



MAX-PLANCK-GESELLSCHAFT



When $(2 + 2) \neq 4$

or

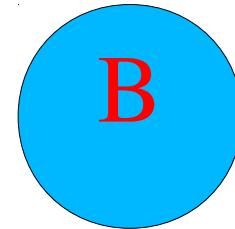
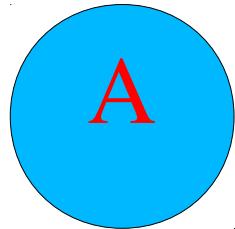
**Van der Waals Interactions in Complex
(and Simple) Materials**

Alexandre Tkatchenko

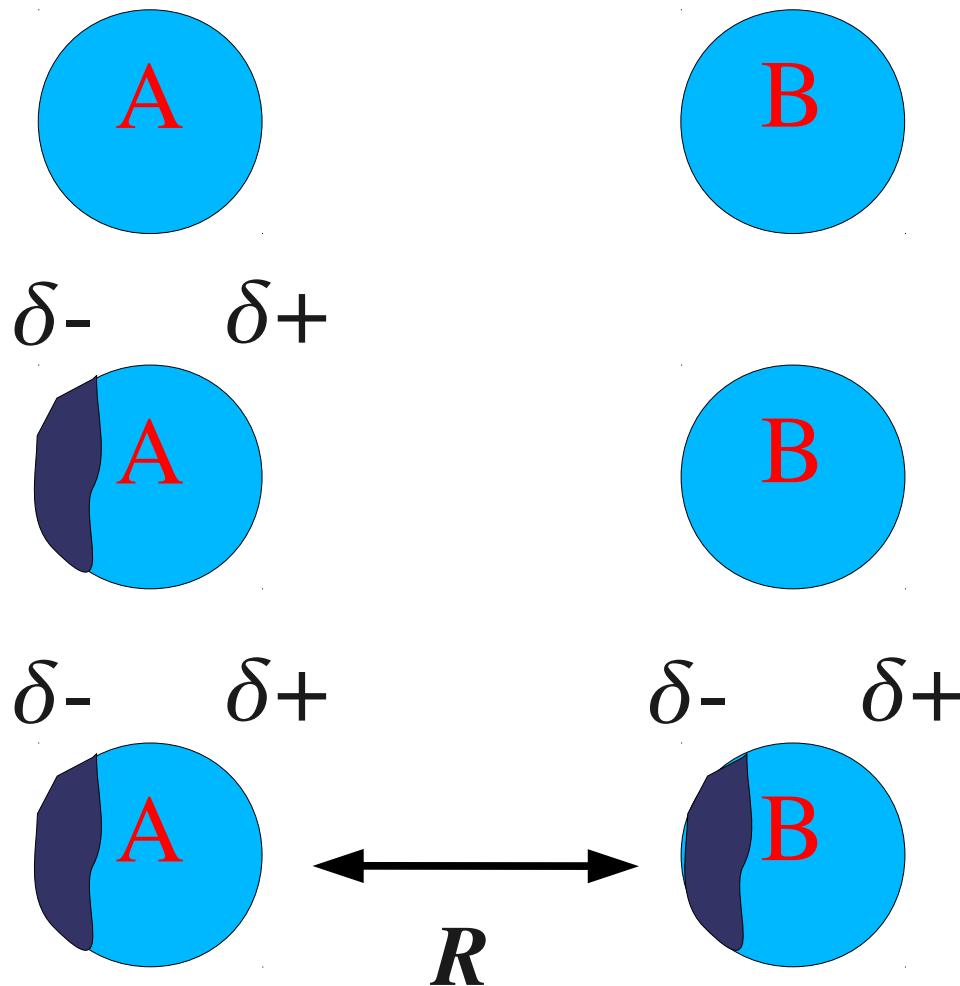
*Fritz-Haber-Institut der Max-Planck-Gesellschaft,
Berlin, Germany*

FPCSS2013, Norderney, Jul 25, 2013

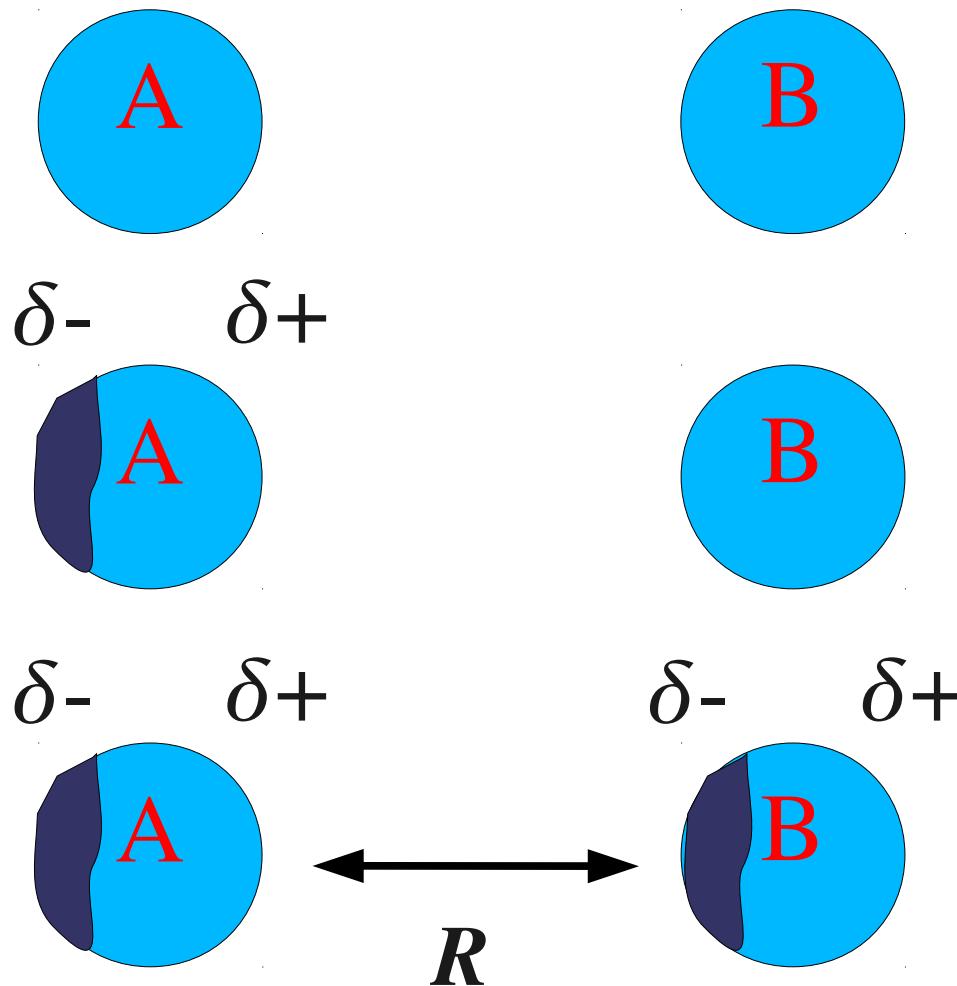
Textbook picture of vdW interactions



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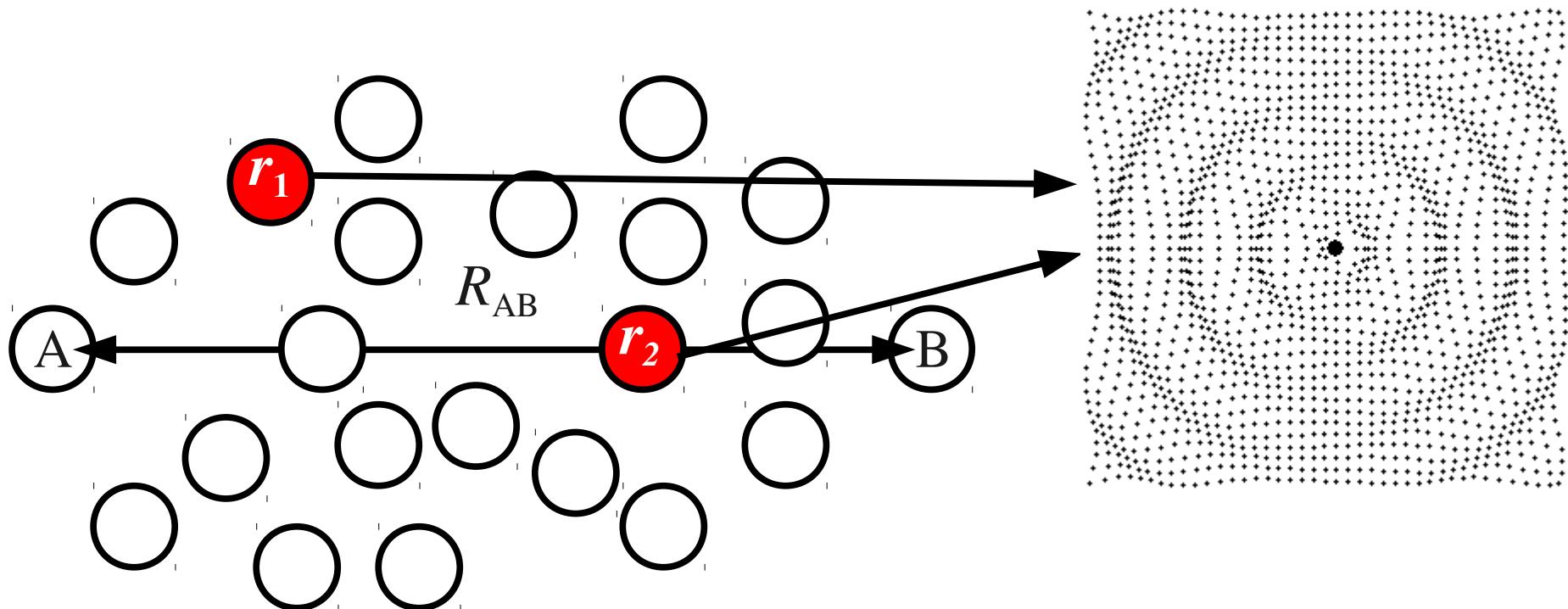
Textbook picture of vdW interactions



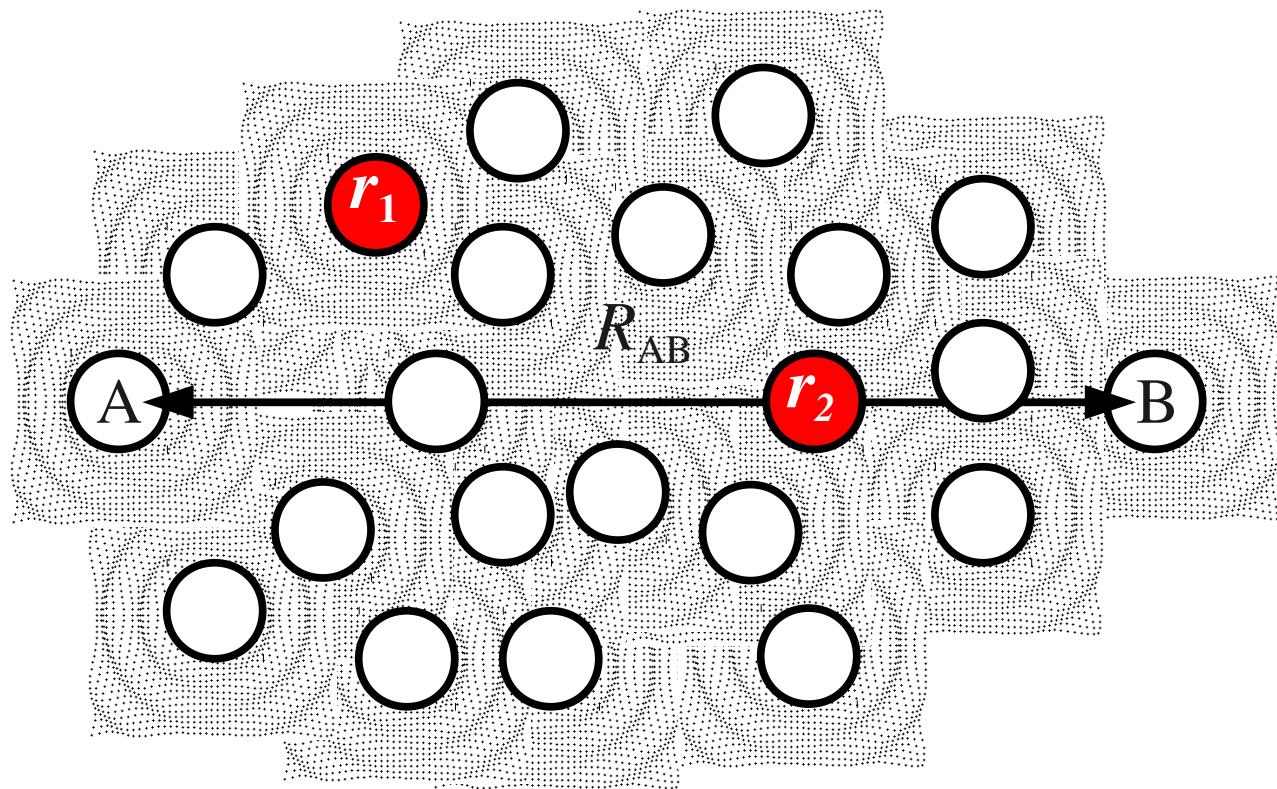
$$E_{\text{vdW}}^{(2)} = -\frac{C_6^{\text{AB}}}{R_{\text{AB}}^6}$$

$$\rightarrow C_6^{\text{AB}} = \frac{3}{\pi} \int \alpha_A(i\omega) \alpha_B(i\omega) d\omega$$

Beyond textbook model of vdW interactions: Electrodynamic response effects



Beyond textbook model of vdW interactions: Electrodynamic response effects



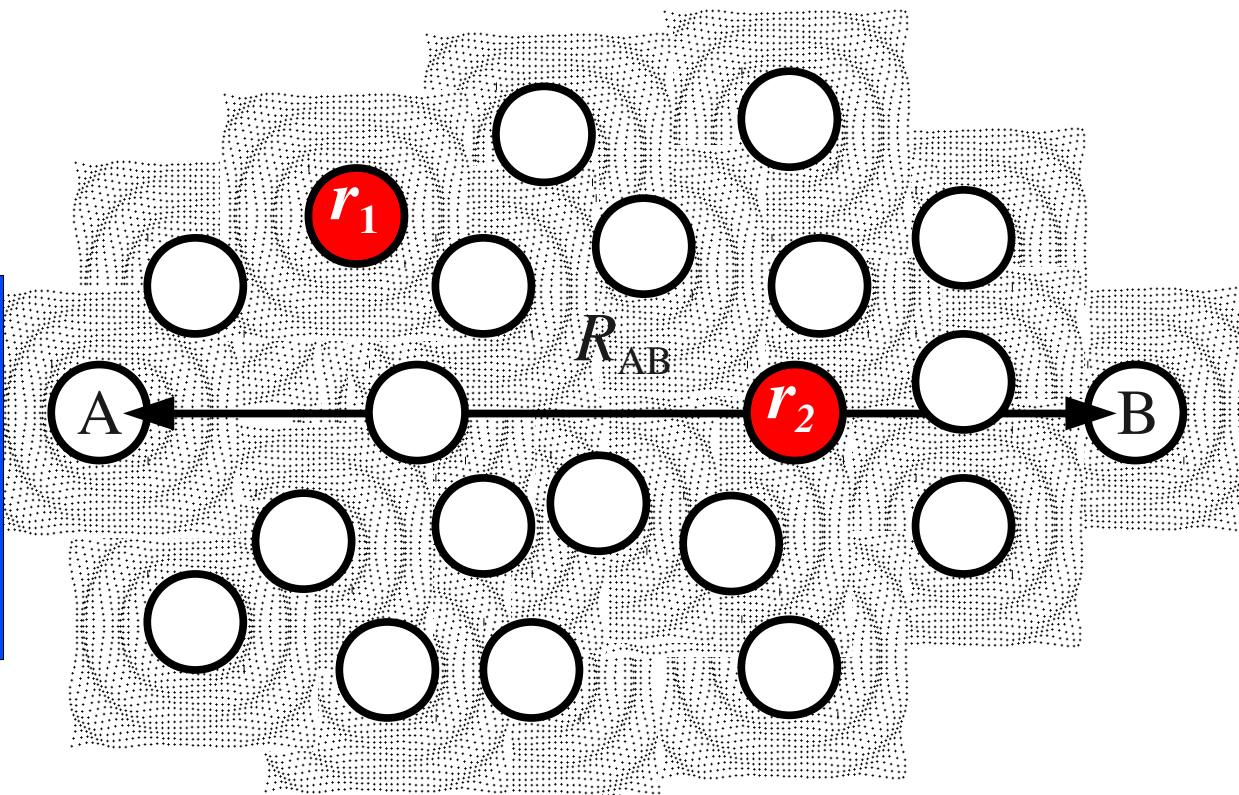
Beyond textbook model of vdW interactions: Electrodynamic response effects

1

Accurate
Microscopic
Modeling of
Coulomb
Response

2

Full (All-Order)
Many-Body
van der Waals
Energy



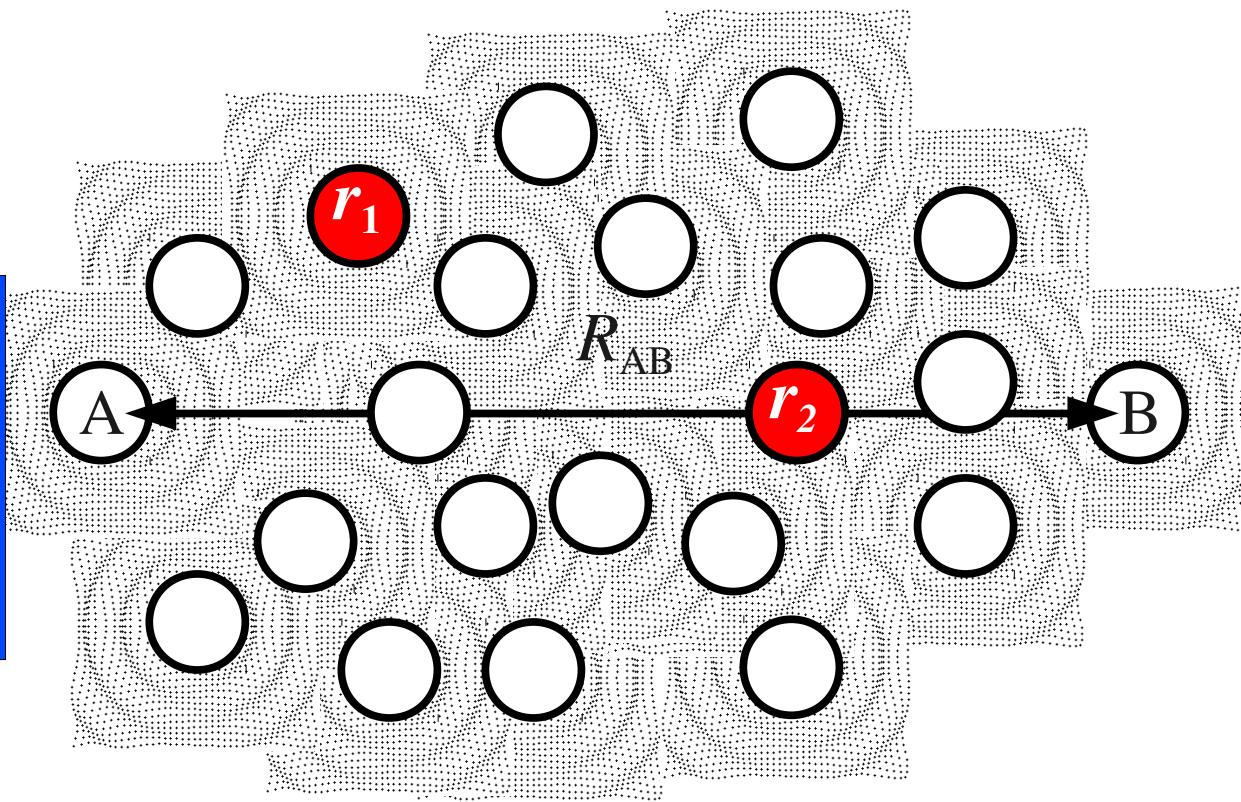
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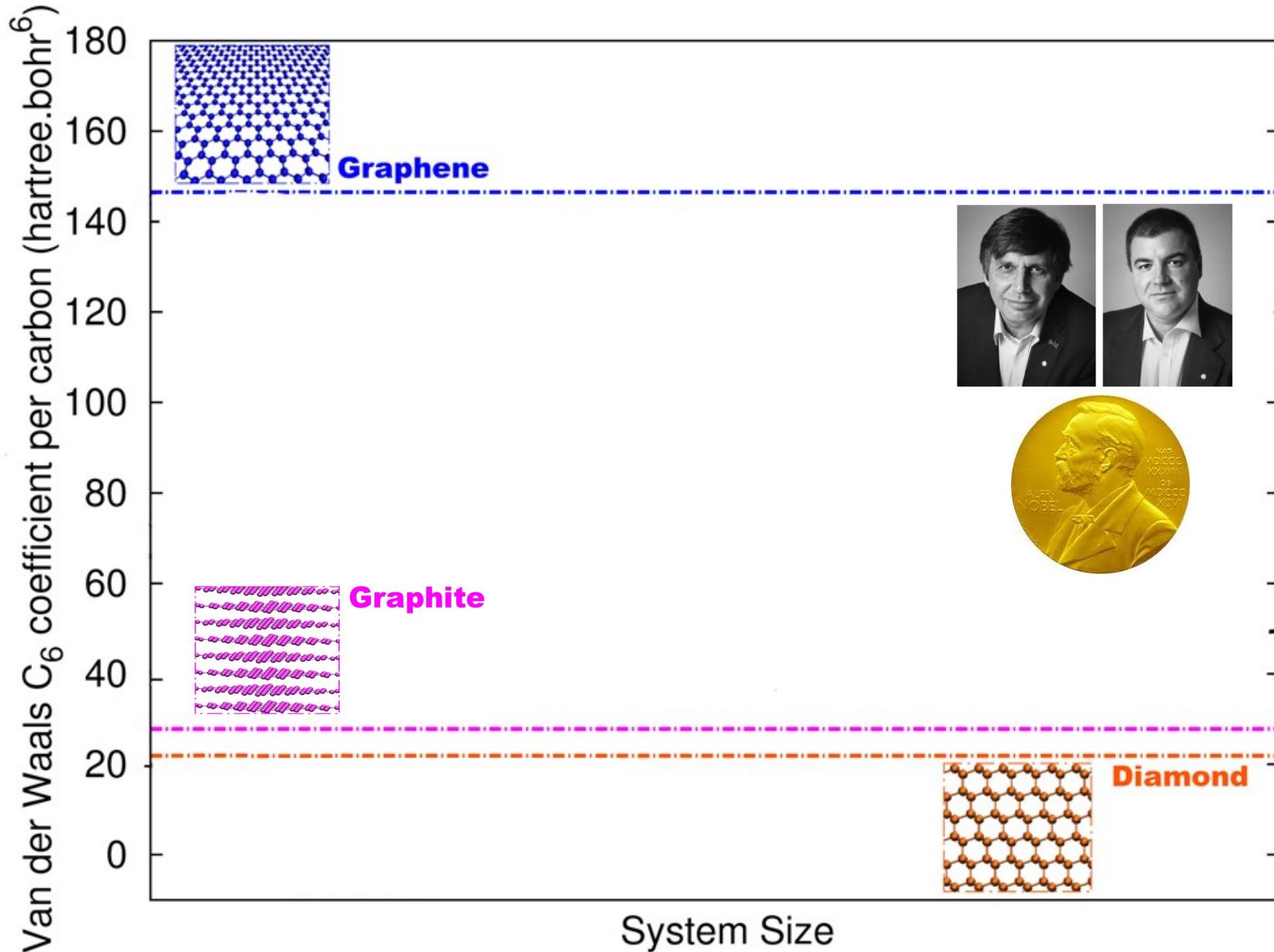
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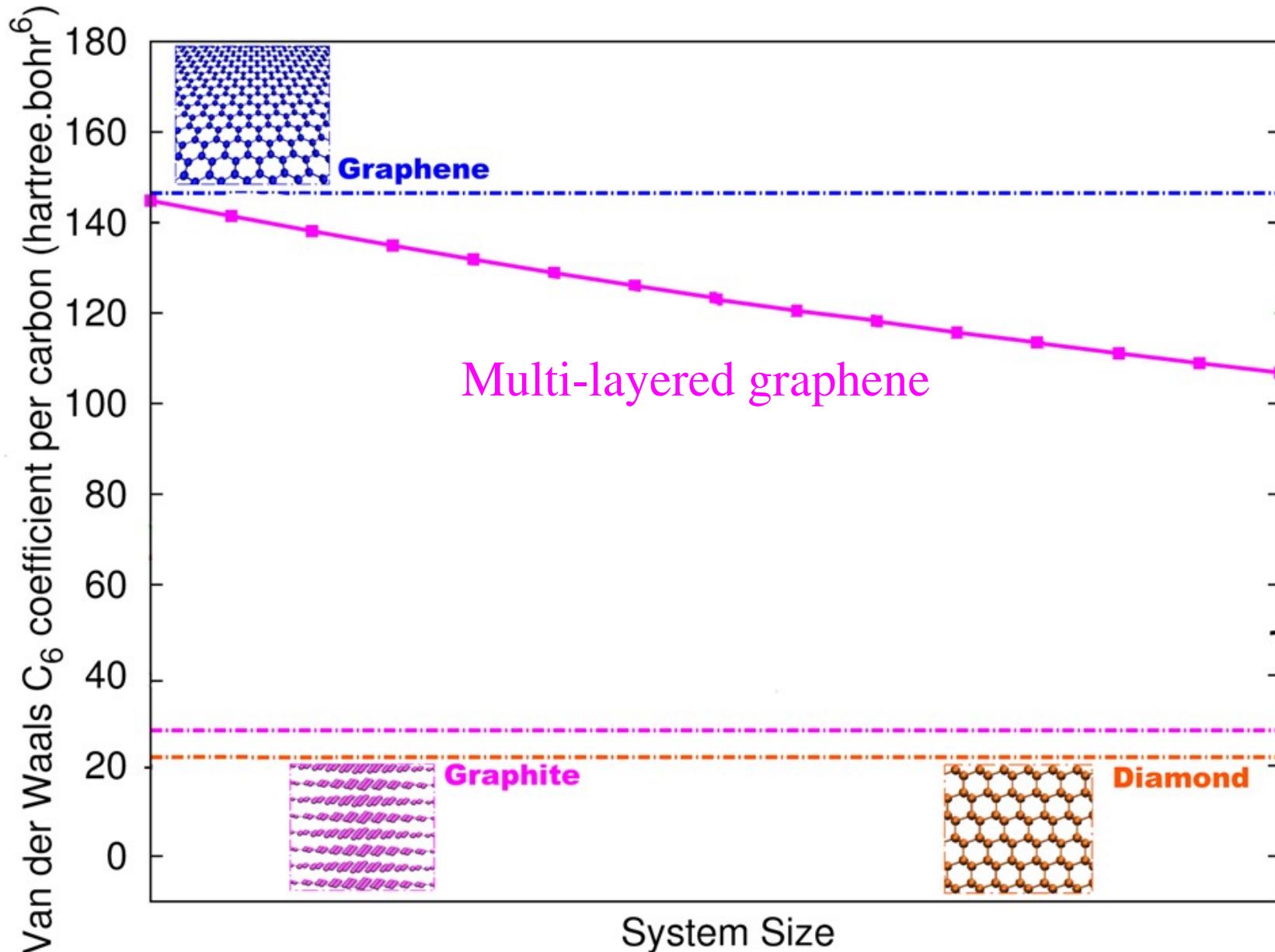
$$E_c = - \int_0^\infty \frac{d\omega}{2\pi} \int_0^1 d\lambda \text{Tr} \left((\chi_\lambda(\mathbf{r}_1, \mathbf{r}_2; i\omega) - \chi_0(\mathbf{r}_1, \mathbf{r}_2; i\omega)) \frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|} \right)$$

We know how to solve the problem, albeit not very efficiently

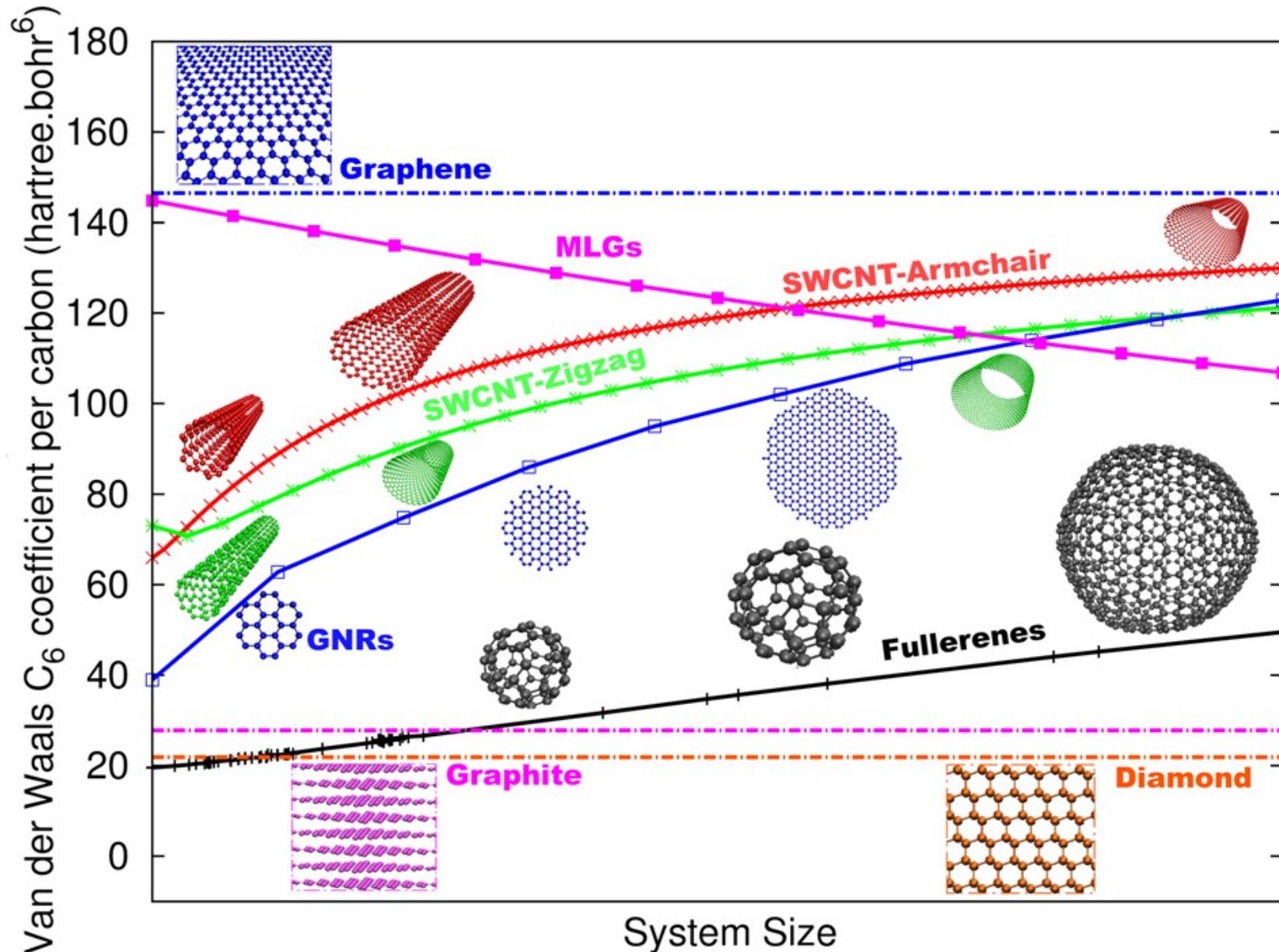
Electrodynamic treatment of vdW interactions: beyond ‘hybridized atoms’



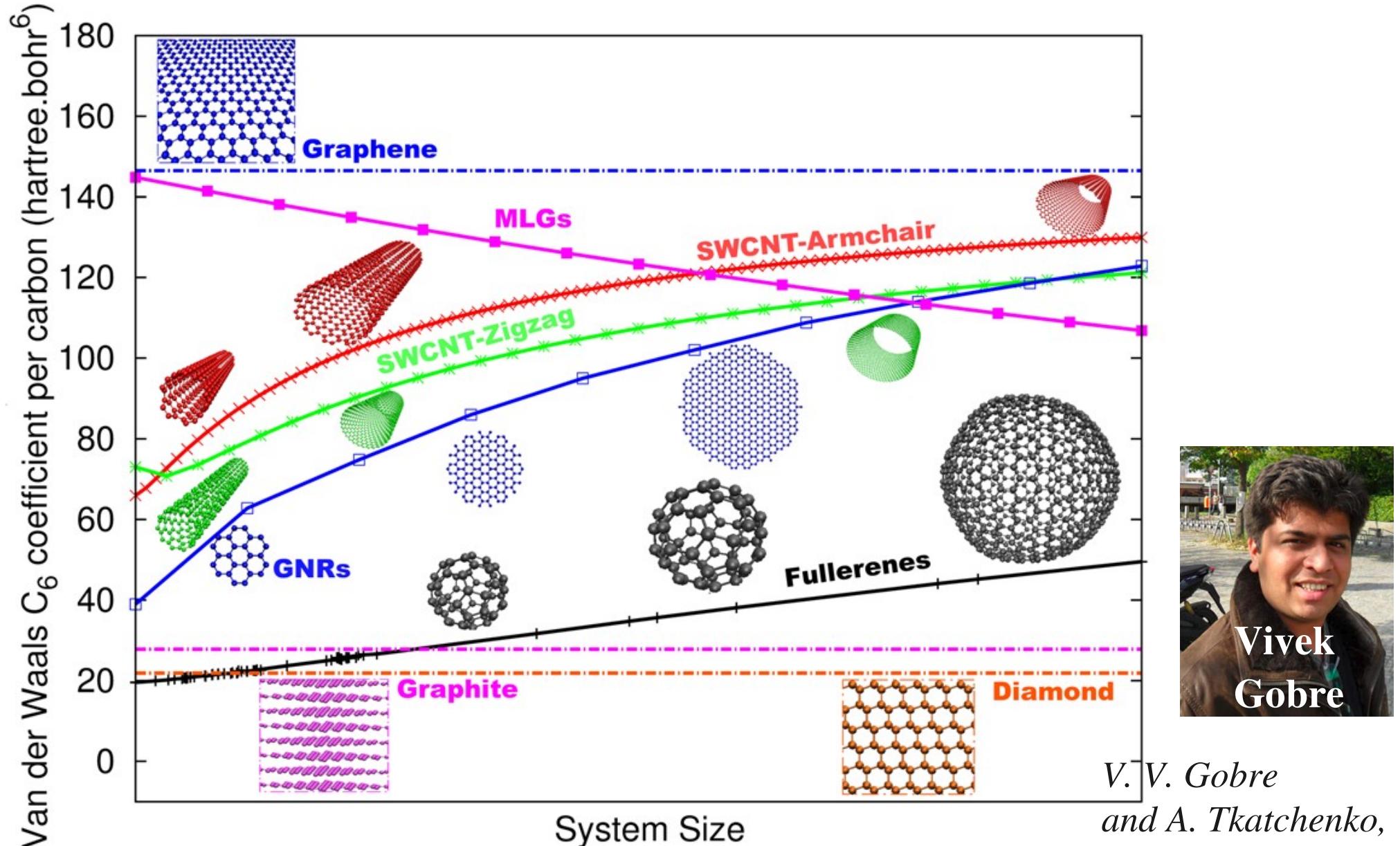
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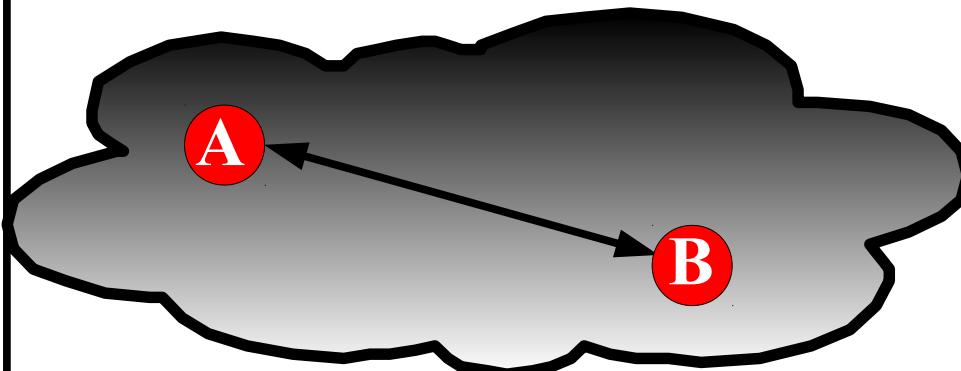
V. V. Gobre
and A. Tkatchenko,
Nature Comm. (2013)

Towards Efficient Many-Body Treatment of vdW Interactions

The conventional approach

(Grimme, Johnson-Becke/Corminboeuf,
Tkatchenko-Scheffler, Langreth-Lundqvist
vdW-DF, Vydrov-van Voorhis, ...)

Effective screening and
two-body energy



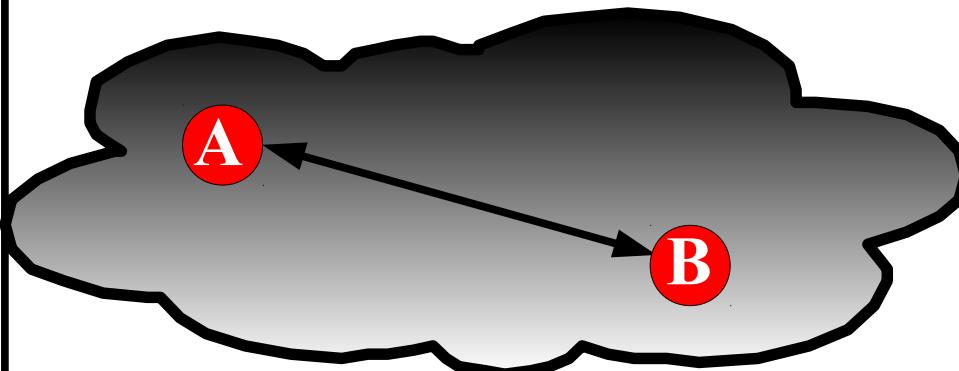
Valid for small molecules *or*
homogeneous dielectrics

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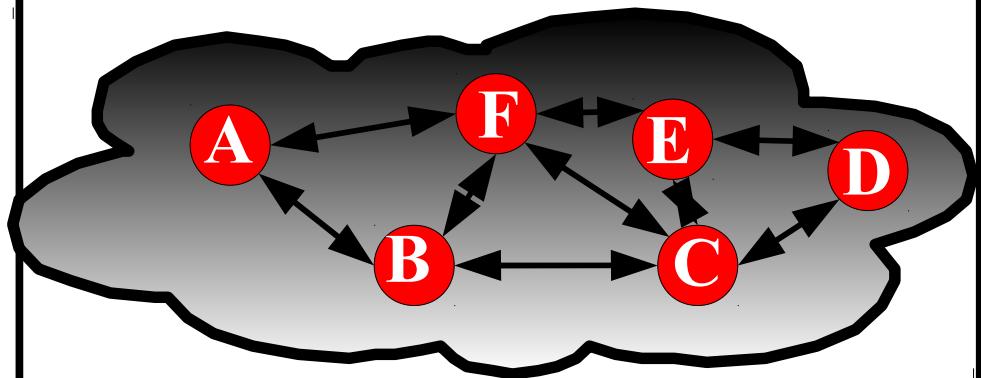
Effective screening and
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The new state-of-the-art:

Full many-body response
and energy for a system of
quantum oscillators (**DFT+MBD**)



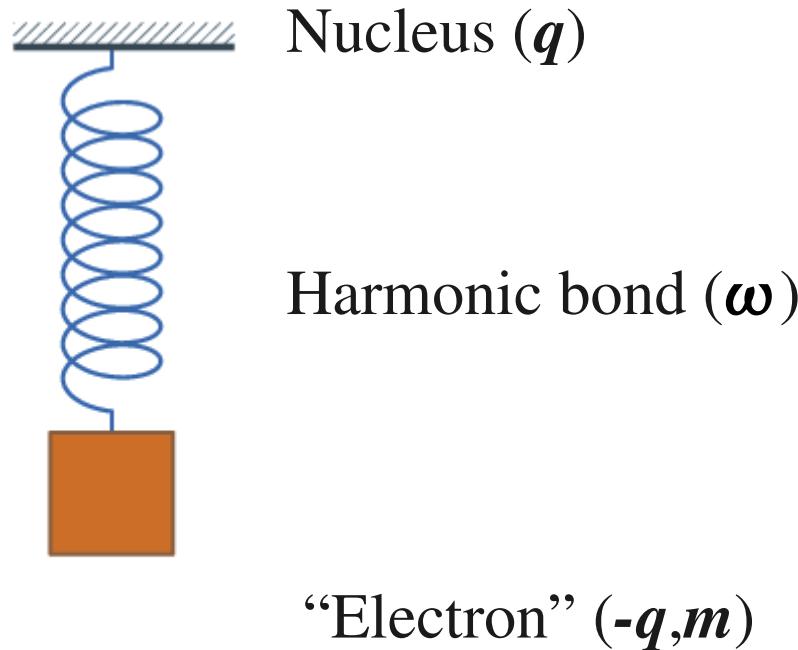
Valid for **small and large
molecules, insulators, metals,
interfaces, ...**

A. Tkatchenko and M. Scheffler, *Phys. Rev. Lett.* (2009).

A. Tkatchenko, R. A. DiStasio Jr., R. Car, M. Scheffler, *Phys. Rev. Lett.* (2012).

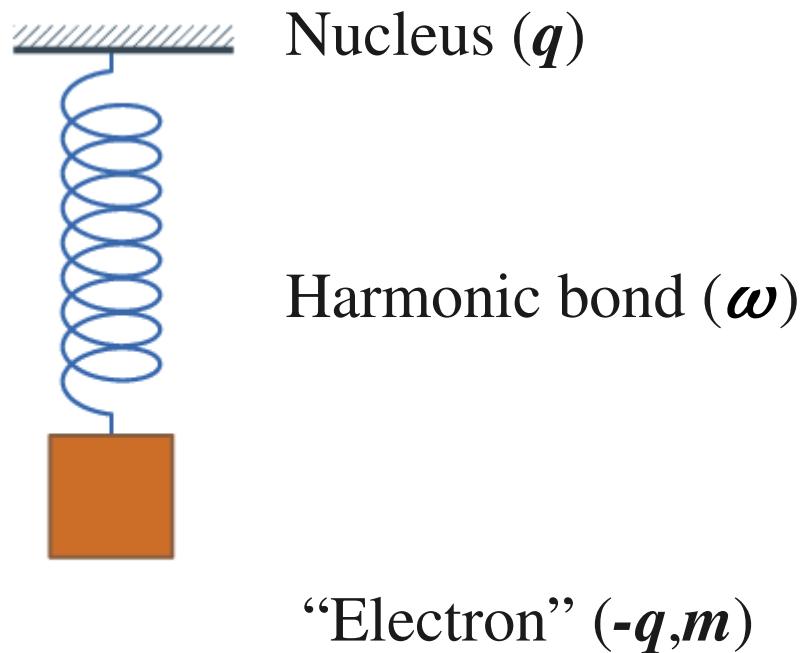
R. A. DiStasio Jr., O. A. von Lilienfeld, A. Tkatchenko, *Proc. Natl. Acad. Sci.* (2012).

The Model: Quantum Harmonic Oscillator (QHO)



Model proposed by *W. L. Bade* (1957); and used by *B. J. Berne*; *A. Donchev*; *M. W. Cole*; *G. Martyna*; *K. Jordan*; and others.

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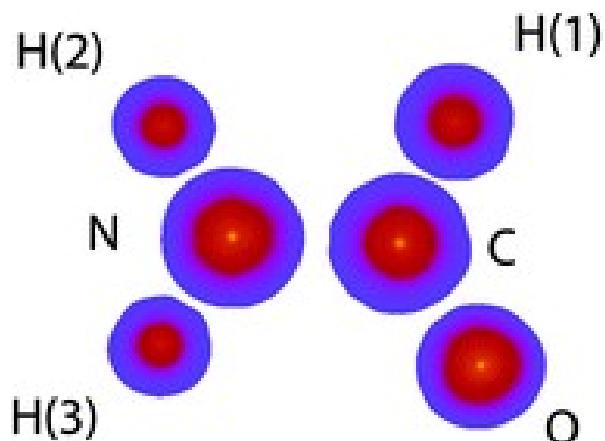
In the dipole approximation:

(α, ω) fully characterize the QHO

Model proposed by *W. L. Bade* (1957); and used by *B. J. Berne*; *A. Donchev*; *M. W. Cole*; *G. Martyna*; *K. Jordan*; and others.

First-Principles Model: QHOs in Molecules and Solids

$$\alpha_A^0 = \alpha_A^0[n(\mathbf{r})]; \quad \omega_A^0 = \omega_A^0[n(\mathbf{r})]$$



$$\alpha_A(i\omega) = \frac{\alpha_A^0}{1 + (\omega/\omega_A^0)^2}$$

α^0 and ω^0 include short-range hybridization

$$C_6 = C_6[n(\mathbf{r})], \quad R_{\text{vdW}} = R_{\text{vdW}}[n(\mathbf{r})]$$

A. Tkatchenko and M. Scheffler, *Phys. Rev. Lett.* (2009)

The Method: DFT+MBD

TS-vdW method

$$\alpha_A^0 = \alpha_A^0[n(\mathbf{r})]; \quad \omega_A^0 = \omega_A^0[n(\mathbf{r})]$$



Self-consistent electrodynamic response (Dyson)

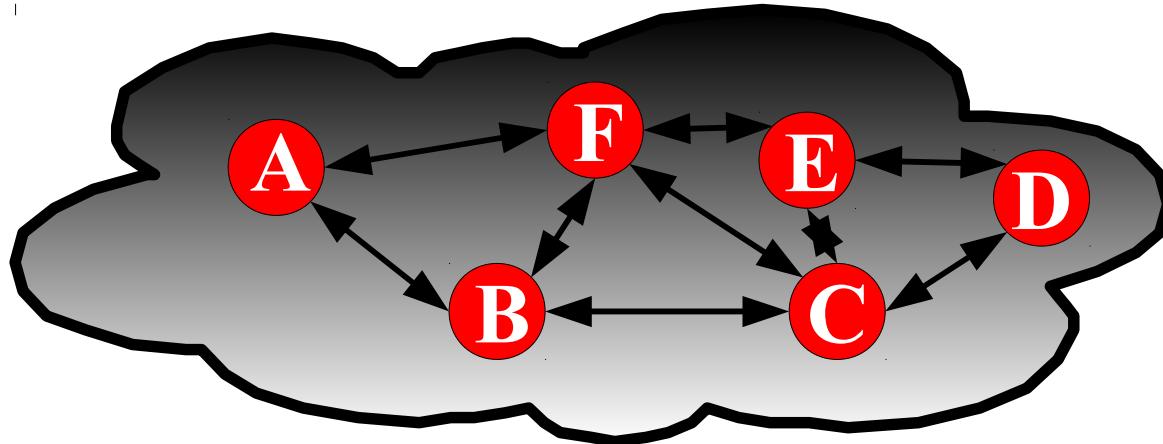
$$\alpha_p(i\omega) = \alpha_p^{\text{TS}}(i\omega) + \alpha_p^{\text{TS}}(i\omega) \sum_{q \neq p}^N \mathcal{T}_{pq} \alpha_q(i\omega)$$



Many-body vdW energy for a system of
coupled oscillators (RPA)

$$H = -\frac{1}{2} \sum_{i=1}^N \nabla_{\chi_i}^2 + \frac{1}{2} \sum_{i=1}^N \omega_i^2 \chi_i^2 + \sum_{i>j=1}^N \omega_i \omega_j \sqrt{\alpha_i \alpha_j} \chi_i T_{ij} \chi_j.$$

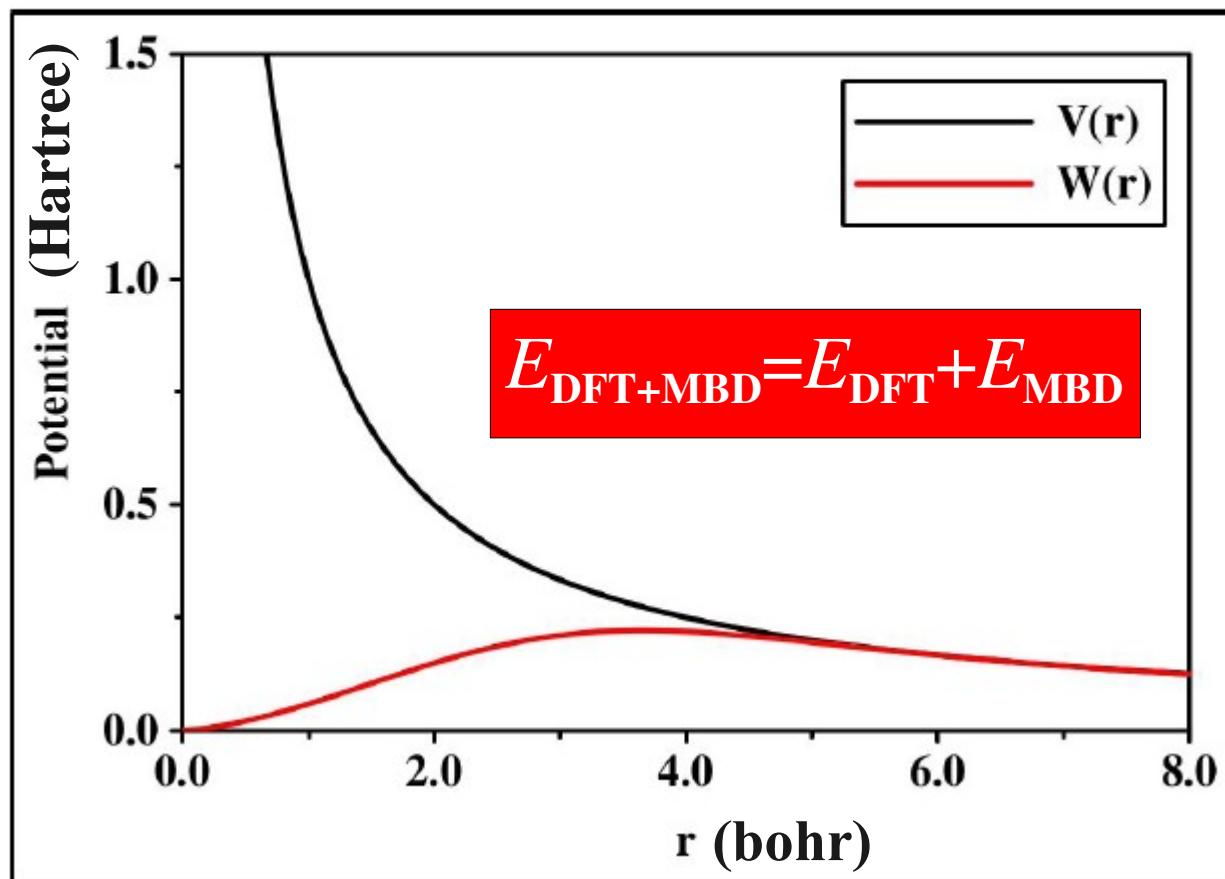
Salient Features of the MBD Method



- Seamless treatment of short-range (quantum) and long-range (classical) electrodynamic response
- Full correlation energy of coupled QHOs is equivalent to the random-phase approximation or ring-CCD (*JCP* 138, 074106 (2013))
- Computes many-body vdW energy to *infinite order*
- Negligible computational cost compared to DFT (**MBD** calculations can be easily done for > 10,000 atoms)

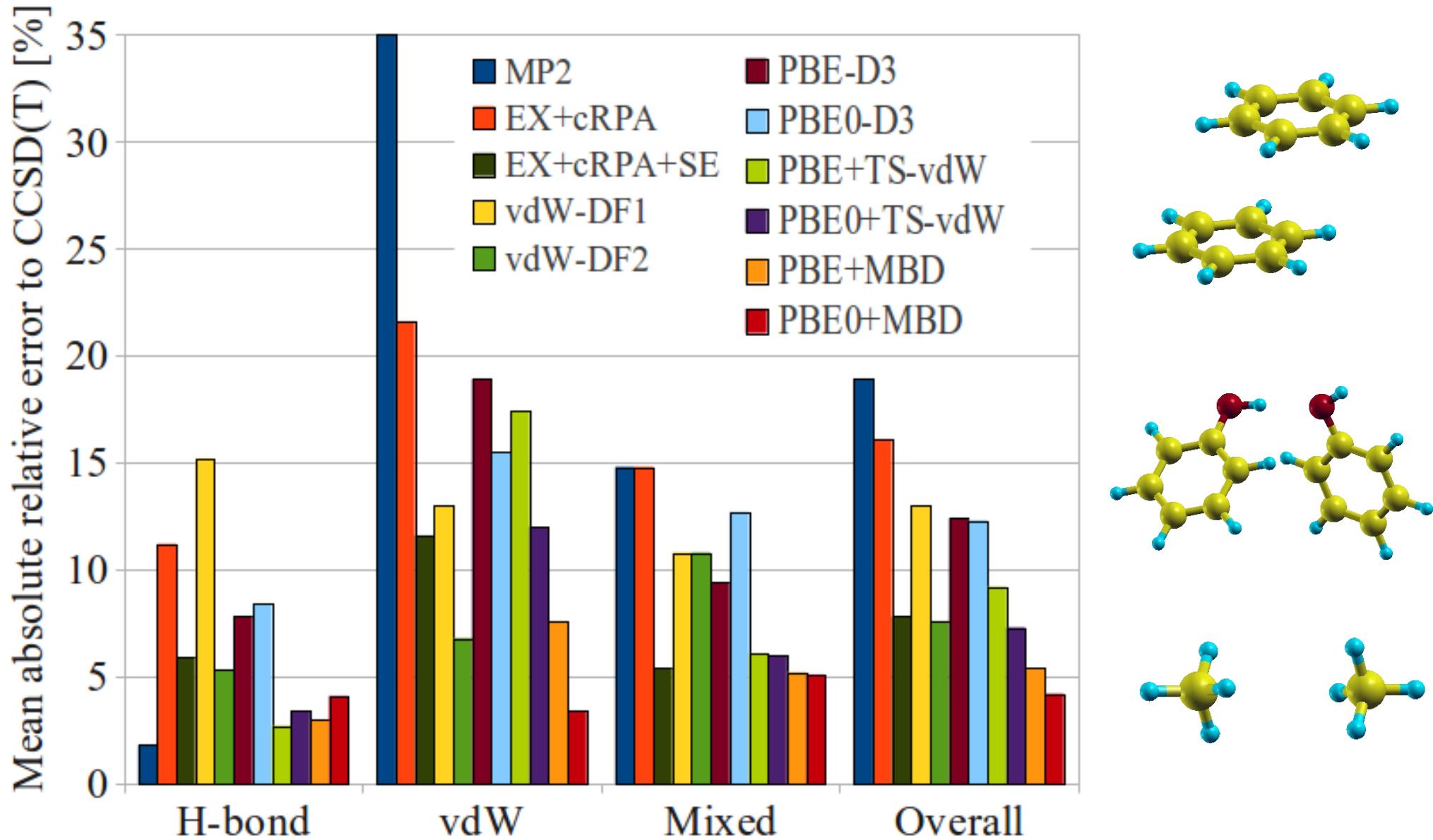
Coupling DFT and MBD by range separation of the Coulomb potential

$$W(r_{pq}) = \left(1 - \exp\left(-\left(r_{pq}/R_{pq}^{\text{vdW}}\right)^\beta\right)\right) / r_{pq}$$



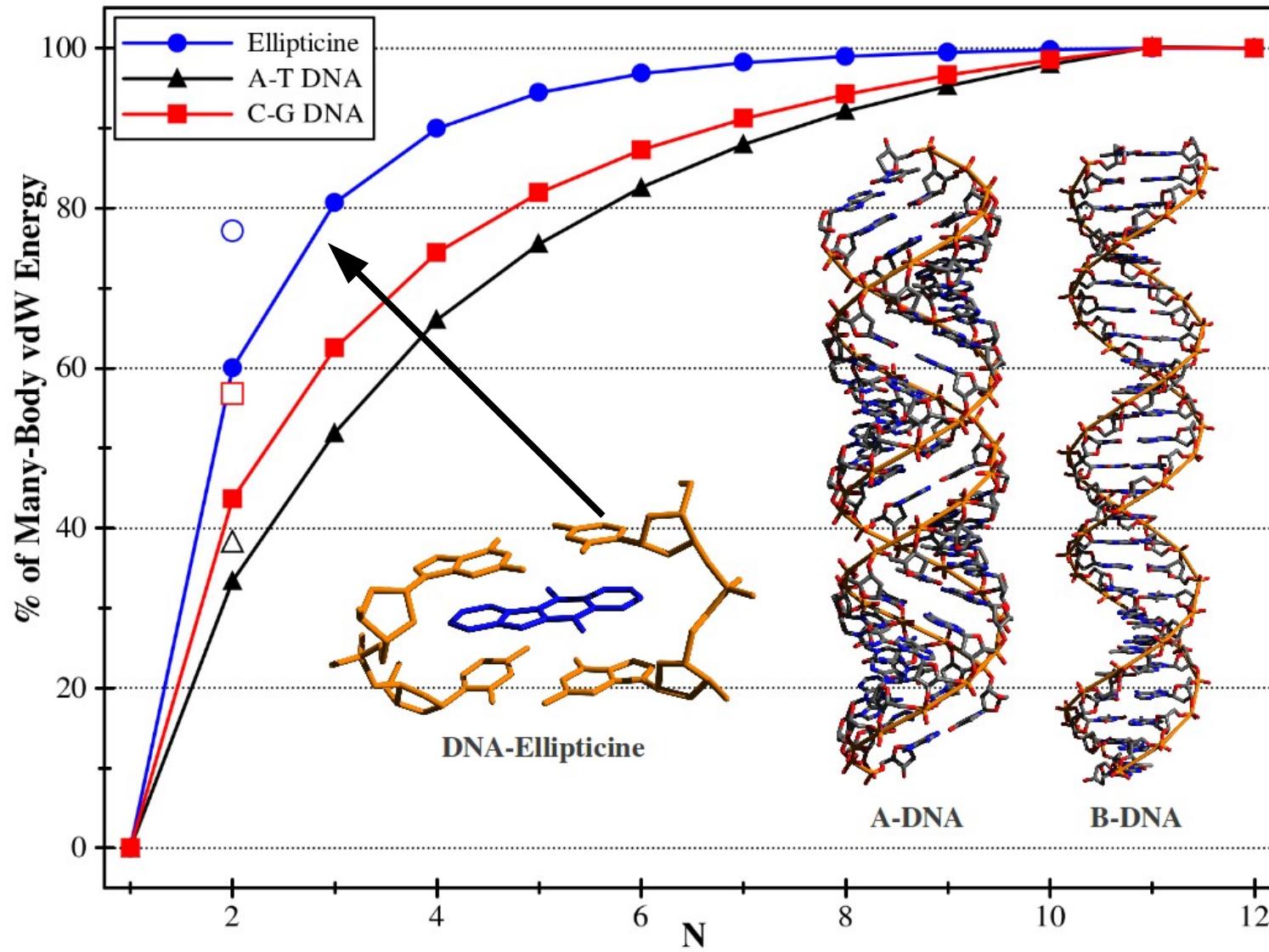
See work on range separation by *A. Savin, H. Stoll, G. Scuseria, K. Hirao, ...*

Performance of DFT+MBD for gas-phase intermolecular interactions



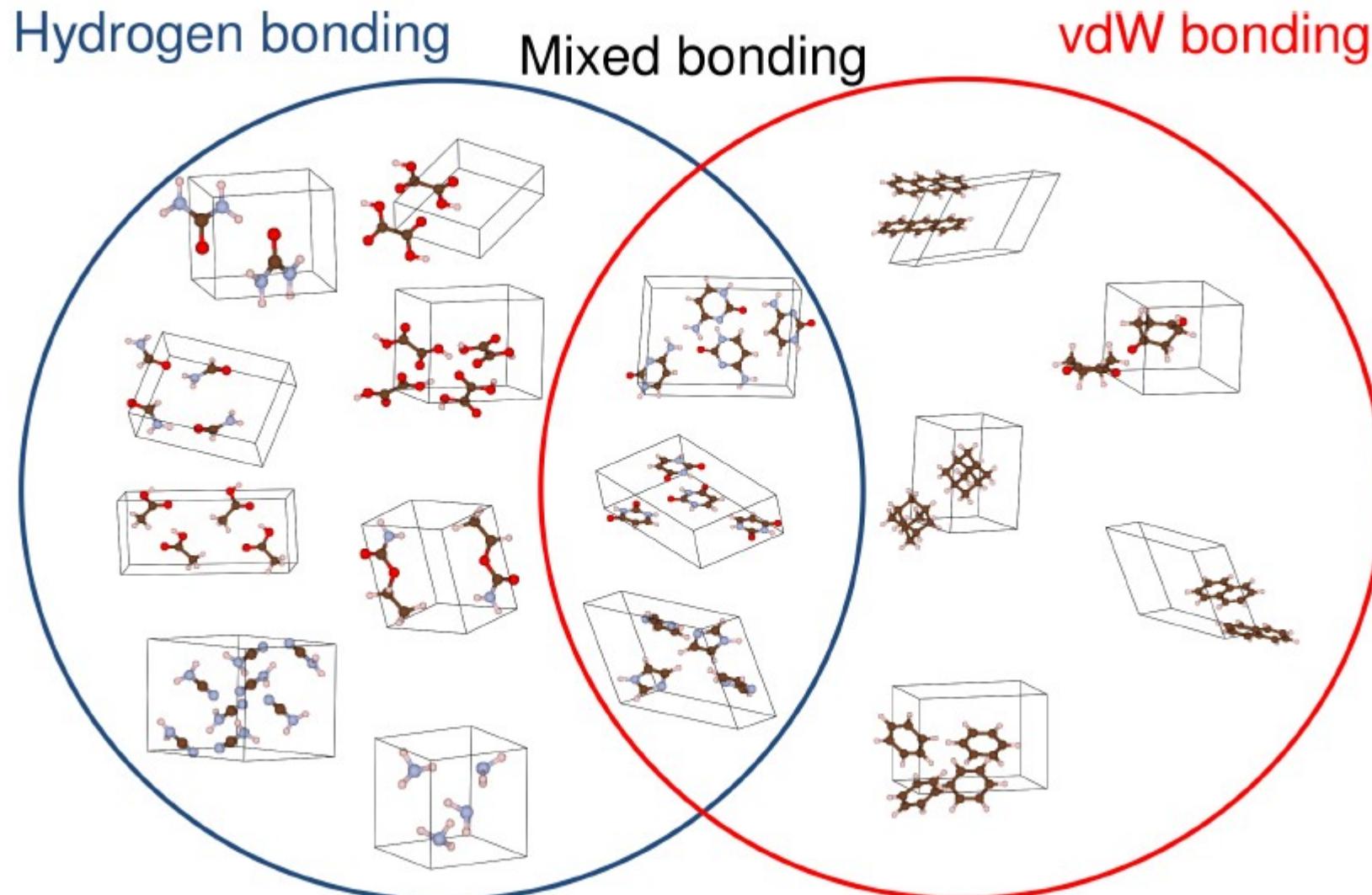
S22 CCSD(T): Jurecka, Sponer, Cerny, Hobza, *PCCP* (2006); Sherrill *et al.*, *JCP* (2010).

Large many-body vdW effects in complex molecular geometries



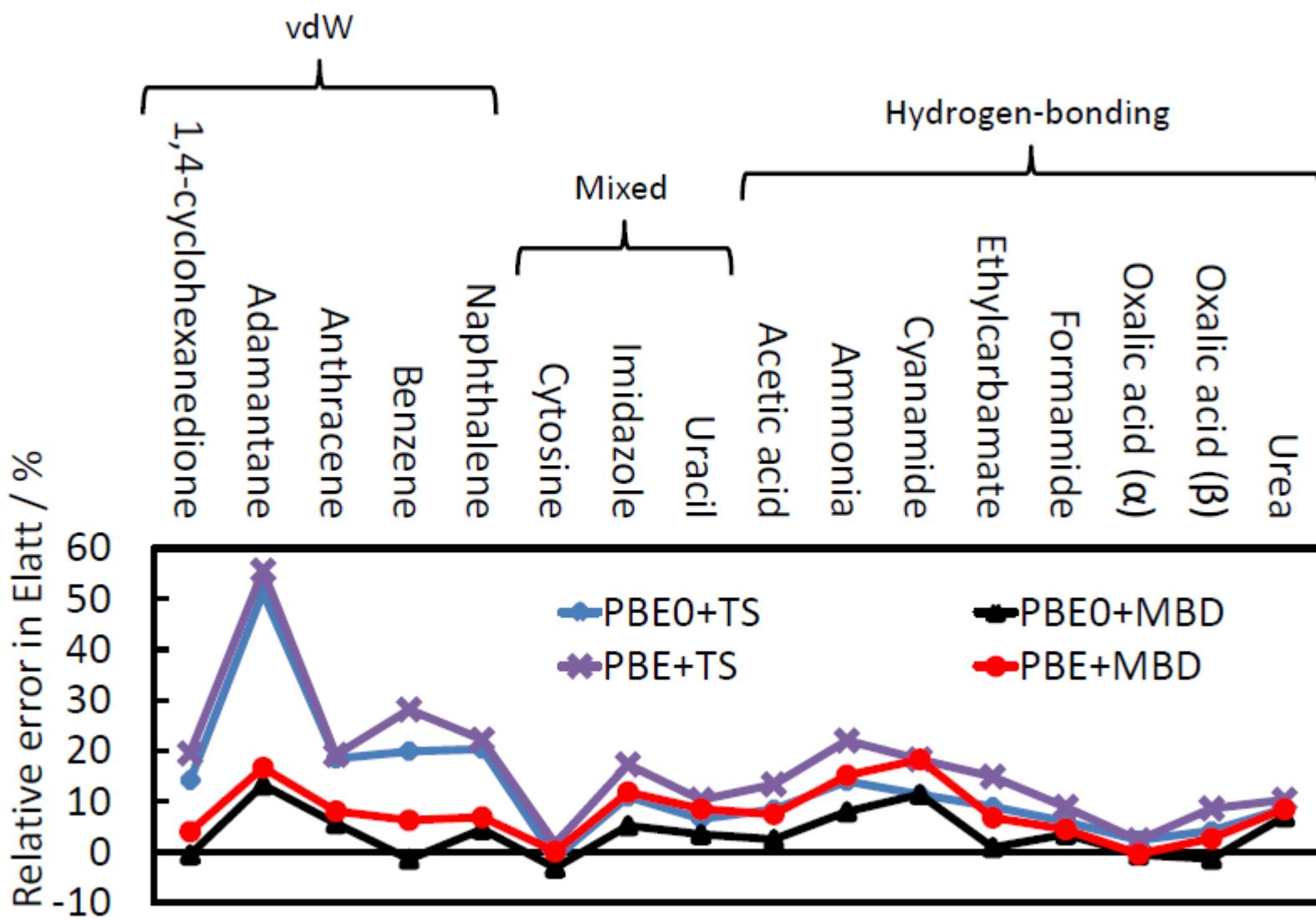
R. A. DiStasio Jr., O. A. von Lilienfeld, and A. Tkatchenko, *PNAS* (2012)

“Chemically Accurate” Predictions for Molecular Materials



A. M. Reilly and A. Tkatchenko, *J. Phys. Chem. Lett.* **4**, 1028 (2013).

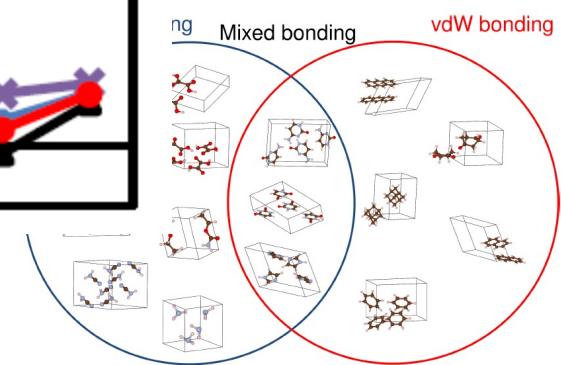
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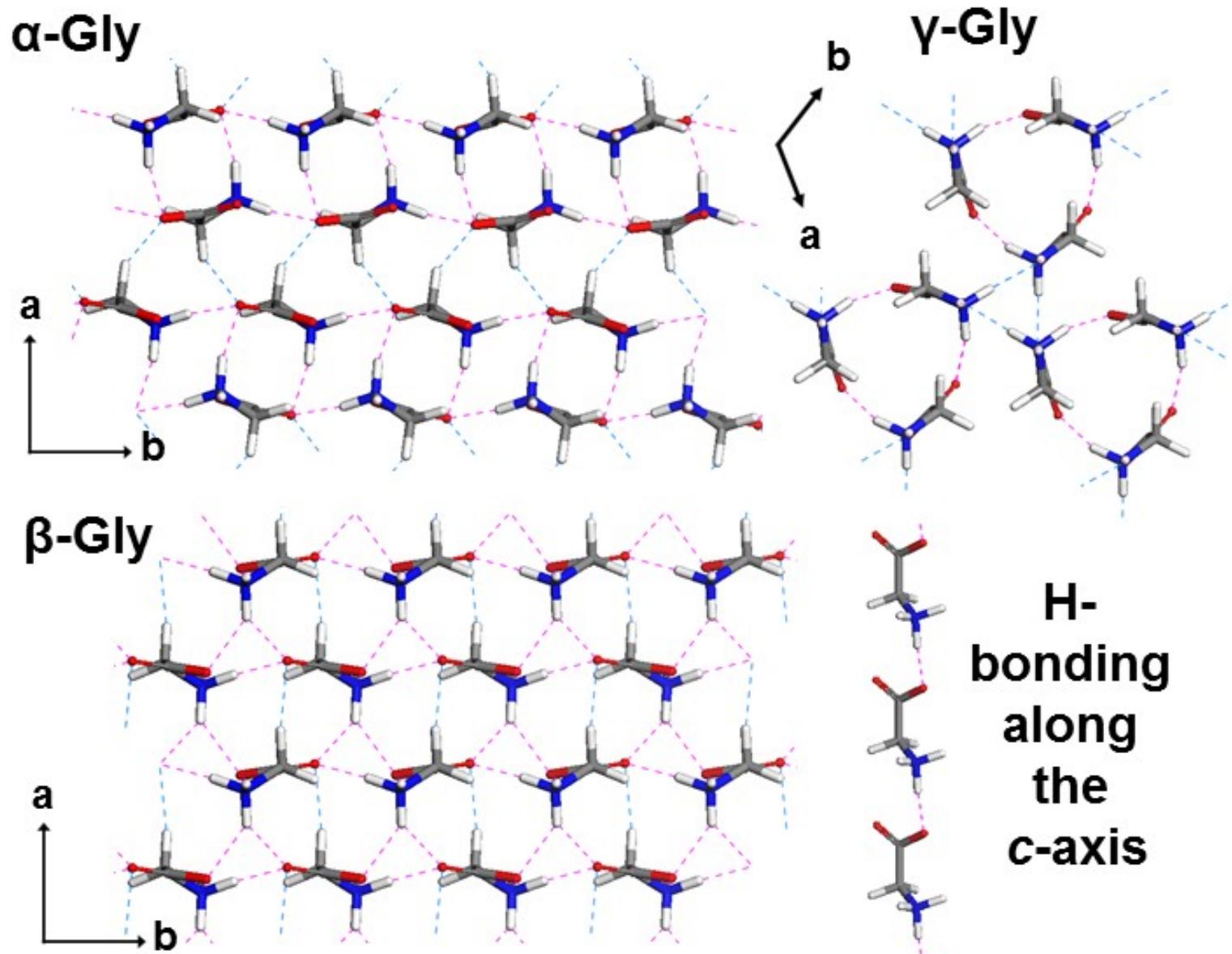
Anthony
Reilly

PBE0+MBD overall accuracy: **0.8 kcal/mol; 4.6%**

A. M. Reilly and A. Tkatchenko, *J. Phys. Chem. Lett.* 4, 1028 (2013).

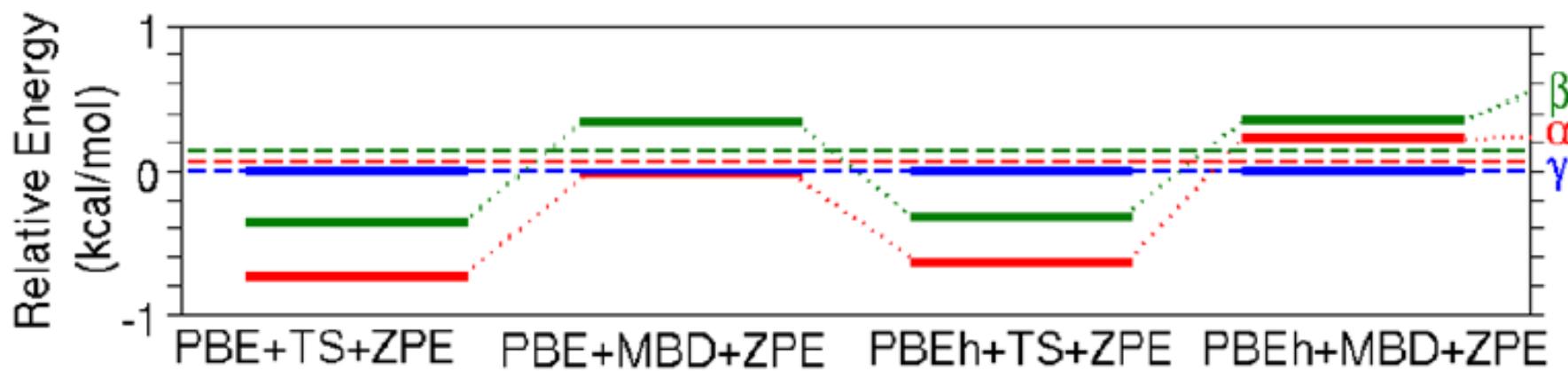


DFT+MBD correctly discriminates between crystal polymorphs: Example of glycine



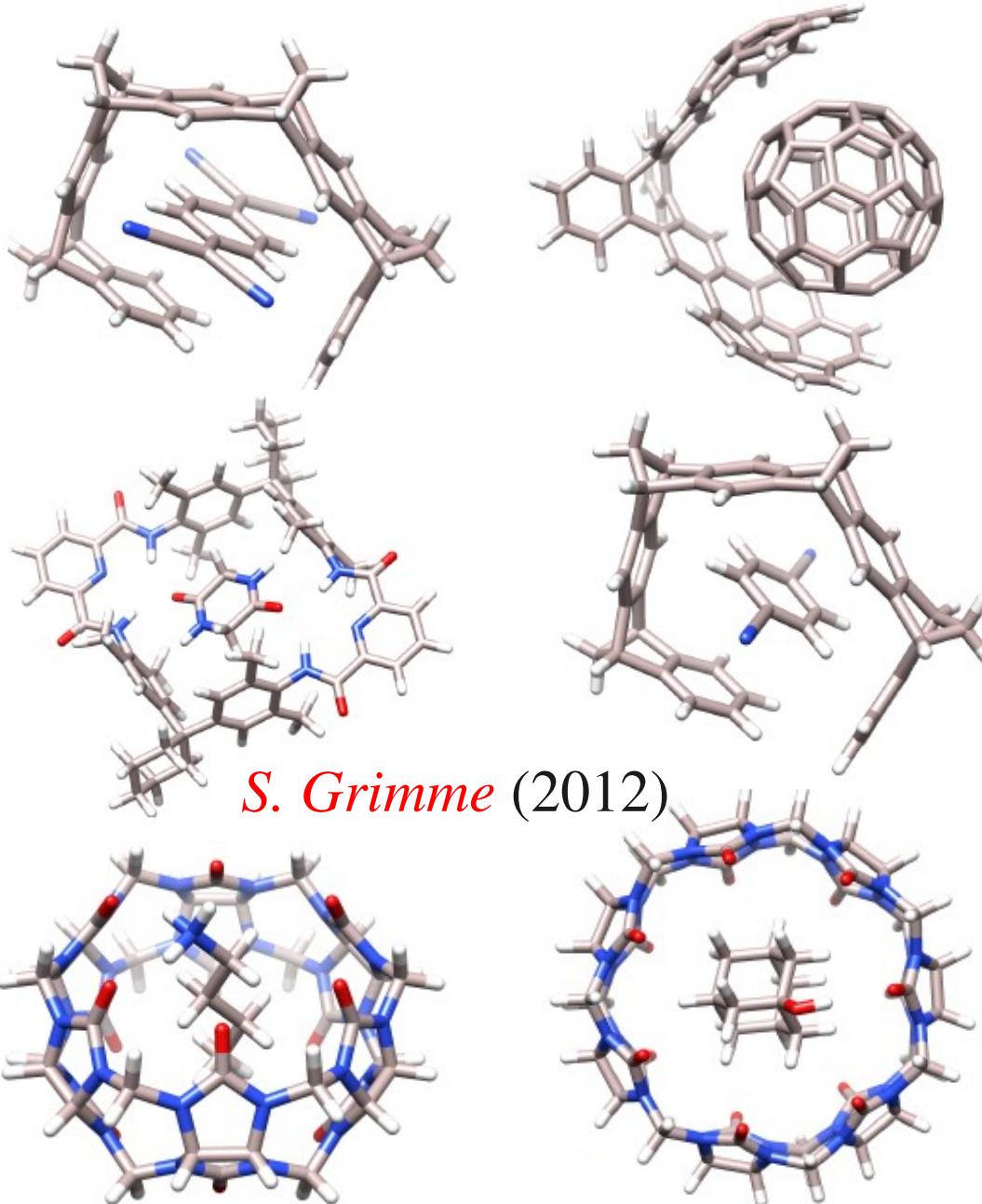
DFT+MBD correctly discriminates between crystal polymorphs: Example of glycine

MBD reduces the error in the unit cell volume to less than 0.8%!



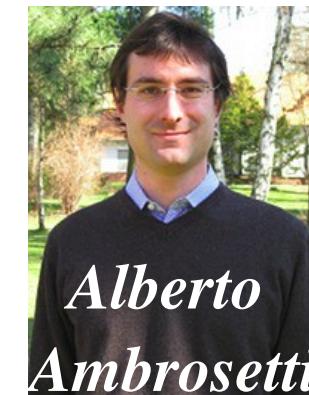
DFT+MBD gives relative energies in excellent agreement with experiment, reaching an accuracy better than 0.3 kcal/mol!

Binding in supramolecular systems: PBE+MBD performance within DMC error bar



MAE of PBE+MBD vs. DMC:
1.6 kcal/mol

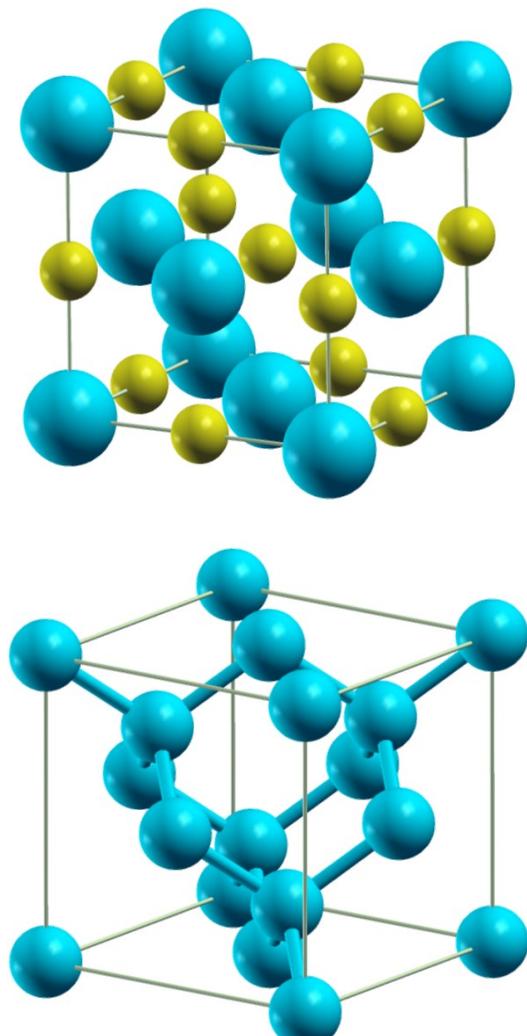
MAE of PBE+MBD vs. “Exp.”
2.9 kcal/mol



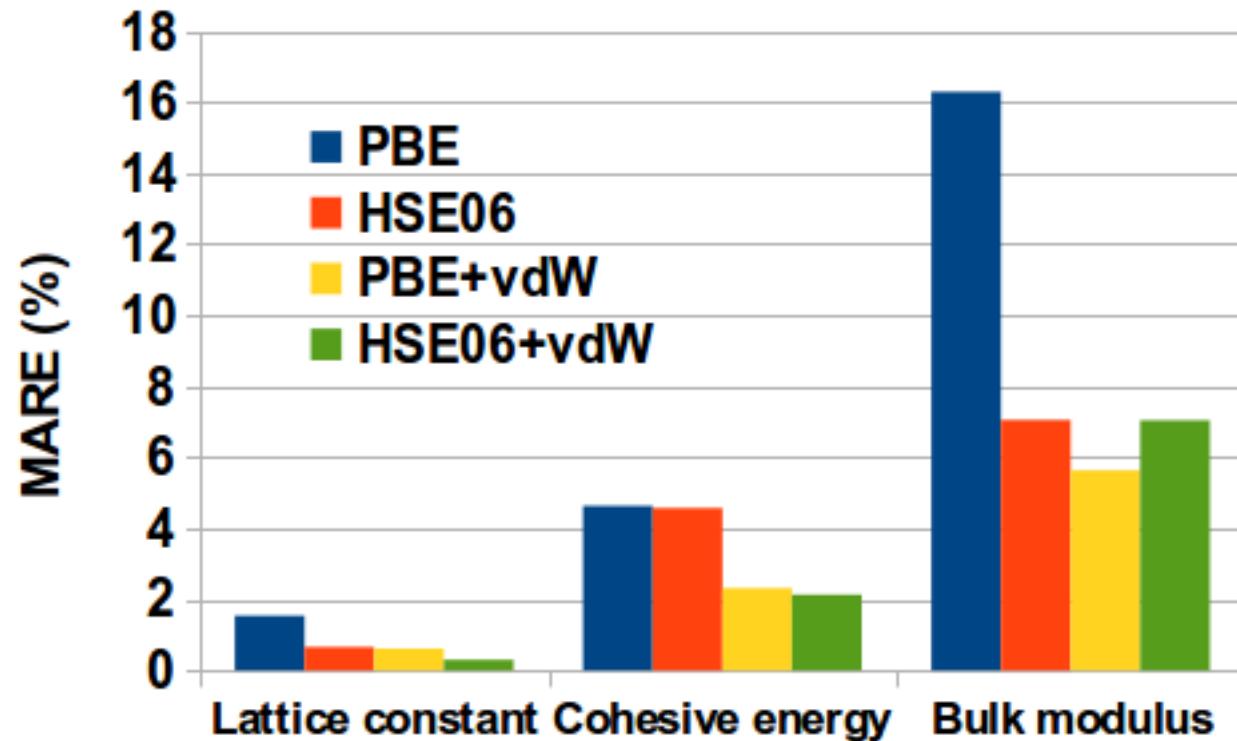


Van der Waals Interactions in Hard Solids

G.X. Zhang



Semiconductors and ionic solids

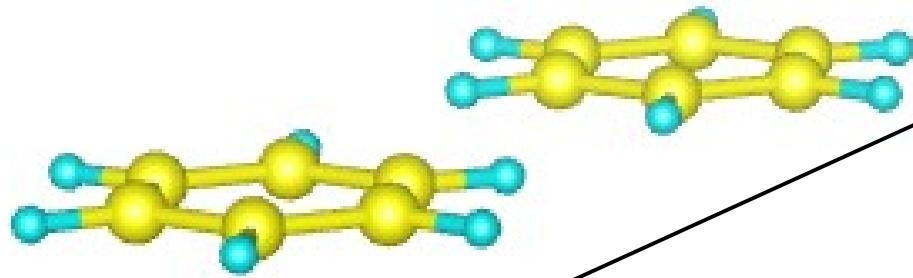


The vdW energy contributes 0.2-0.3 eV/atom to the cohesive energy, 3-17 GPa to the bulk modulus, and 0.01-0.15 Å to the lattice constant

G.-X. Zhang, A. Tkatchenko, J. Paier, H. Appel, M. Scheffler,
Phys. Rev. Lett. (2011).

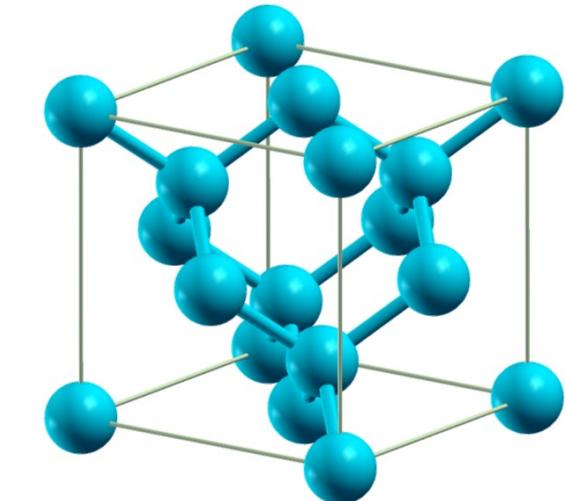
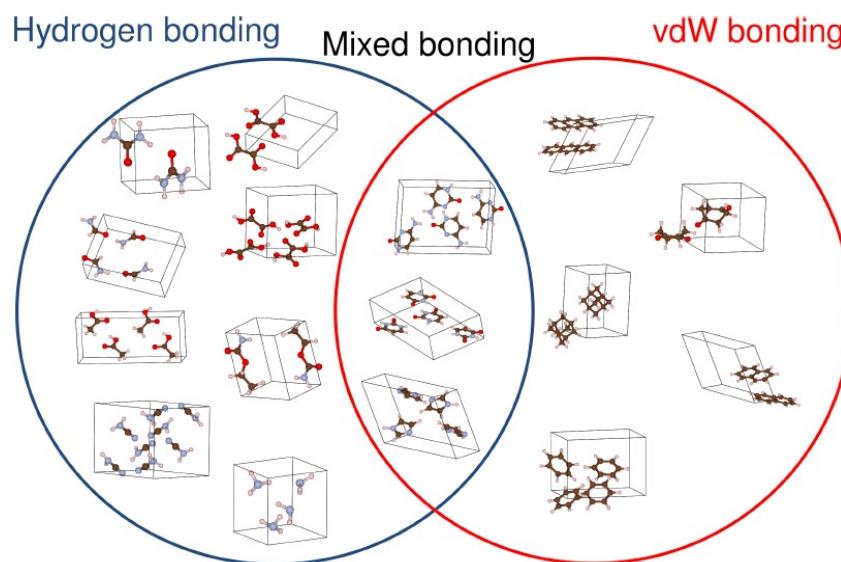
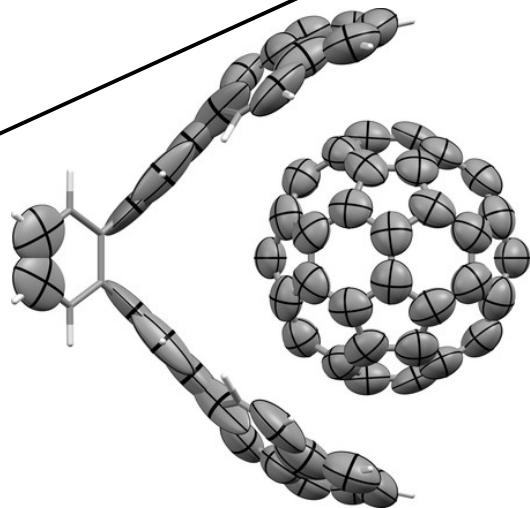
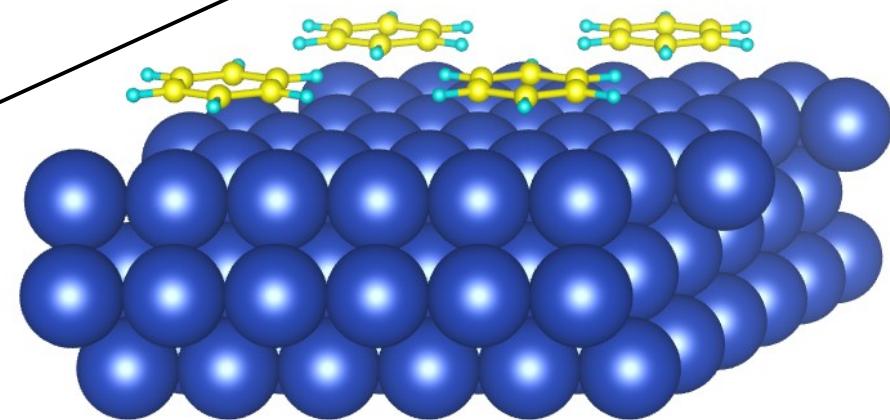
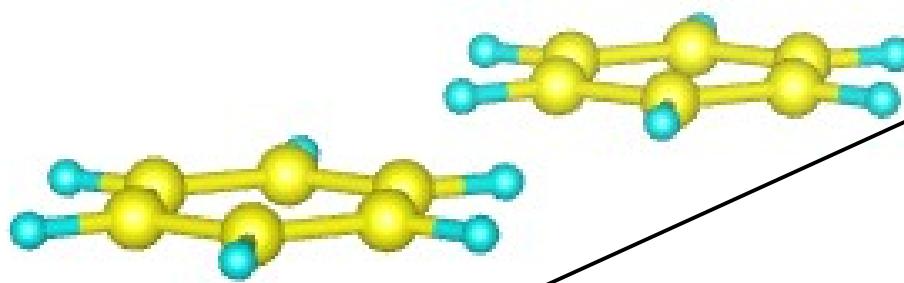
Summary

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Summary

$$C_{6,\text{eff}}^{ii} = C_{6,\text{eff}}^{ii}[n(\mathbf{r})]$$

Van der Waals interactions arise due to collective fluctuations of the electron density (they are part of long-range electron correlation)

Now, it is possible to compute the many-body vdW energy accurately and efficiently for complex (non-metallic) materials

... Perhaps it is time to move beyond rather simplified empirical pairwise models ...

$$E_c = - \int_0^\infty \frac{d\omega}{2\pi} \int_0^1 d\lambda \text{Tr} \left((\chi_\lambda(\mathbf{r}_1, \mathbf{r}_2; i\omega) - \chi_0(\mathbf{r}_1, \mathbf{r}_2; i\omega)) \frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|} \right)$$