equations

$$H[n]\phi_i = \varepsilon_i\phi_i \quad \text{with} \quad n(\vec{r}) = \sum_i |\Psi_i(\vec{r})|^2$$

Outline

How to get to the Born-Oppenheimer surface?

- **Direct diagonalization**
- **Iterative diagonalization**
 - **Method**
 - **Scaling**
 - **Improvements**
 - **Self-consistency**

How to move on the Born-Oppenheimer surface?

- **Equilibrium geometry (T=0K)**
- **Molecular dynamics**

Basis set representation

Basis set expansion: $|\varphi_i\rangle = \sum_{\mu} c_{i\mu} |\mu\rangle$ $\langle\mu|\nu\rangle = \delta_{\mu\nu}$ **orthogonal**

$\sum_{\mu} |\mu\rangle\langle\mu| = \hat{1}$ **complete**

KS-equations: $H \varphi_i = \varepsilon_i \varphi_i \rightarrow \sum_{\mu} \langle\nu|H|\mu\rangle\langle\mu|\varphi_i\rangle = \varepsilon_i \langle\nu|\varphi_i\rangle$

\Rightarrow **Solve eigenvalue problem:** $\hat{H} \vec{Y} = \varepsilon \vec{Y}$

direct diagonalization: $\hat{U}^T \hat{H} \hat{U} = \hat{1} \vec{\varepsilon}$

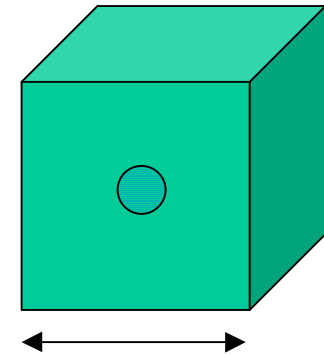
$$\begin{pmatrix} U_{11} & \cdots & \cdots & U_{N1} \\ \vdots & \ddots & & \vdots \\ \vdots & & \ddots & \vdots \\ U_{1N} & \cdots & \cdots & U_{NN} \end{pmatrix} \begin{pmatrix} H_{11} & \cdots & \cdots & H_{1N} \\ \vdots & \ddots & & \vdots \\ \vdots & & \ddots & \vdots \\ H_{N1} & \cdots & \cdots & H_{NN} \end{pmatrix} \begin{pmatrix} U_{11} & \cdots & \cdots & U_{1N} \\ \vdots & \ddots & & \vdots \\ \vdots & & \ddots & \vdots \\ U_{N1} & \cdots & \cdots & U_{NN} \end{pmatrix} = \begin{pmatrix} \varepsilon_{11} & \cdots & \cdots & 0 \\ \vdots & \ddots & & \vdots \\ \vdots & & \ddots & \vdots \\ 0 & \cdots & \cdots & \varepsilon_{NN} \end{pmatrix}$$

- Problems:**
- N^3 scaling
 - complete Hamiltonian has to be saved

Discussion: Plane wave basis

Example: cubic cell with $a=5\text{\AA}$

| E_{cut} | n_{max} | N_{PW} |
|------------------|------------------|-----------------|
| 5 | 14 | 534 |
| 10 | 20 | 1510 |
| 15 | 25 | 2774 |
| 20 | 28 | 4271 |
| 40 | 40 | 12081 |
| 100 | 64 | 47752 |



Realistic systems have basis sets with $10^{4\dots 6}$ functions!

\Rightarrow direct diagonalization inefficient/impossible

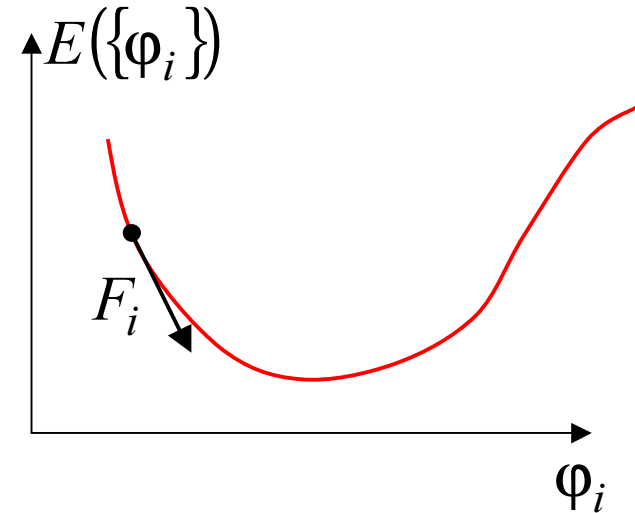
Iterative diagonalization

Basic idea:

Total energy gradient:

$$F_i = -\frac{\delta E(\{\varphi_i\})}{\delta \varphi_i} \quad \text{with} \quad \langle \varphi_i | \varphi_j \rangle = \delta_{ij}$$

Residuum: $|F_i\rangle = (H - \varepsilon_i)|\varphi_i\rangle$



⇒ Find equation of motion (EOM) which minimizes total energy!

First order EOM: $-|F_i\rangle = \alpha|\dot{\varphi}_i\rangle$ α “friction” parameter

or: $(H - \varepsilon)|\varphi_i\rangle = -\alpha|\dot{\varphi}_i\rangle$

interesting aspect: time dependent Schrödinger equation for $\alpha \rightarrow i\hbar$

Scaling of the iterative algorithm

Equation of motion: $(H - \varepsilon)|\varphi_i\rangle = -\alpha|\dot{\varphi}_i\rangle = -\alpha\frac{\varphi_i^{(n+1)} - \varphi_i^{(n)}}{\Delta t} = -\tilde{\alpha}\Delta\varphi_i^{(n)}$

Plane wave basis set: $\sum_{G'} \langle G|(H - \varepsilon)|G'\rangle \langle G'|\varphi_i\rangle = -\tilde{\alpha}\langle G|\Delta\varphi_i^{(n)}\rangle$

$O(N^2M)$ operations
 M number of states ($\ll N$)

Trick: Use locality of the contributions of the Hamiltonian in real and reciprocal space!

Hamilton Operator: $H = -\nabla^2 + V^{eff}(\vec{r})$ real space $\xleftrightarrow{\text{FFT}}$ reciprocal space
 $O(N \ln(N))$ operations

local in reciprocal space

$$\begin{pmatrix} -G_1^2 & \dots & \dots & 0 \\ \vdots & \ddots & & \vdots \\ \vdots & & \ddots & \vdots \\ 0 & \dots & \dots & -G_N^2 \end{pmatrix} \begin{pmatrix} \varphi_i(G_1) \\ \vdots \\ \vdots \\ \varphi_i(G_N) \end{pmatrix}$$

$O(N)$ operations

local in real space

$$\begin{pmatrix} V^{eff}(\vec{r}_1) & \dots & \dots & 0 \\ \vdots & \ddots & & \vdots \\ \vdots & & \ddots & \vdots \\ 0 & \dots & \dots & V^{eff}(\vec{r}_N) \end{pmatrix} \begin{pmatrix} \varphi_i(\vec{r}_1) \\ \vdots \\ \vdots \\ \varphi_i(\vec{r}_N) \end{pmatrix}$$

$O(N)$ operations

Construction of the initial wave functions

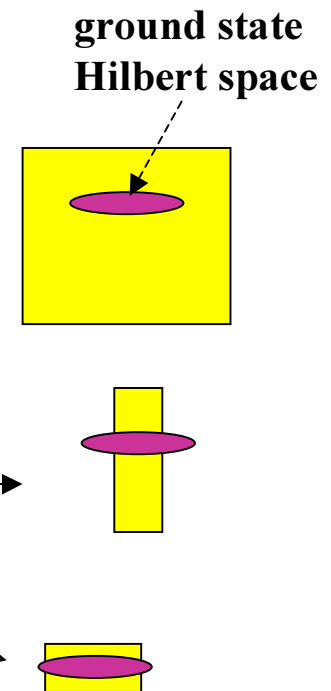
Iterative schema: $\varphi_i^{(n+1)} = F\left(\left\{\varphi_i^{(n)}\right\}\right)$

Problem: How to construct $\left\{\varphi_i^{(0)}\right\}$?

$$\Psi_{\vec{k}}(\vec{r}) = \sum_{\vec{G}} \underline{\underline{c_{\vec{k}}(\vec{G})}} e^{i(\vec{G}+\vec{k})\vec{r}}$$

Solutions:

- Direct diagonalization (not possible for large systems)
- Random numbers
- Direct diagonalization in PW subset
- Direct diagonalization in LCAO basis set



Diagonalization in LCAO basis set

Expand atomic orbitals in plane waves:
$$\mu_{\vec{k}}(\vec{r}) = \sum_{\vec{G}} \mu_{\vec{k}}(\vec{G}) e^{i(\vec{G}+\vec{k})\vec{r}}$$

μ ... atomic s, p, d orbitals

Formally fully equivalent to wave functions:
$$\Psi_{\vec{k}}(\vec{r}) = \sum_{\vec{G}} c_{\vec{k}}(\vec{G}) e^{i(\vec{G}+\vec{k})\vec{r}}$$

Construct hamiltonian and overlap matrix:
$$H_{\mu\nu}(\vec{k}) = \langle \mu_{\vec{k}} | \hat{H} | \nu_{\vec{k}} \rangle$$
$$S_{\mu\nu}(\vec{k}) = \langle \mu_{\vec{k}} | \nu_{\vec{k}} \rangle$$

Solve generalized eigenvalue problem:
$$\left\{ H_{\mu\nu}(\vec{k}) - \varepsilon_i S_{\mu\nu}(\vec{k}) \right\} \langle \mu_{\vec{k}} | \psi_i \rangle = 0$$

Number of atomic orbitals \ll Number of plane waves (approx. 1:100)

 Eigenvalue problem can be easily solved!

Convergence Criteria (I)

iterative solution: $\varphi^{(n+1)} = \varphi^{(n)} - (H - \varepsilon)|\varphi^{(n)}\rangle\Delta t$ $\varepsilon = \langle\varphi^{(n)}|H|\varphi^{(n)}\rangle$

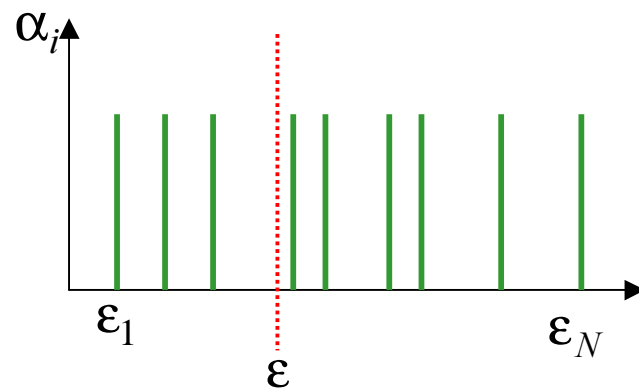
exact solution: $(H - \varepsilon)|\chi_i\rangle = 0$

representation of the trial wave function in the χ_i set:

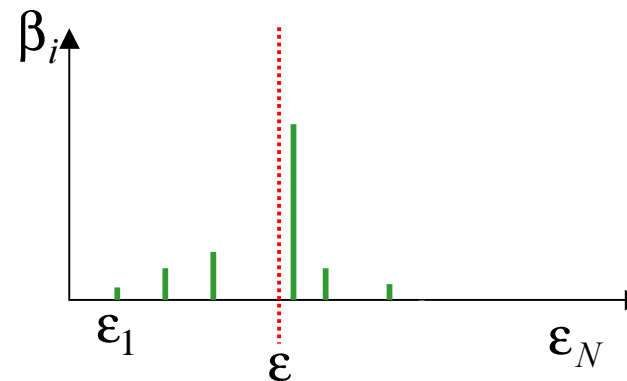
$$|\varphi^{(n)}\rangle = \sum_i \alpha_i |\chi_i\rangle \quad |\varphi^{(n+1)}\rangle = \sum_i \beta_i |\chi_i\rangle$$

→ $\beta_i = \alpha_i - (\varepsilon - \varepsilon_i)\Delta t$

Before iteration step:

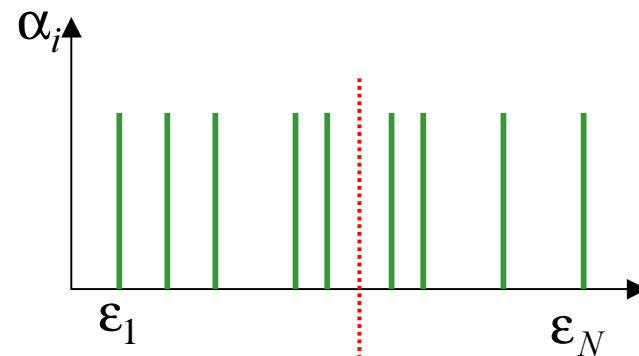


After “ideal” iteration step:



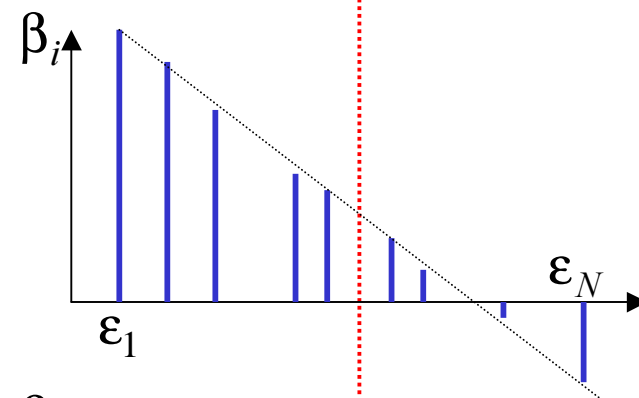
Convergence Criteria (II)

Before iteration step:



EOM: $\beta_i = \alpha_i - (\epsilon - \epsilon_i)\Delta t$

After iteration step:

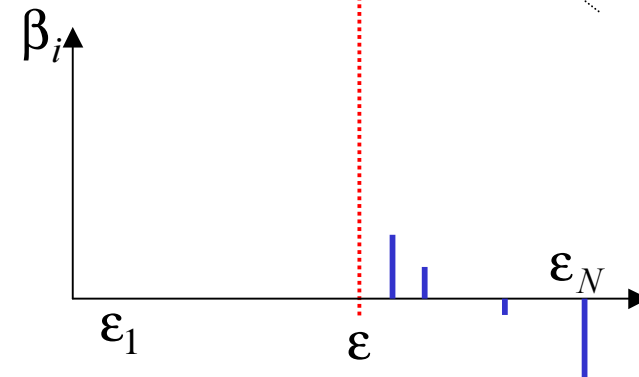


optimum time step:

$$0 \leq \beta_N = \alpha_N - (\epsilon - \epsilon_N)\Delta t_{\max}$$

$$\Delta t_{\max} = \frac{1}{\epsilon_N - \epsilon_1} = \frac{1}{\underbrace{\epsilon_{\max} - \epsilon_{\min}}_{\text{spectral radius}}}$$

After orthogonalization:



- **Orthogonalization essential for each iteration step**
- **Convergence rate decreases with increasing number of PW**

Preconditioning

Residual error: $(H - \varepsilon) \underbrace{|\chi + \Delta\varphi\rangle}_{=: |\varphi\rangle} = (H - \varepsilon) |\Delta\varphi\rangle = |\Delta\tilde{\varphi}\rangle$

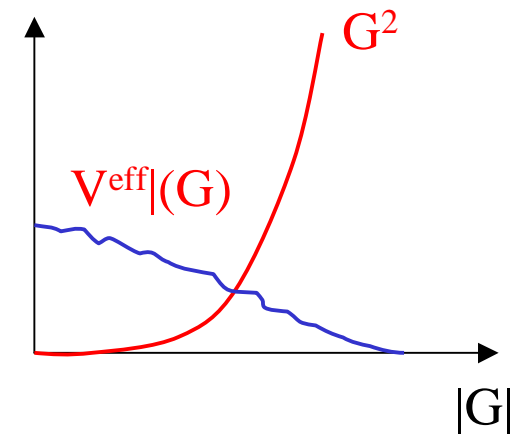
we need this
we get this

→ $(H - \varepsilon)^{-1} |\Delta\tilde{\varphi}\rangle = |\Delta\varphi\rangle$

Hamiltonian in PW basis:

$$H_{GG'} = -G^2 \delta_{GG'} + V^{eff}(G - G')$$

Kinetic energy dominates for high wave numbers!:



Partition Hamiltonian:

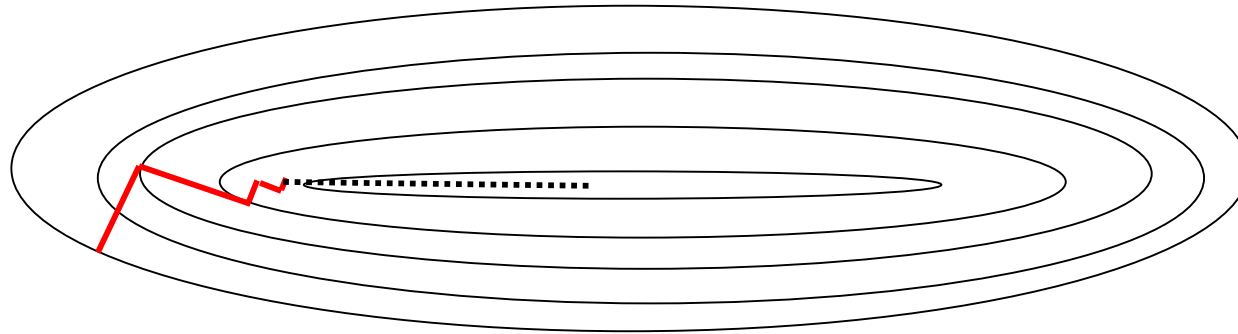
$$\hat{H} = \hat{D} + \hat{L}$$

diagonal matrix
all diagonal elements are zero

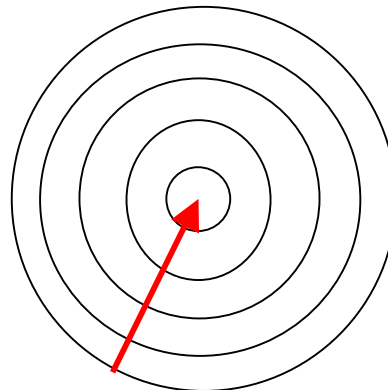
For high wave numbers: $\hat{H} \approx \hat{D}$ **→** $(\hat{D} - \varepsilon)^{-1} \Delta\tilde{\varphi} = \Delta\varphi$

Preconditioning: Geometric interpretation

Without preconditioning:



With preconditioning:



Williams-Soler Algorithm

Equation of motion: $(H - \varepsilon_i)|\phi_i\rangle = -\alpha|\dot{\phi}_i\rangle$

In plane-wave basis set:

$$\alpha \dot{c}_i(G) = \underbrace{\left[G^2 + V^{eff}(G-G) \right]}_{\text{diagonal part}} c_i(G) - \sum_{\substack{G' \\ G' \neq G}} \underbrace{V^{eff}(G-G')}_{\text{off-diagonal part}} c_i(G')$$

EOM for a single state: $\dot{c}_i(G) = \omega_G^2 c_i(G) - B_G$

Assumption: B_G not time dependent

$$\Rightarrow c_i(G, t = \Delta t) = -\frac{B_G}{\omega_G^2} + \left[c_i(G, t = 0) + \frac{B_G}{\omega_G^2} \right] \exp(-\omega_G^2 \Delta t)$$

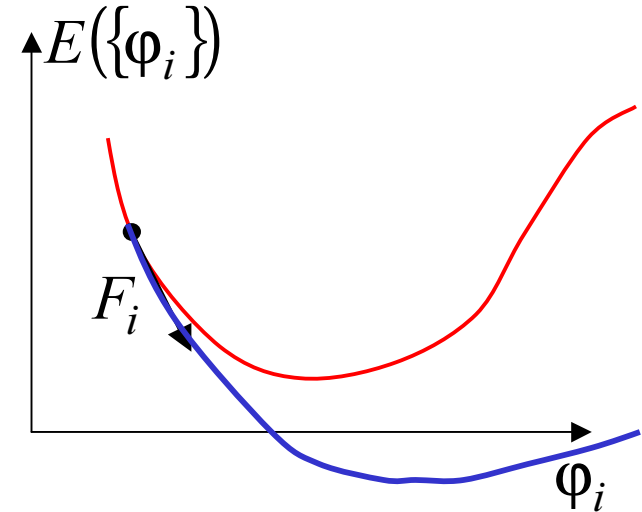
Improves significantly the convergence rate for high PW energy cutoffs!

Higher order EOM's

Goal: Find total-energy minimum more efficiently

Strategies:

- minimize number of iterations
- minimize computational effort for each step



Two approaches

Use higher derivatives

$$D_{ij} = \frac{\partial^2 E(\{\varphi_i\})}{\partial \varphi_i \partial \varphi_j}, \dots$$

high computational effort!

Use higher order in time

$$\varphi_i = F(\dot{\varphi}_i, \ddot{\varphi}_i, \dots)$$

equivalent:

$$\varphi_i^{(n+1)} = \tilde{F}(\varphi_i^{(n)}, \varphi_i^{(n-1)}, \dots)$$

no additional computational effort!

Second order equation of motion

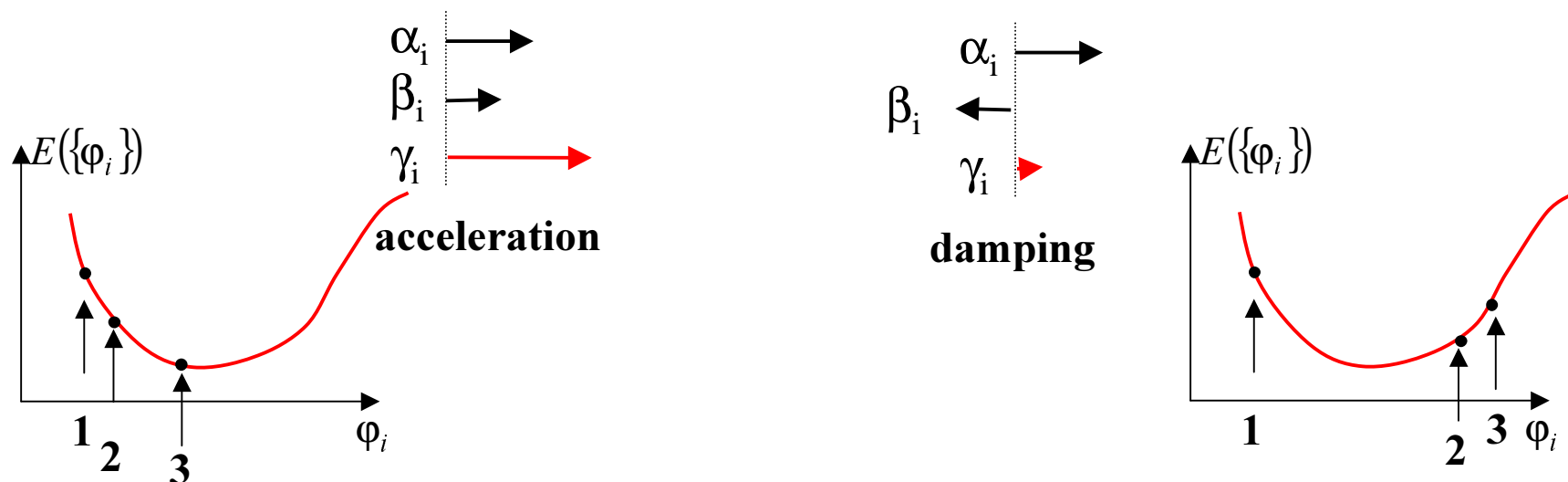
Finite differences:
$$\varphi^{(n+1)} = \varphi^{(n)} - (H - \varepsilon) \left| \varphi^{(n)} \right\rangle \Delta t_1 - (H - \varepsilon) \left| \varphi^{(n-1)} \right\rangle \Delta t_2$$

Expansion with respect to eigenfunctions χ_i :

$$\left| \varphi^{(n-1)} \right\rangle = \sum_i \alpha_i \left| \chi_i \right\rangle \quad \left| \varphi^{(n)} \right\rangle = \sum_i \beta_i \left| \chi_i \right\rangle \quad \left| \varphi^{(n+1)} \right\rangle = \sum_i \gamma_i \left| \chi_i \right\rangle$$

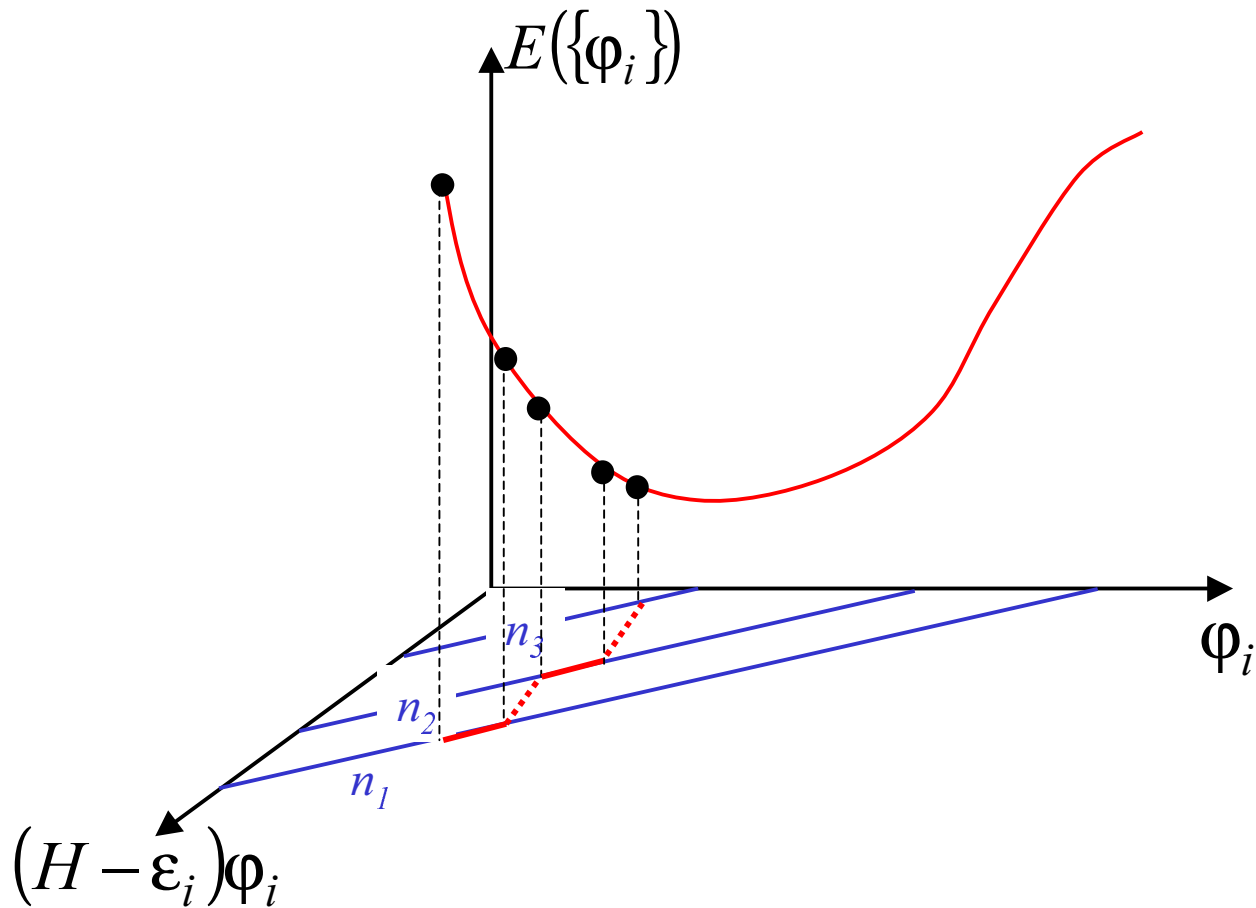
→
$$\gamma_i = \beta_i - \beta_i (\varepsilon - \varepsilon_i) \Delta t_1 - \alpha_i (\varepsilon - \varepsilon_i) \Delta t_2$$

Two possible scenarios



Direct vs iterative methods

Self-consistency: $(H - \varepsilon_i)|\varphi_i\rangle = 0$ with $H = H(\{\varphi_i\})$



Improved convergence rate by enforcing self-consistency at each iteration step!

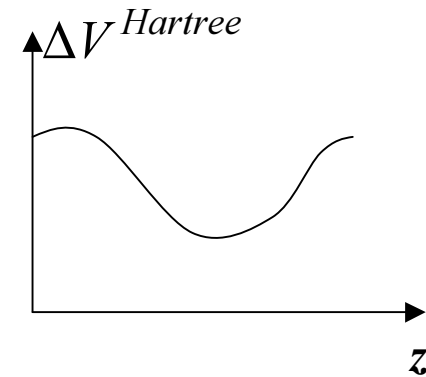
Instabilities: Charge sloshing

Hartree potential:

$$V^{Hartree}(G) = \frac{4\pi e\rho(G)}{|G|^2}$$

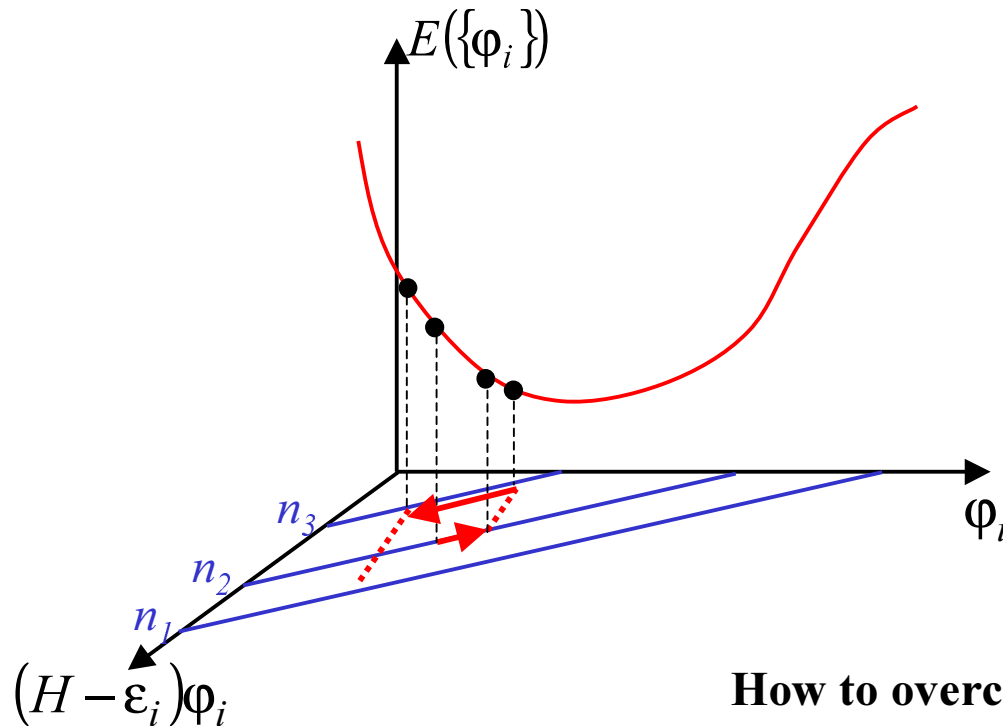
Change in Hartree potential:

$$\Delta V^{Hartree}(G) = \frac{4\pi e\Delta\rho(G)}{|G|^2}$$



Instability increases with increasing system size:

$$G_{\min} \propto \frac{1}{l_{\max}}$$



How to overcome the instabilities?

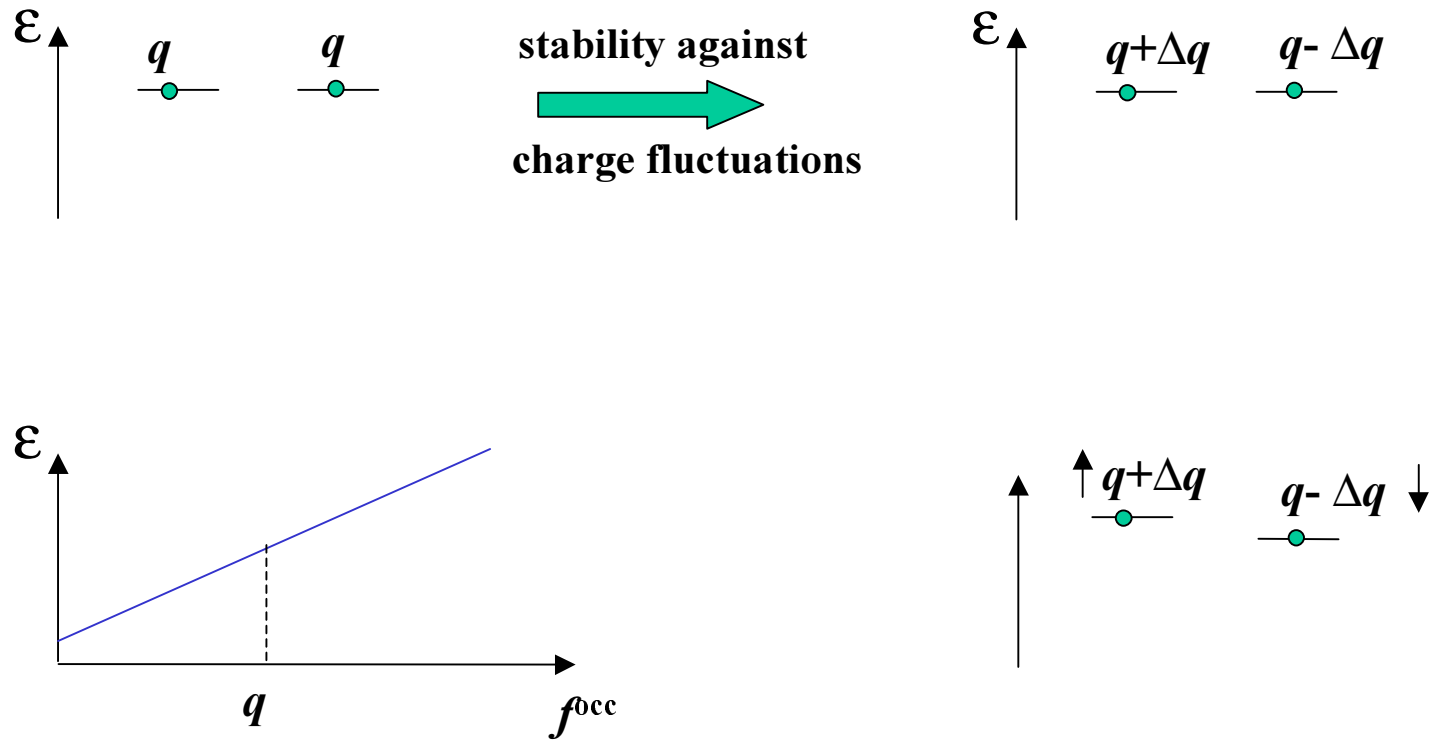
- Use higher order algorithms
- Mixing of the charge density, e.g.

$$\rho(\vec{r}) = \lambda_{mix}\rho^{old}(\vec{r}) + (1 - \lambda_{mix})\rho^{new}(\vec{r})$$

Instabilities: Metallic systems (I)

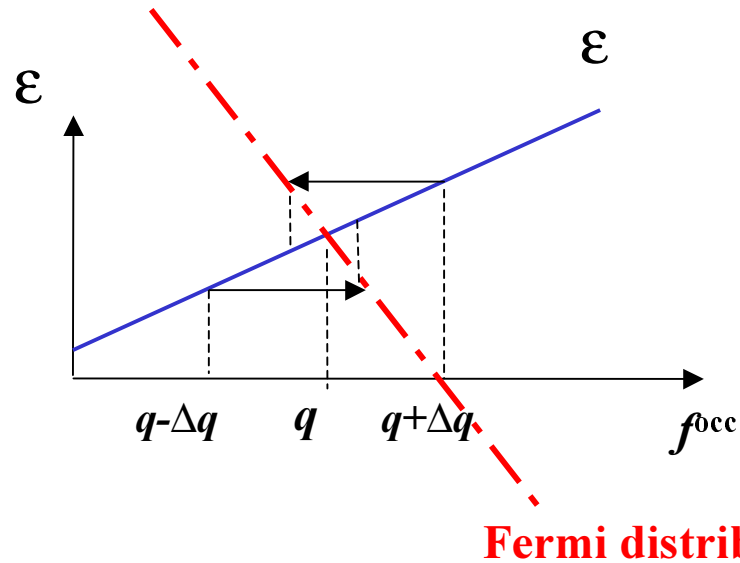
Fractionally occupied states:

Consider two degenerate states close to the Fermi level:



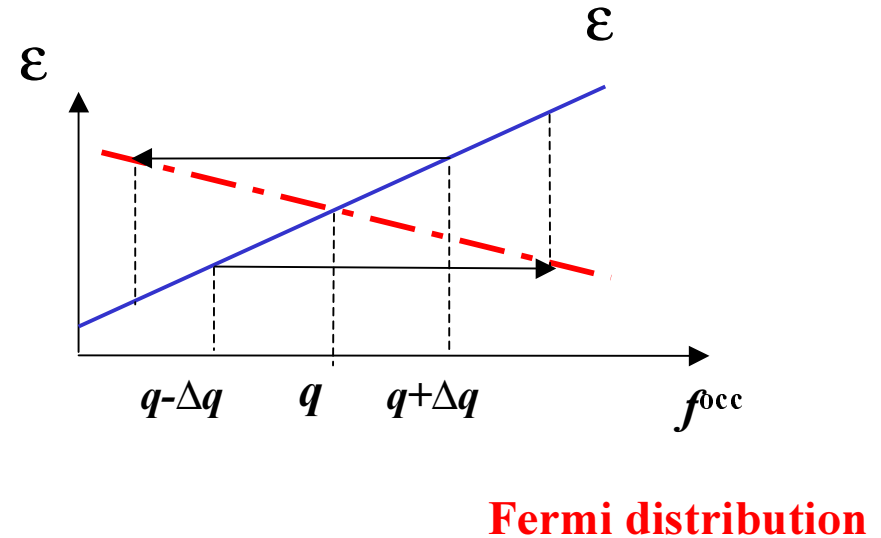
Instabilities: Metallic systems (II)

High electronic temperatures:



Stable against charge fluctuations!

Low electronic temperatures:

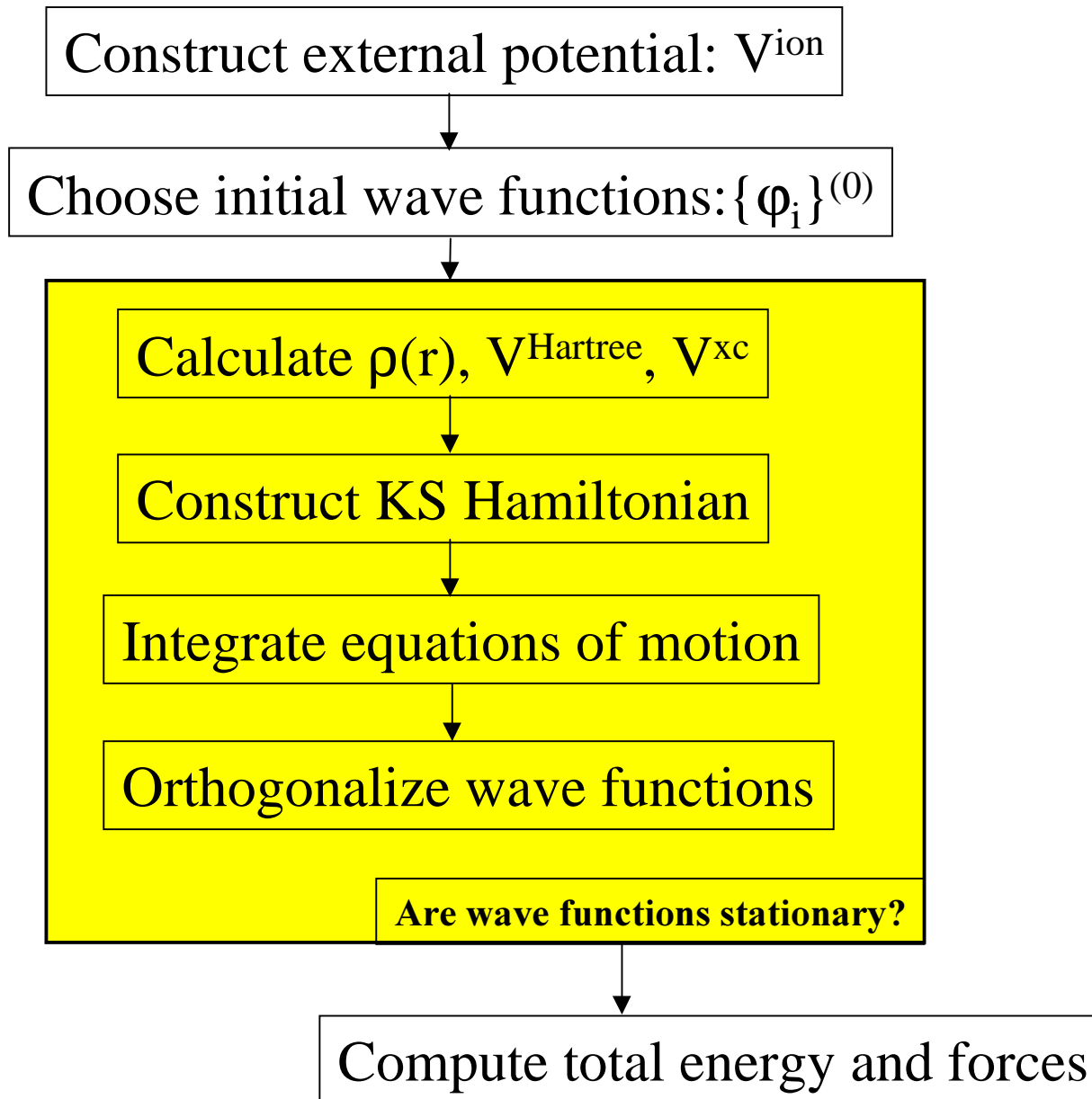


Unstable against charge fluctuations!

How can we avoid these instabilities?

- Improve k-point sampling
- damp occupation numbers

How to get down to the BO surface



Hellmann-Feynman forces

Hellmann-Feynman theorem:

$$\vec{F}_I = -\frac{\partial E(\{\vec{R}_I\})}{\partial \vec{R}_I}$$
$$= -\sum_i \langle \varphi_i | \frac{\partial H}{\partial \vec{R}_I} | \varphi_i \rangle$$

This theorem is valid only when: $(H - \varepsilon_i)|\varphi_i\rangle = 0$

Forces are not variational with respect to the wave functions
 \Rightarrow good convergence required

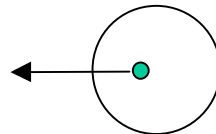
Ionic motion

Basic idea: Simultaneous motion of electrons and ions per iteration step
fictitious dynamics for electrons keeps wave functions on BO surface

Car-Parrinello Lagrangian:
$$L = \sum_i \frac{1}{2} \mu \langle \dot{\phi}_i | \dot{\phi}_i \rangle + \sum_I \frac{1}{2} M_I \dot{\vec{R}}_I^2 - E[\{\phi_i\}, \{\vec{R}_I\}]$$

Equations of motions:
$$M_I \ddot{\vec{R}}_I = - \frac{\partial E[\{\phi_i\}, \{\vec{R}_I\}]}{\partial \vec{R}_I} \quad \mu \ddot{\phi}_i = - \frac{\partial E[\{\phi_i\}, \{\vec{R}_I\}]}{\partial \phi_i}$$

First order EOM:



Second order EOM:



How do we move the atoms

Two schemes

Molecular dynamics

Integrate equations of motion:

$$M_I \ddot{\vec{R}}_I = \vec{F}_I$$

Structure optimization

2nd order EOM with damping

$$\mu_I \ddot{\vec{R}}_I + \lambda_I \dot{\vec{R}}_I = \vec{F}_I$$

effective mass

damping/friction