# The FHI-aims All-Electron Framework for Molecular and Materials Simulations

#### <u>Volker Blum</u>

Department of Mechanical Engineering & Materials Science - Duke University, Durham, NC

http://aims.pratt.duke.edu



Frontiers of Advanced Electronic Structure Methods - Hefei, June 16, 2016

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All-Electron Theory for Large Systems: FHI-aims



Extending the Reach of DFT & Many-Body Theory

Tunable Electronic Structure in Organic-Inorganic Hybrids



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### Many Individuals Contributed to This Work - Thanks!



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Dr.William Huhn Spin-Orbit Coupling "Energy Materials"



Victor Yu Electronic Structure Infrastructure (NSF)

Björn Lange Jan Kloppenburg Tiago Botari





Dr. Raul Laasner Nuclear Spin States & NMR



Tong Zhu PV Materials GW & RPA



Garnett Liu Perovskites Excitonic Effects



FHI-aims team and collaborators: Matthias Scheffler (Berlin), Xinguo Ren (Hefei), over 100 individuals with contributions to the project. Development continues in Berlin, Hefei, Munich, Helsinki, London, Duke, Argonne, and many more.

### So, We Have a Theory to Model "Everything"

Schrödinger (Dirac) Equation

 $\hat{\mathcal{H}}\Psi = E\Psi$ 



P.A.M. Dirac

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<u>Perfect recipe for parameter-free modeling ...</u> ... but how do we make it practical?

P.A.M. Dirac





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 $\hat{\mathcal{H}}\Psi = E\Psi$ 



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P.A.M. Dirac

The approximation is accurate but <u>feasible systems are too small</u>



Feasible systems large enough but <u>approximations are not great</u>

(How?) can we work towards having both?



Quantum chemistry & many-body theory:



 $E_{tot} \leq \langle \psi | H | \psi \rangle$  ... successive refinement of  $\psi$ 

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Density functional theory: (Hohenberg-Kohn 1964, Kohn-Sham 1965)

 $E_{tot} = E[n(r)] = T_s[n] + V[n] + V_{es}[n] + E_{xc}[n]$ 

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• Key practical approximation:  $E_{xc}$ 



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"Perdew's ladder"

to exact solution

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cost, accurac) • Key practical approximation:  $E_{xc}$ response / many-body terms: RPA, SOSEX, ... hybrid functionals: non-local exchange meta-GGAs:  $\nabla^2 n(r)$ ,  $\nabla^2 \phi(r)$ Generalized gradient approximations (GGAs):  $|\nabla n(r)|$ Local-density approximation (LDA): n(r)

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<u>Universality:</u> Cover full space of materials and chemistry

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 <u>Accuracy:</u> Hierarchy of methods: DFT and "beyond" (GW, MP2, RPA, rPT2, ...)
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Accuracy: Hierarchy of methods: DFT and "beyond" (GW, MP2, RPA, rPT2, ...) All-electron

Efficiency: Scalable (system size, number of CPUs)

Efficient (1,000s of atoms), but do not sacrifice accuracy!

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"Basis set expansion":  $\psi_k(m{r}) = \sum_i c_{ki} \varphi_i(m{r})$ 

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$$\begin{split} h_{ij} &= \langle \varphi_i | \hat{h}_{\rm KS} | \varphi_j \rangle \\ s_{ij} &= \langle \varphi_i | \varphi_j \rangle \end{split}$$

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• Plane waves (VASP, abinit, ...)

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 <u>Many</u> others: "Augmented plane waves", wavelets, finite elements,

numeric atom-centered functions, ...

Our choice!  $\varphi_{i[lm]}(\boldsymbol{r}) = \frac{u_i(r)}{r} \cdot Y_{lm}(\Omega)$ 

### Numeric Atom-Centered Basis Functions: Some Advantages

$$arphi_{i[lm]}(oldsymbol{r}) = rac{u_i(r)}{r} \cdot Y_{lm}(\Omega)$$



Many popular implementations: DMol<sup>3</sup> (Delley), FPLO (Eschrig *et al.*) PLATO (Horsfield *et al.*), ADF, BDF, PAOs (Siesta, Conquest, OpenMX<sup>2</sup>, Fireball, ...)

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- "Naturally" all-electron
- Strictly localizable O(N) computational scaling for most expensive steps
- Rather compact basis sets for converged results

We have a basis set library for all elements (I-I02), from fast qualitative to meV-converged calculations. (total energy, DFT).

Blum, Gehrke, Hanke, Havu, Havu, Ren, Reuter, Scheffler, Computer Physics Communications 180, 2175 (2009)

#### Accuracy in Community Wide Benchmark - "Delta Test"

<u>S. Cottenier and coworkers (Ghent University):</u> <u>https://molmod.ugent.be/deltacodesdft</u>

E(V) for 71 elemental solids - Reference: Full-Potential LAPW (Wien2k).

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	Compari	ng Solid State DFT Codes, Basis Se	ets and Potentials   Center for Molecular Mod	eling	
+ 🔝 https	🔒 molmod.ug	ent.be/deltacodesdft			C Reader
	Con	nparing Solid State DFT Codes, Basis Set	s and Potentials   Center for Molecular Modeling		
Code	Version		Electron treatment	∆-value	Authors
WIEN2k	13.1	$\Delta = \frac{1}{\Delta}$	0.6 meV	0 meV/atom	S. Cottenier
FHI-aims 🗗	081213	tier2 numerical orbitals	all-electron (relativistic atomic_zora scalar)	0.2 meV/atom	ASE [2]
Exciting 🗗	developi version	FHI-aims, really 1 $\Delta=0$	tight / <i>tier</i> 2 basis sets .2 meV (!)	• 0.2 meV/atom	Exciting [10]
FHI-aims 🗗	081213	Dr. Mar	rcin Dulak, DTU	2) 0.4 meV/atom	ASE [2]
CASTEP	<sup>8.0</sup>	plane waves producibility in Density	-Functional Calculations of Soli	0.5 ds, meV/atom	CASTEP [7]
ABINIT 🗗	7.7.3	K. Lejaeghere, 68 Science <b>35</b>	coauthors!, S. Cottenier, I, aad3000 (2016).	0.6 meV/atom	F. Jollet and M Torrent
FHI-aims 🗗	081213	tight numerical orbitals	all-electron (relativistic atomic_zora scalar)	0.6 meV/atom	ASE [2]
VASP	5.2.12	plane waves	PAW 2012	0.7 meV/atom	K. Lejaeghere

### Similar High Precision Achievable for Band Structures



# All-Electron Accuracy for Large Systems: FHI-aims



http://aims.fhiberlin.mpg.de

- <u>Accurate</u>, efficient quantum mechanics for molecules, materials
- Numerical Foundation: <u>Numeric atom-centered basis functions</u> Seamless from "light" to basis-converged results, from light to heavy elements
- Non-periodic and periodic structure models, <u>same framework</u>
- Density Functional Theory (semilocal, global and rangeseparated hybrid functionals, van der Waals corrections)
- Beyond DFT (non-periodic): RPA, GW, MP2, TD-LDA, ...
- Scalable (1,000s of atoms on 1(00),000s of CPU cores)
- <u>Properties</u> (charged and neutral excitations, IR spectra, transport, dielectric functions, dynamics & transition states, ...)

~100 contributors to date, over 100 licensing groups. Active development in Berlin, Duke, Munich, Hefei, Helsinki, London, Argonne, ... New stable release "160328"

Blum, Gehrke, Hanke, Havu, Havu, Ren, Reuter, Scheffler, Computer Physics Communications **180**, 2175(2009) Ren, Rinke, Blum, Wieferink, Tkatchenko, Sanfilippo, Reuter, Scheffler, New J. Phys. **14**, 053020 (2012)

### ... But How to Push an All-Electron Approach to Large Systems?

I. Kohn-Sham DFT Eigenvalue Problem  $(O(N^3))$ 

$$\underbrace{\underline{h}}\underline{c}_{k} = \epsilon_{k} \underline{\underline{s}} \underline{c}_{k}$$

2. Exact Exchange & Hybrid Functionals

$$E_{\mathbf{x}}^{\mathrm{HF}} = \frac{1}{2} \sum_{ij\sigma} D_{ij}^{\sigma} K_{ij}^{\sigma} = \frac{1}{2} \sum_{ijkl} D_{ij}^{\sigma} D_{kl}^{\sigma} (ik|lj)$$

3. Many-Body Perturbation Theory "Beyond DFT"

$$\sum_{\sigma}^{GW}(\mathbf{r},\mathbf{r}',\epsilon) = \frac{i}{2\pi} \int d\epsilon' \left[ G_{\sigma}(\mathbf{r},\mathbf{r}',\epsilon+\epsilon') W(\mathbf{r},\mathbf{r}',\epsilon') e^{i\eta\epsilon'} \right]$$

Cost Accuracy? Difficulty!

# Scalability: 3-Layer Graphene on 3C-SiC(0001)



Structure: Nemec et al., PRL **III**, 065502 (2013).

Benchmark: Björn Lange, Duke University 2014

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### Parallel Eigenvalue Solvers - the Problem



IBM BlueGene (MPG, Garching) 16384 CPU cores

### Parallel Eigenvalue Solvers - the Problem





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### Parallel Eigenvalue Solvers - the Problem






#### Two-stage reduction algorithm, ELPA:



Larger fraction of efficient matrix-matrix operations Efficient compute kernels for added backtransform step ④

Note: Hard to beat dense linear algebra for small to midsized problems, many EVs

Auckenthaler, Blum, Bungartz, Huckle, Johanni, Krämer, Lang, Lederer, Willems, Parallel Computing 37, 783 (2011) A. Marek, V. Blum, R. Johanni, V. Havu, B. Lang, T. Auckenthaler, A. Heinecke, H.-J. Bungartz, H. Lederer, The Journal of Physics: Condensed Matter 26, 213201 (2014).

#### ELPA, Two-Step Solver



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I) Further reduce prefactor of dense linear algebra based approach?
 → new hardware? (GPU, MIC)

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 $O(N^3)$  approach focused on density matrix, naturally iterative in s.c.f.  $\rightarrow$  lower prefactor Recently reimplemented by Corsetti, Comput. Phys. Commun. 185, 273 (2014).

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All these are feasible but optimal choice depends on <u>system size</u>, <u>system character</u>, <u>required electronic structure output</u>, ... complex. Can we simplify this task?

#### ... what to do?

I) Further reduce prefactor of dense linear algebra based approach?
 → new hardware? (GPU, MIC)

2) Switch to "traditional" O(N) approach (density matrix based)?
 → well understood, but still some system-specificity

 Orbital minimization method
 O(N<sup>3</sup>) approach focused on density matrix, naturally iterative in s.c.f. → lower prefactor Recently reimplemented by Corsetti, Comput. Phys. Commun. 185, 273 (2014).

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# ELSI: Solving or Circumventing the Eigenvalue Problem

<u>... in one infrastructure</u> <u>& for "any" code?</u>

... and many other "stakeholders" from the community



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Fabiano Corsetti,

Lin Lin, Berkeley London





Alvaro Vazquez-Mayagoitia, ANL

Work in Progress:

"ELSI" - Electronic Structure Infrastructure (NSF-SI2)



#### http://elsi-interchange.org

... and many other "stakeholders" from the community



#### ... But How to Push an All-Electron Approach to Large Systems?

I. Kohn-Sham DFT Eigenvalue Problem  $(O(N^3))$ 

$$\underbrace{\underline{h}}\underline{c}_{k} = \epsilon_{k} \underline{\underline{s}} \underline{c}_{k}$$

2. Exact Exchange & Hybrid Functionals

$$E_{\mathbf{x}}^{\mathrm{HF}} = \frac{1}{2} \sum_{ij\sigma} D_{ij}^{\sigma} K_{ij}^{\sigma} = \frac{1}{2} \sum_{ijkl} D_{ij}^{\sigma} D_{kl}^{\sigma} (ik|lj)$$

3. Many-Body Perturbation Theory "Beyond DFT"

$$\sum_{\sigma}^{GW}(\mathbf{r},\mathbf{r}',\epsilon) = \frac{i}{2\pi} \int d\epsilon' \left[ G_{\sigma}(\mathbf{r},\mathbf{r}',\epsilon+\epsilon') W(\mathbf{r},\mathbf{r}',\epsilon') e^{i\eta\epsilon'} \right]$$

Cost Accuracy? Difficulty!

### Beyond DFT-LDA/GGA: Two-Electron Integrals

<u>Common bottleneck: Two-electron interactions, e.g.:</u>

$$\int d^3r d^3r' \phi_i(r) \phi_j(r') \frac{1}{|r-r'|} \phi_k(r') \phi_l(r)$$

Ren, Rinke, Blum, Wieferink, Tkatchenko, Sanfilippo, Reuter, Scheffler, New J. Phys. 14, 053020 (2012)

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$$\int d^3r d^3r' \phi_i(r) \phi_j(r') \frac{1}{|r-r'|} \phi_k(r') \phi_l(r)$$

Early Solution: "Resolution of the identity"

[Boys and Shavitt (1950s), Whitten (1974), Dunlap et al. (1979), Vahtras et al. (1993), many others]

<u>I. Expand pair products in</u> smaller auxiliary basis set  $\{P_{\mu}\}$ 

$$arphi_i(m{r})arphi_j(m{r}) = \sum_\mu C^\mu_{ij} P_\mu(m{r})$$

2. Use "Coulomb metric" V to cancel linear error terms

$$C_{ij}^{\mu} = \sum_{\nu} (ij|\nu) V_{\nu\mu}^{-1}$$
$$(ij|\nu) = \int d^3r d^3r' \frac{\varphi_i(\boldsymbol{r})\varphi_j(\boldsymbol{r}) P_{\nu}(\boldsymbol{r}')}{|\boldsymbol{r} - \boldsymbol{r}'|}$$

→ The rest is matrix algebra - works with NAO basis sets! Hartree-Fock, hybrid functionals, MP2, GW, RPA, LR-TDLDA, ...

Ren, Rinke, Blum, Wieferink, Tkatchenko, Sanfilippo, Reuter, Scheffler, New J. Phys. 14, 053020 (2012)

#### Localized "Resolution of Identity" (RI) for Two-Electron Terms

$$(ij|kl) = \int d^3r d^3r' \frac{\varphi_i(\boldsymbol{r})\varphi_j(\boldsymbol{r}')\varphi_k(\boldsymbol{r})\varphi_l(\boldsymbol{r}')}{|\boldsymbol{r}-\boldsymbol{r}'|}$$

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u} (ij|
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Problem: Full RI-V delocalizes C across entire system

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Problem: Full RI-V delocalizes C across entire system

Solution: For each  $C_{ij}^{\mu}$ , restrict  $\mu$ only to atoms *I* and *J* at which *i* and *j* are centered!



Ihrig, Wieferink, Zhang, Ropo, Ren, Rinke, Scheffler, Blum, New J. Phys. **17**, 093020 (2015)

# Hybrid Functionals: Scalability, Large Systems

Levchenko, Ren, Wieferink, Rinke, Johanni, Blum, Scheffler, Comp. Phys. Commun. **192**, 60-69 (2015).

O(N) scaling implementation, localized resolution of identity.

Note: Localized RI also works for MP2, RPA, GW, etc. (but not O(N))



Zincblende GaAs

<u>Computational Scaling of Periodic GaAs</u>, HSE06 Hybrid Density-Functional Theory Time per iteration (seconds) 4096 1024 ▶ 128 atoms ■ 256 atoms 256 ← 512 atoms ▲ 1024 atoms 4096 16 64 256 1024 Number of CPUs

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# Scaling Limitations for Many-Body Theory: GW

<u> $G_0W_0$ : Widely used to obtain accurate quasiparticle energies</u> (molecules and materials)

Central: Self-Energy

$$\left(\Sigma_{\sigma}^{GW}(\mathbf{r},\mathbf{r}',\epsilon) = \frac{i}{2\pi} \int d\epsilon' \left[G_{\sigma}(\mathbf{r},\mathbf{r}',\epsilon+\epsilon')W(\mathbf{r},\mathbf{r}',\epsilon')e^{i\eta\epsilon'}\right]\right)$$

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However,  $G_0$ ,  $W_0$  not easy to converge known by different terminologies in different communities:

- "Slow-converging sums over states"
- "Slow convergence with basis set size"
- "Slow convergence of the electron-electron cusp" (Quantum Chemistry)

 $\leftrightarrow$  Large basis set requirement <u>and</u> formal  $O(N^4)$  scaling: Doubly challenging.



<u>Igor Ying Zhang</u>, Xinguo Ren, Patrick Rinke, Volker Blum, and Matthias Scheffler, New Journal of Physics **15**, 123033 (2013).



Igor Ying Zhang, Xinguo Ren, Patrick Rinke, Volker Blum, and Matthias Scheffler, New Journal of Physics **15**, 123033 (2013).



Tong Zhu



#### Can We Extrapolate Smaller (Cheaper) Basis Sets?



<u>2Z-3Z Extrapolation (*T. Zhu*</u>): <u>5Z</u> quality, but at the price of <u>3Z</u>. Works for valence and low-lying conduction levels, light-element molecules. However, not for core states, unbound states; restricted to H-Ar. Summary

$$\hat{\mathcal{H}}\Psi = E\Psi$$



P.A.M. Dirac



High-accuracy platform for predictive molecular, materials simulations Scalable to large systems, advanced electronic structure approximations

Ongoing frontier: Accurate, affordable many-body perturbation theory towards excited states for real materials



Crystalline tunable organic-inorganic hybrid materials - predictive approach to truly "new" materials, close integration with experiment





National Science Foundation







MAX-PLANCK-GESELLSCHAFT

Goal: Element-dependent, *transferable* basis sets from fast qualitative to meV-converged total energy accuracy (ground-state DFT)

from fast qualitative to meV-converged total energy accuracy (ground-state DFT)

Can't we have the computer pick good basis sets for us?

Robust iterative selection strategy: (e.g., Delley 1990)

Initial basis {*u*}<sup>(0)</sup>: Occupied free atom orbitals *u*free

from fast qualitative to meV-converged total energy accuracy (ground-state DFT)



from fast qualitative to meV-converged total energy accuracy (ground-state DFT)



from fast qualitative to meV-converged total energy accuracy (ground-state DFT)



### Iterative Selection of NAO Basis Functions

<u>"Pool" of trial basis functions:</u> 2+ ionic *u(r)* Hydrogen-like *u(r)* for z=0.1-20

**Optimization target:** 

Non-selfconsistent symmetric dimers, averaged for different d

Pick basis functions one by one, up to complete total energy convergence

Blum et al., Comp. Phys. Commun. **180**, 2175-2196 (2009)

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Systematic hierarchy of	—						
<i>automated</i> construction based on <i>dimers</i>		Au	0	С	Н		
	,	[Xe] + 6s5d4f	$[\mathrm{He}] + 2s2p$	$[\mathrm{He}] + 2s2p$	1s	minimal	
		$\mathrm{Au}^{2+}(6\mathrm{p})$	H(2p, 1.8)	$\mathrm{H}(2p,\!1.7)$	H(2s,2.1)	Tier 1	
		$\mathrm{H}(4f,\!7.4)$	H(3d, 7.6)	H(3d, 6.0)	$\operatorname{H}(2p,\!3.5)$		
"First tier (level)"		$Au^{2+}(6s)$	H(3s, 6.4)	H(2s, 4.9)			
		${\rm H}(5g,\!10)$					
		${\rm H}(6h,\!12.8)$					
<pre>     "Second tier"     "Third tier" </pre>	J	H(3d, 2.5)					
		H(5f, 14.8)	H(4f, 11.6)	$\mathrm{H}(4f,\!9.8)$	H(1s, 0.85)	Tier 2	
		H(4d, 3.9)	H(3p, 6.2)	H(3p, 5.2)	$\rm H(2p,\!3.7)$		
		$\operatorname{H}(3p,\!3.3)$	H(3d, 5.6)	H(3s, 4.3)	$\mathrm{H}(2s,\!1.2)$		
		H(1s, 0.45)	$\operatorname{H}(5g,\!17.6)$	$\mathrm{H}(5g{,}14.4)$	$\mathrm{H}(3d,\!7.0)$		
		$\mathrm{H}(5g,\!16.4)$	H(1s, 0.75)	H(3d, 6.2)			
		$\mathrm{H}(6h,\!13.6)$					
		$H(4f, 5.2)^{*}$	$\mathcal{O}^{2+}(2p)$	H(2p, 5.6)	H(4f, 11.2)	Tier 3	
		H(4d, 5.0)	H(4f, 10.8)	H(2s, 1.4)	H(3p, 4.8)		
	}	•••	•••	•••	•••		
•••							

# Accuracy: (H<sub>2</sub>O)<sub>2</sub> Hydrogen Bond Energy



# Using Numeric Atom-Centered Basis Functions: Pieces

Numerical Integration

$$h_{ij} = \int d^3 r \varphi_i(\mathbf{r}) \hat{h}_{\mathrm{KS}} \varphi_{j}(\mathbf{r})$$

• Electron density update

$$egin{aligned} n(oldsymbol{r}) &= \sum_k f_k |\psi_k(oldsymbol{r})|^2 \ v_{ ext{es}}(oldsymbol{r}) &= \int d^3 r' rac{n(oldsymbol{r}')}{|oldsymbol{r}-oldsymbol{r}'|} \end{aligned}$$

- All-electron electrostatics
- Eigenvalue solver

$$\underline{\underline{h}}\,\underline{\underline{c}}_k = \epsilon_k\,\underline{\underline{s}}\,\underline{\underline{c}}_k$$
# Using Numeric Atom-Centered Basis Functions: Pieces

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• Eigenvalue solver

• Relativity?

$$\underline{\underline{h}}\,\underline{\underline{c}}_k = \epsilon_k\,\underline{\underline{s}}\,\underline{\underline{c}}_k$$

- needed for heavy elements
- Periodic systems?

need suitable basis, electrostatics

$$(ij|kl) = \int d^3r d^3r' \frac{\varphi_i(\boldsymbol{r})\varphi_j(\boldsymbol{r}')\varphi_k(\boldsymbol{r})\varphi_l(\boldsymbol{r}')}{|\boldsymbol{r}-\boldsymbol{r}'|}$$

• Coulomb operator?

### Numeric Atom-Centered Basis Functions: Integration

$$h_{ij} = \int d^3 r \varphi_i(\boldsymbol{r}) \hat{h}_{\mathrm{KS}} \varphi_j(\boldsymbol{r})$$

• Discretize to integration grid:

$$\int d^3r f(\boldsymbol{r}) \to \sum_{\boldsymbol{r}} w(\boldsymbol{r}) f(\boldsymbol{r})$$

... but even-spaced integration grids are out: f(r) strongly peaked near all nuclei!

# Numeric Atom-Centered Basis Functions: Integration

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- Discretize to integration grid:  $\int d^3r f(\mathbf{r}) \rightarrow \sum_{\mathbf{r}} w(\mathbf{r}) f(\mathbf{r})$ ... but even-spaced integration grids are *out*:  $f(\mathbf{r})$  strongly peaked near all nuclei!
- Overlapping atom-centered integration grids:
  - Radial shells (e.g., H, light: 24; Au, tight: 147)
  - Specific angular point distribution ("Lebedev") exact up to given integration order *l* (50, 110, 194, 302, .... points per shell)



Pioneered by Becke JCP 88, 2547 (1988), Delley, JCP 92, 508 (1990), MANY others!

#### All-Electron Integrals: Rather Benign for NAOs

$$\int d^3r \phi_{1s}(\boldsymbol{r}) \hat{H} \phi_{1s}(\boldsymbol{r}) = \int dr \left[f(r)\right] \times \text{angular integral.}$$



#### All-Electron Integrals: Rather Benign for NAOs



# Overlapping Atom-Centered Grids: "Partitioning of Unity"

Becke, 1988

$$h_{ij} = \int d^3 r \varphi_i(\boldsymbol{r}) \hat{h}_{\rm KS} \varphi_j(\boldsymbol{r})$$

#### • <u>Rewrite to atom-centered integrands:</u>

$$\begin{split} & \left( \int d^3 r f(\boldsymbol{r}) = \sum_{\text{atoms}} \int d^3 r p_{\text{atom}}(\boldsymbol{r}) f(\boldsymbol{r}) \right) \\ & \underline{\text{exact:}} \qquad \sum_{\text{atoms}} p_{\text{atom}}(\boldsymbol{r}) = 1 \\ & \text{through} \quad p_{\text{atom}}(\boldsymbol{r}) = \frac{g_{\text{atom}}(\boldsymbol{r})}{\sum_{\text{atom'}} g_{\text{atom'}}(\boldsymbol{r})} \end{split}$$

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through 
$$p_{\text{atom}}(\mathbf{r}) = \frac{g_{\text{atom}}(\mathbf{r})}{\sum_{\text{atom'}} g_{\text{atom'}}(\mathbf{r})}$$

• e.g.: 
$$g_{
m atom} = rac{
ho_{
m atom}(r)}{r^2}$$
 (Delley 1990)

many alternatives: Becke 1988, Stratmann 1996, Koepernik 1999, ...



#### Integration in Practice: Large Systems, Small Errors!



Fully extended Polyalanine peptide molecule Ala20, DFT-PBE (203 atoms)



# Hartree Potential (Electrostatics): Overlapping Multipoles

$$v_{
m es}(oldsymbol{r}) = \int d^3r' rac{n(oldsymbol{r}')}{|oldsymbol{r}-oldsymbol{r}'|}$$

• Partitioning of Unity: (same trick as used for integrals)

$$n(m{r}) = \sum_{\mathrm{atoms}} p_{\mathrm{atom}}(m{r}) n(m{r})$$



atoms

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• Classical electrostatics:

$$v_{
m es}(\boldsymbol{r}) = \sum_{
m atoms} \sum_{lm}^{l_{
m max}} v_{
m atom, lm}(|\boldsymbol{r} - \boldsymbol{R}_{
m atom}|) Y_{lm}(\Omega_{
m atom})$$

#### **Electrostatics:** Multipole expansion

$$egin{aligned} v_{ ext{es}}(m{r}) &= \sum_{ ext{atoms}} \sum_{m{lm}}^{m{l_{ ext{max}}}} v_{ ext{atom},m{lm}} (|m{r}-m{R}_{ ext{atom}}|) Y_{m{lm}}(\Omega_{ ext{atom}}) \ . \end{aligned}$$

Polyalanine Ala<sub>20</sub>, DFT-PBE (203 atoms)  $\alpha$ -helical vs. extended: Total energy convergence with  $l_{\max}$ E(l)-E(l=12) [eV]0.4  $E_{tot}(\alpha helix)$  $E_{tot}$ (extended) Energy difference 0.2 Q 0 -0.2 (b)(a) (c)-0.4 23456 4 5 6 23456 2 3 8 12 8 12 8 12 Maximum angular momentum *l* in multipole expansion extended Ω