

Hibrid kinetic schemes for modeling complex fluids

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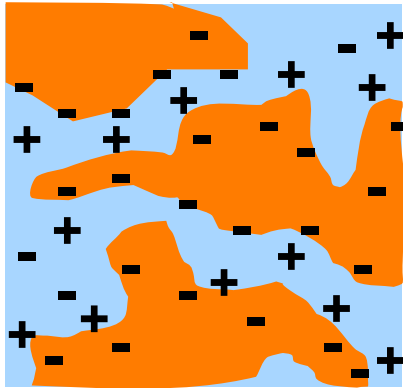
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1. Introduction

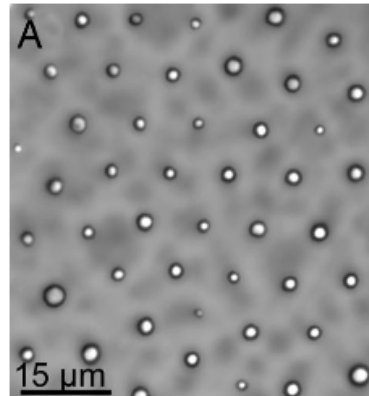
Complex fluids: Interfaces, fluctuations, instabilities

Porous media



Rocks, membranes

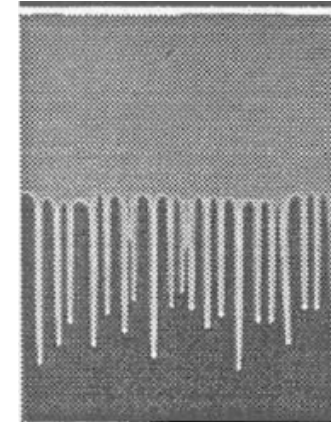
Colloidal systems



Surfactant free w/o emulsions

Leunissen et al., PNAS, 104, 2585 (2007)

Capillary flows



Dynamic wetting instabilities

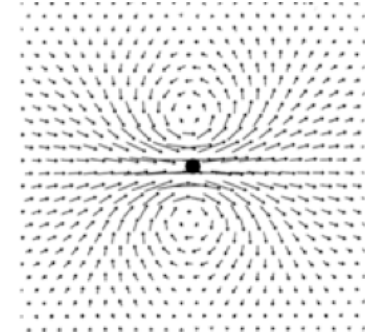
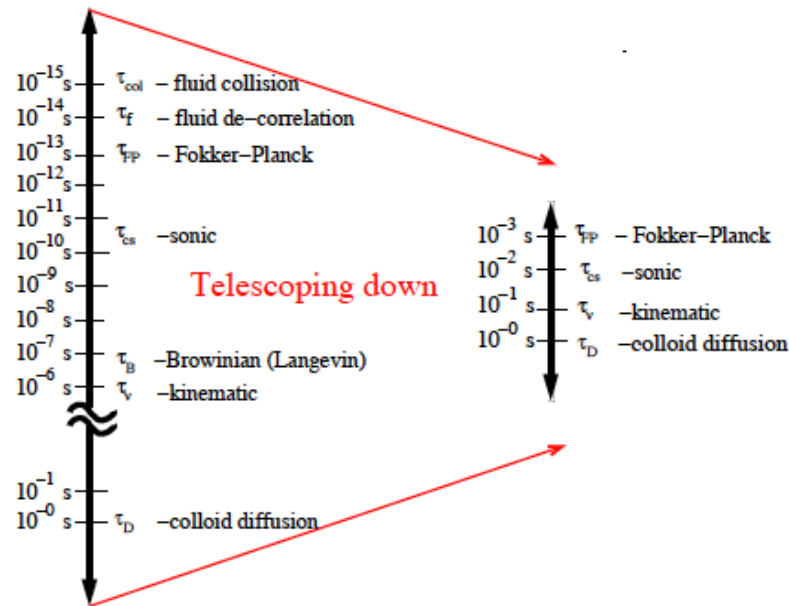
Fluid charged over Debye screening length

$k^{-1} \sim 1-10$ nm in water

Fluid/solid interactions for generic fluids

→ Coupling between ionic and solvent flows

1. Introduction



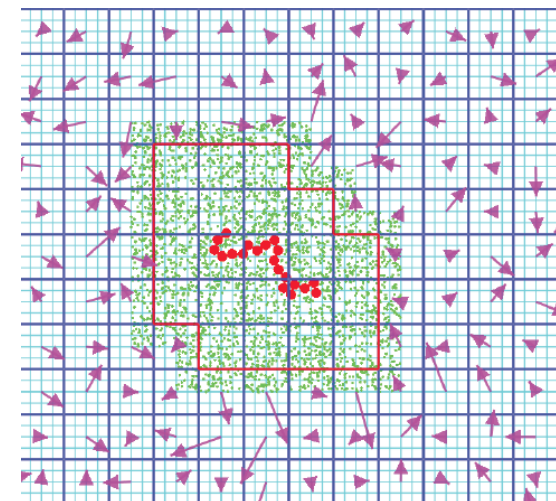
Kinetic theory

Systematic approach

Identify conserved variables

Particle based methods
 Kinetic approaches
 Fluctuating hydrodynamics +
 embedded particles

Hybrid approaches
 systematic?



1. Introduction

Why coarse-grained?

MODEL

→ No molecular structure : Local fields (densities, fluxes, ...)

→ Evolution : Principle

Composition = (simple) DDFT

Hydrodynamic description of the solvent

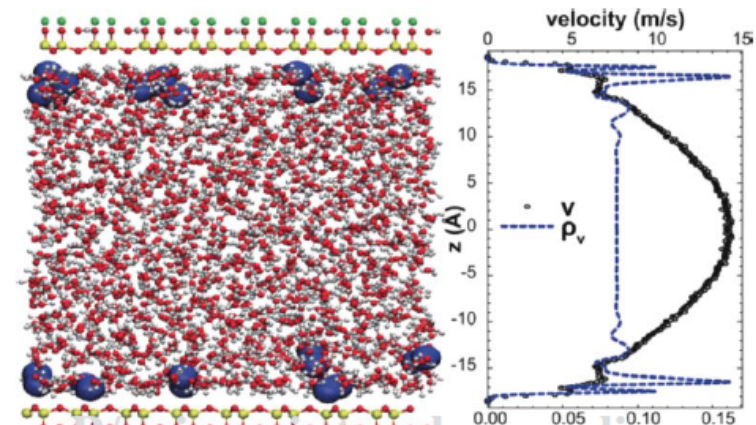
→ Evolution : Numerics

Lattice-based algorithms

Strict charge conservation

LIMITATIONS

→ No correlations



1. Introduction

Simulating coupled solvent and ionic dynamics

Goals

Describe correctly

- Advection, diffusion and migration of ions
- Oil / Water interface, solvation of ions
- Forces acting on the fluid
- Boundary conditions

Compute observables

- Collective properties : $\rho(r)$, $\rho_{\pm}(r)$, $v(r)$, κ , σ , ...
- Individual trajectories ? $D(t)$, D_e

Tools

Lattice Boltzmann, Link flux and Moment Propagation⁵

2. Kinetic models

Models motivated from kinetic theory:

Statistical description of exact dynamical evolution of a system

Long tradition back to Boltzmann: solid starting point

Fokker-Planck / diffusion equations as a limit

Why is such an approach useful for complex fluids?

-systematic procedure to integrate out degrees of freedom
in an operational/systematic way

-underlying solid theoretical treatments

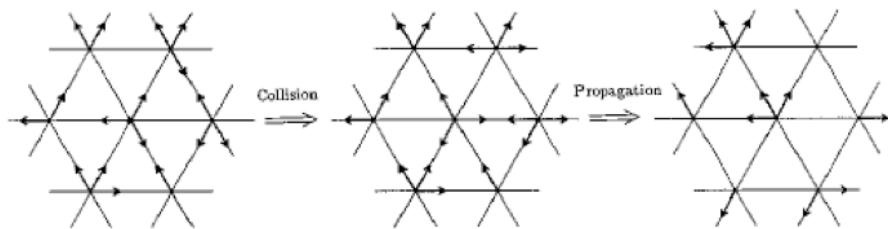
2. Lattice Gases

1986: Lattice Gas Cellular Automaton (LGCA)

Individual particles in lattice nodes

Discrete positions/velocities
Each time step: local collision

advection to nearest neighbours
completely local dynamics



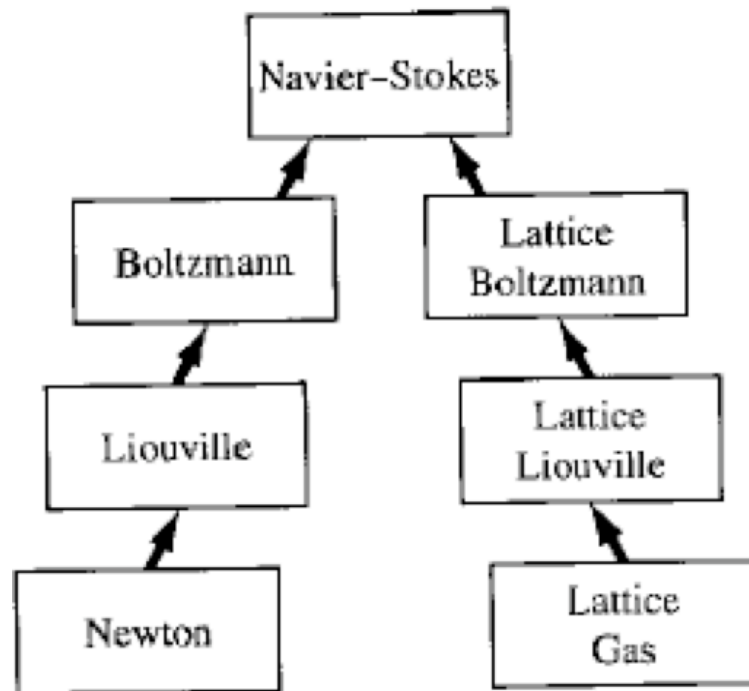
Large expectations: High Re \rightarrow turbulence

Drawbacks: No Galilean invariance (for all lattices)
Pressure dependence on velocity
Difficult to reach low viscosities
Need of large statistics (noise)



Frisch/Hasslacher/Pomeau

2. Lattice Gases



Solid theoretical basis

Parallel to basic
statistical physics

“Lattice kinetic theory”

From Fermi to Boltzmann

2. Lattice Boltzmann

$$f_i(\vec{r} + \vec{c}_i, t + 1) = f_i(\vec{r}, t)^* = f_i(\vec{r}, t) + \sum_j L_{ij} [f_j(\vec{r}, t) - f_j^{eq}(\vec{r}, t)] + F_i$$

Lattice occupation
density

Euler (explicit) update
Simplest rule
Improved integrators

No Boolean variables

Velocity averages
Hydrodynamic variables

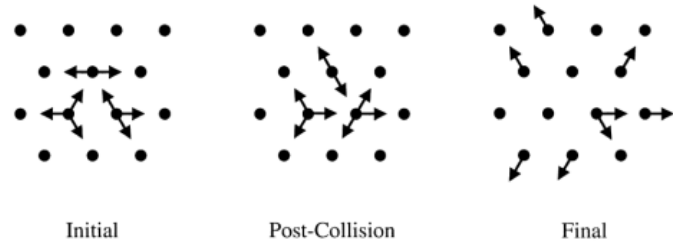
$$\begin{aligned}\sum_i f_i(\vec{r}, t) &= \rho(\vec{r}, t) \\ \sum_i f_i(\vec{r}, t) \vec{c}_i &= \rho \vec{u}(\vec{r}, t) = \vec{j} \\ \sum_i f_i(\vec{r}, t) \vec{c}_i \vec{c}_i &= \rho \vec{u} \vec{u}(\vec{r}, t) + \vec{P}(\vec{r}, t)\end{aligned}$$

Lattice geometry DnQm

Symmetry requirements to avoid lattice anisotropies

3. Lattice Boltzmann Model

Lattice kinetic model: "microscopic" dynamics



$$f_i(r + c_i, t + 1) = f_i(r, t) - \omega [f_i(r, t) - f_i^{eq}(r, t)]$$

$$\sum f_i = \rho$$

Conserved variables
Proper symmetries

Hydrodynamic equations

$$\sum f_i c_i = \rho v$$

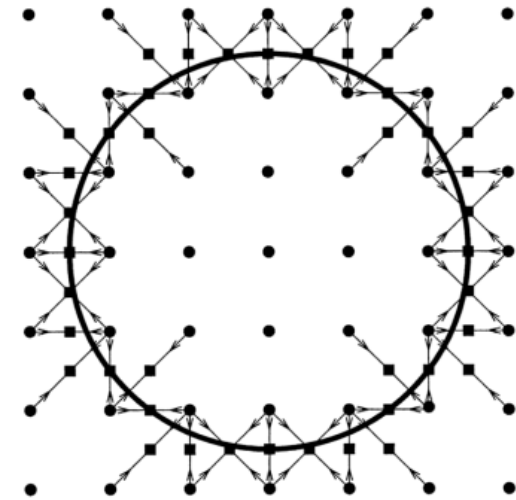
$$\sum f_i c_i c_i = \rho v v + P$$

Colloid
rigid hollow surface

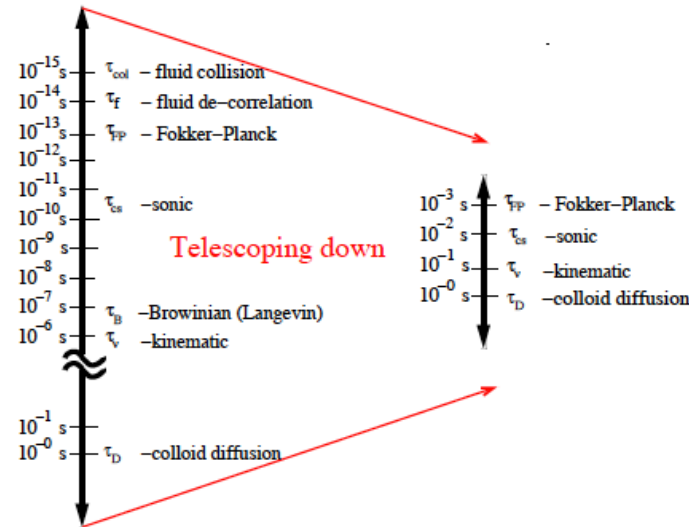
collision
bounce-back

molecular dynamics

Hybrid scheme:
Pre-selection of relevant degrees of freedom



3. Lattice Boltzmann Model



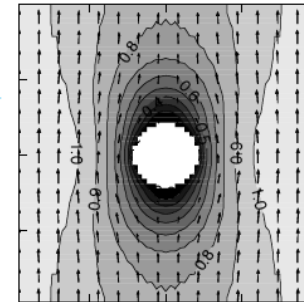
How to bridge scales?
reduce separation

$$\tau = \sigma / c_s \ll \tau = \sigma^2 / \nu \ll \tau = \sigma^2 / D$$

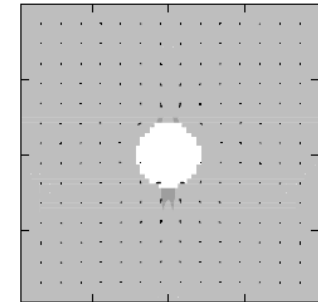
Importance to keep proper hierarchy
length
time scales

Impossibility to capture real parameters
how far can we take it?
slow speed of sound
wide interfaces
finite Reynolds numbers

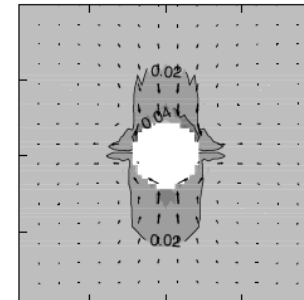
a) $Re = 3 \times 10^{-6}$



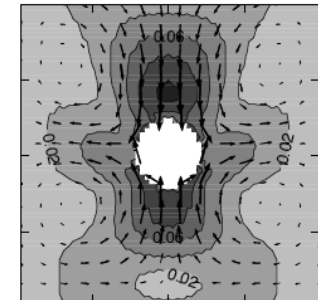
b) $Re = 0.008$



c) $Re = 0.08$



d) $Re = 0.8$



Effect of finite Re
 $Re = U\sigma/\nu \ll 1$ $Pe = U\sigma/D < 1$

Sedimenting sphere

3. Lattice Boltzmann Model

Thermodynamics of non ideal mixtures

Free energy $\beta f = \sum_{k=s,\pm} \rho_k (\ln \Lambda^3 \rho_k - 1) + \beta f^{ex}$

Mean-field excess term $\beta f^{ex} = \frac{1}{2} \beta e \psi \sum_k \rho_k z_k$

$$\beta f^{ex} = -\frac{1}{2} B \phi^2 + \frac{1}{4} B \phi^4 + \frac{\kappa}{2} (\nabla \phi)^2$$

Other possibilities, e.g. Landau-Ginzburg (oil/water, polymers)

Chemical potential and pressure

$$\beta \mu_k = \frac{\delta(\beta f)}{\delta \rho_k} = \ln \rho_k + \beta \mu_k^{ex} \quad \rightarrow \quad \beta \mu_k = \ln \rho_k + \beta z_k e \psi$$

$$\beta \nabla p = \sum_k \rho_k \beta \nabla \mu_k = \beta \nabla p^{id} + \sum_k \rho_k \beta \nabla \mu_k^{ex} \quad \text{Gibbs-Duhem}$$

3. Lattice Boltzmann Model

Link flux method for solutes

Principle

Integrated conservation law $\partial_t \int_{V_0} \rho_k dV = - \oint_{A_0} \mathbf{j}_k \cdot \mathbf{n} dA$

Algorithm

$$n_k(\mathbf{r}, t + \Delta t) - n_k(\mathbf{r}) = -A_0 \sum_i j_{ki}(\mathbf{r})$$

Link fluxes

Discretized $\mathbf{j}_k = -D_k e^{-\beta\mu_k^{ex}} \nabla \left[\rho_k e^{+\beta\mu_k^{ex}} \right]$

$$j_{ki}(\mathbf{r}) = -d_k \frac{e^{-\beta\mu_k(\mathbf{r})} + e^{-\beta\mu_k(\mathbf{r}+\mathbf{c}_i\Delta t)}}{2} \times \left[\frac{n_k(\mathbf{r} + \mathbf{c}_i\Delta t)e^{\beta\mu_k(\mathbf{r}+\mathbf{c}_i\Delta t)} - n_k(\mathbf{r})e^{\beta\mu_k(\mathbf{r})}}{\|\mathbf{c}_i\Delta t\|} \right]$$

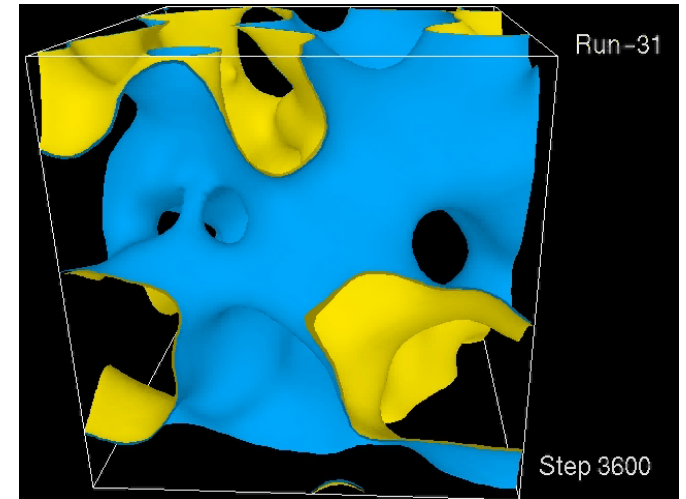
→ $j_{ki}(\mathbf{r}) = -j_{ki'}(\mathbf{r} + \mathbf{c}_i\Delta t)$

→ Boltzmann distribution in equilibrium

4. Dynamics of complex fluids

Controlling materials through external fields
shearing
gravity
electric/magnetic fields

Use internal capabilities of materials



Temperature quench

internal structure

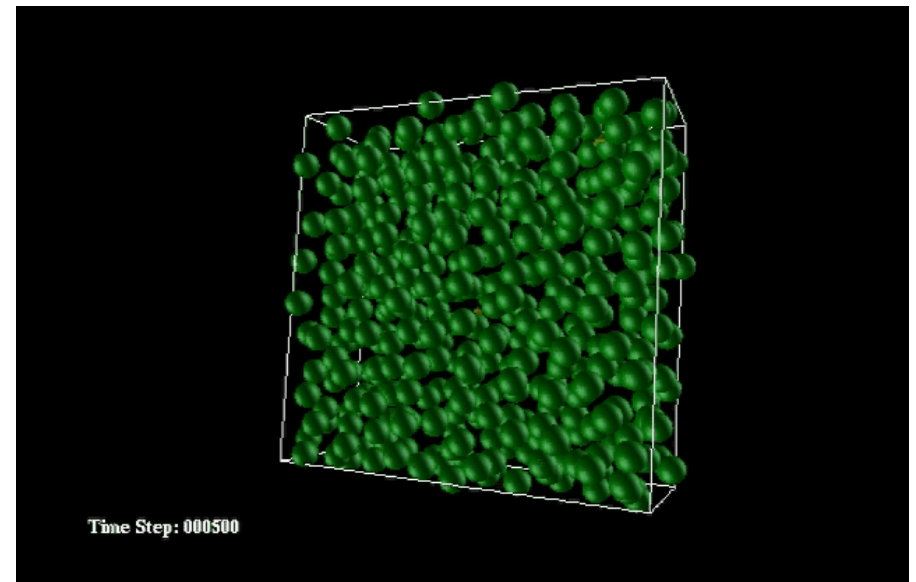
length/time scale competition

emerging structures

Spontaneous internal motion

actuated/internal propulsion

interaction with actuating fields



K. Stratford et al. Science (2005)

4. Dynamics of complex fluids

$$L_0 = \eta^2 / (\rho\sigma)$$

material parameters

$$T_0 = \eta^3 / (\rho\sigma^2)$$

Universal scaling

Viscous regime

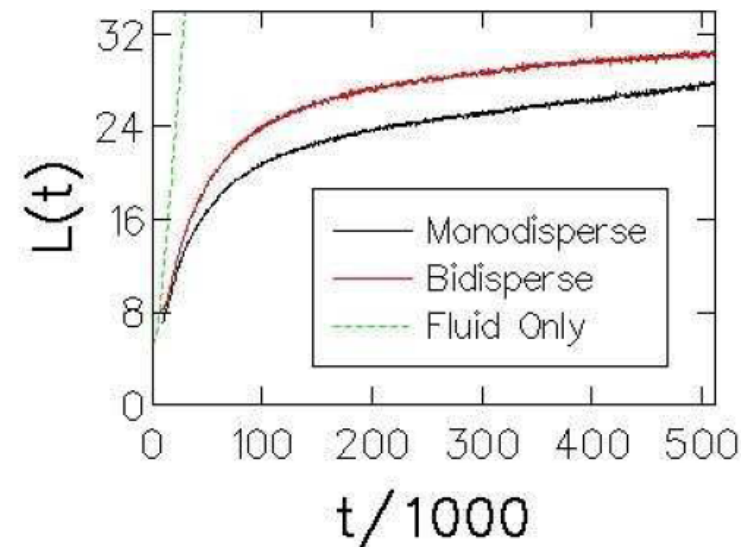
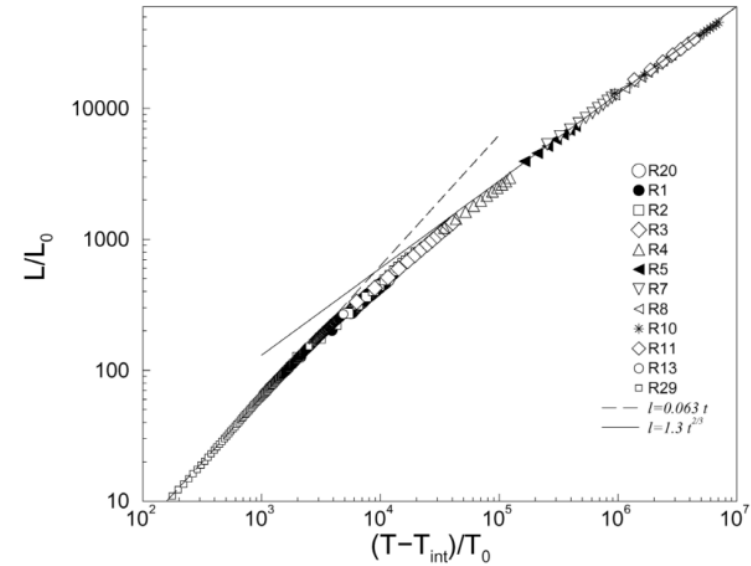
$$L / L_0 \sim A(T / T_0)^1$$

Inertial regime

$$L / L_0 \sim B(T / T_0)^{2/3}$$

Crossover regime

Colloids arrest phase separation
gel formation



5. Mixtures of oil, water and ions

Free energy model

$$F[\phi, \rho_+, \rho_-] = \int d\mathbf{r} \left[-\frac{1}{2}B\phi^2 + \frac{1}{4}B\phi^4 + \frac{\kappa}{2}(\nabla\phi)^2 \right] + \sum_{\alpha=\pm} \int d\mathbf{r} \rho_{\alpha}(\mathbf{r}) \left[k_B T \left(\ln \frac{\rho_{\alpha}(\mathbf{r})}{\rho_s} - 1 \right) + f_{\alpha}(\mathbf{r}) + \frac{z_{\alpha}e}{2}\psi(\mathbf{r}) \right]$$

and parametrize ε and f_{α} as a function of composition

Contains

Immiscibility of solvents

Solvation of ions $f_{\pm} \propto \Delta\mu_{\pm} = \mu_{\pm}^o - \mu_{\pm}^w$

Electrostatics, including without ions (dielectrophoresis)

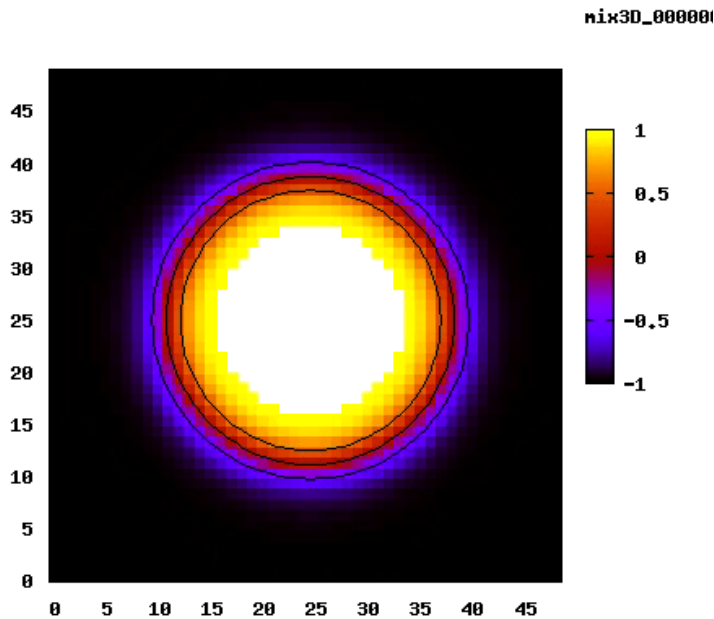
Dynamics of composition parameter = Cahn-Hilliard

5. Mixtures of oil, water and ions

Uncharged oil/water mixtures

Differential permittivity

Applied E field

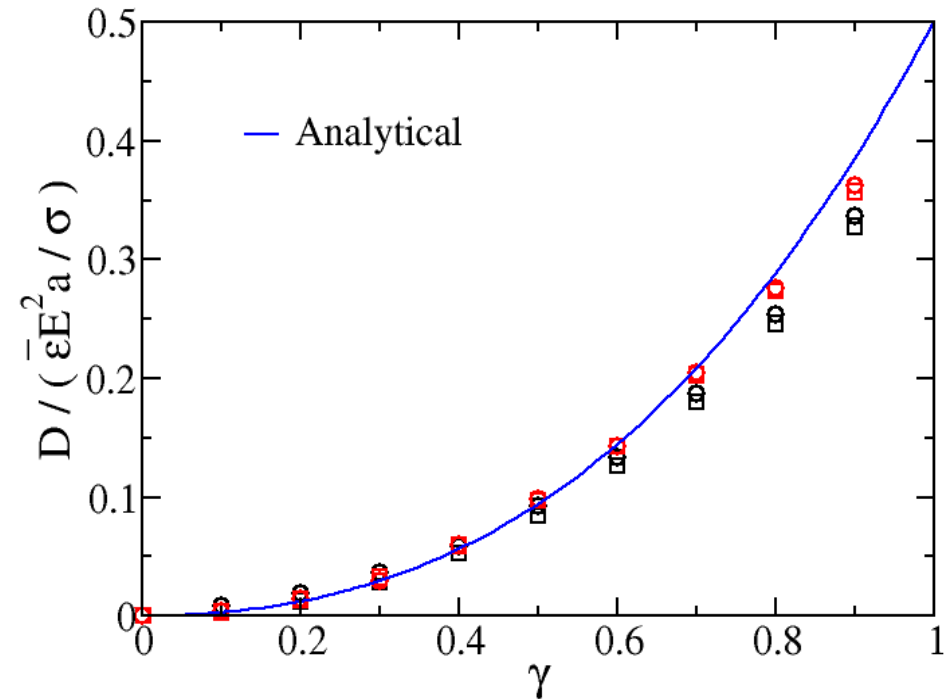


Oil in water droplet deformation

Transient hydrodynamic flow

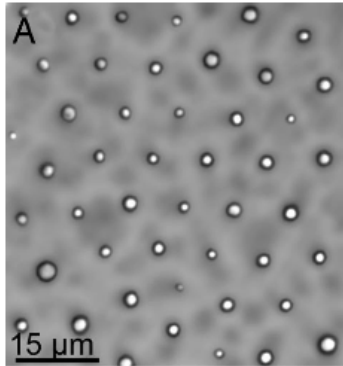
Small deformations

$$D = \frac{\bar{\epsilon} E^2 a}{4\sigma} \times \gamma^2 (1 + \gamma)$$



with $\gamma = \frac{\epsilon_w - \epsilon_o}{\epsilon_w + \epsilon_o}$

5. Mixtures of oil, water and ions



Anions and cations with different solubility

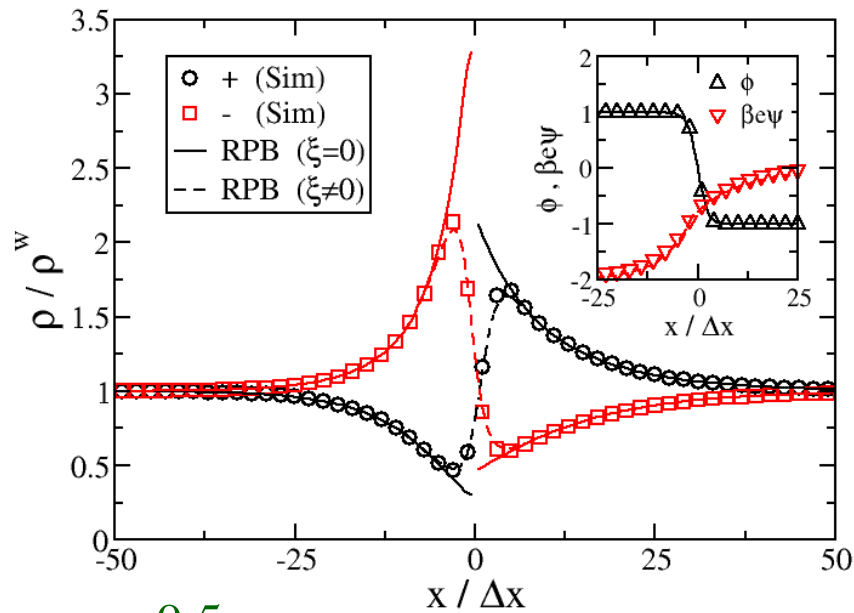
→ Spontaneous charge separation

Leunissen et al., PNAS, 104, 2585 (2007)

Electroosmotic flow

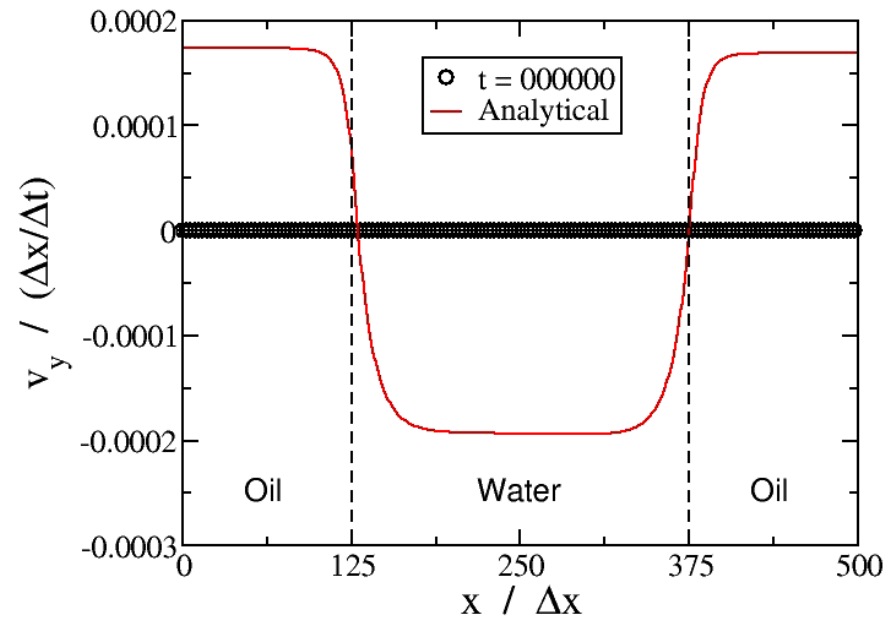
$\gamma = 0.5$; $\beta\Delta\mu_+ = +2$; $\beta\Delta\mu_- = -2$

Ionic profiles



$\gamma = 0.5$

$\Delta\mu_{\pm} = \pm 2 k_B T$: hydrophilic cation & hydrophobic anions



6. Dynamic wetting instability

Non linear regime

$$\theta_E = 90^\circ$$



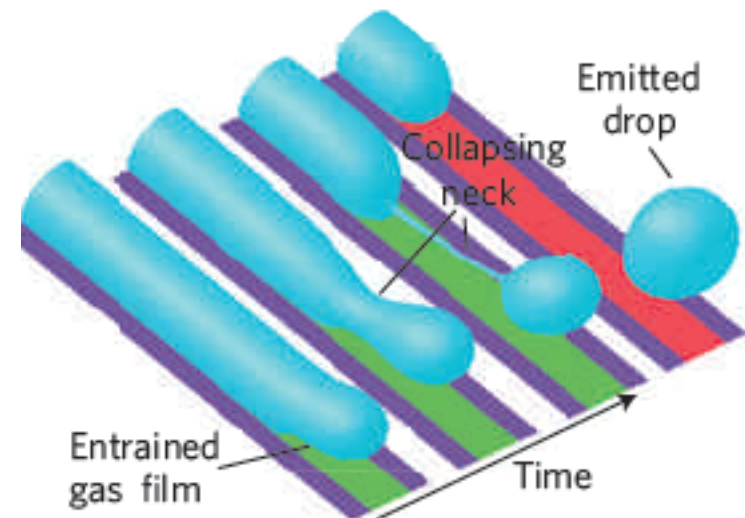
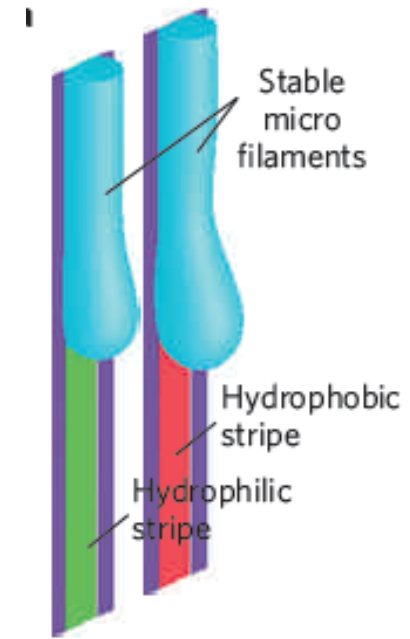
Forced thin film or rivulet

Stabilized by incoming flow

Localized by wetting mismatch

Due to substrate
hydrophobicity/hydrophilicity

Different from Rayleigh-Plateau



6. Dynamic wetting instability

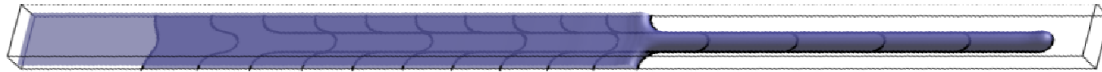
Non linear regime

$$\theta_E = 90^\circ$$

Base

$$L \sim t$$

Tip



Forced thin film or rivulet

Stabilized by incoming flow

Localized by wetting mismatch

Due to substrate

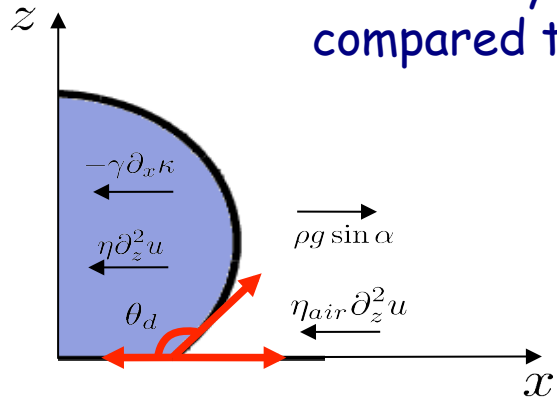
hydrophobicity/hydrophilicity

Different from Rayleigh-Plateau



6. Dynamic wetting instability

Instability induced when forcing on protusion too large compared to friction at contact line



Capillary pressure

Gravity

+

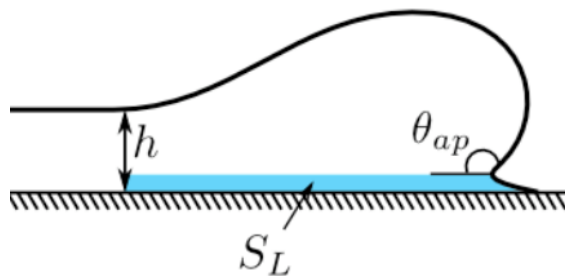
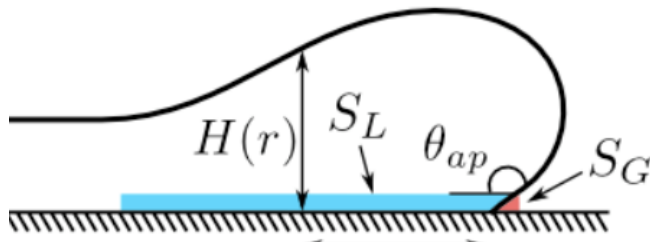
Dissipation in the film

Dissipation in the air

Contact line force

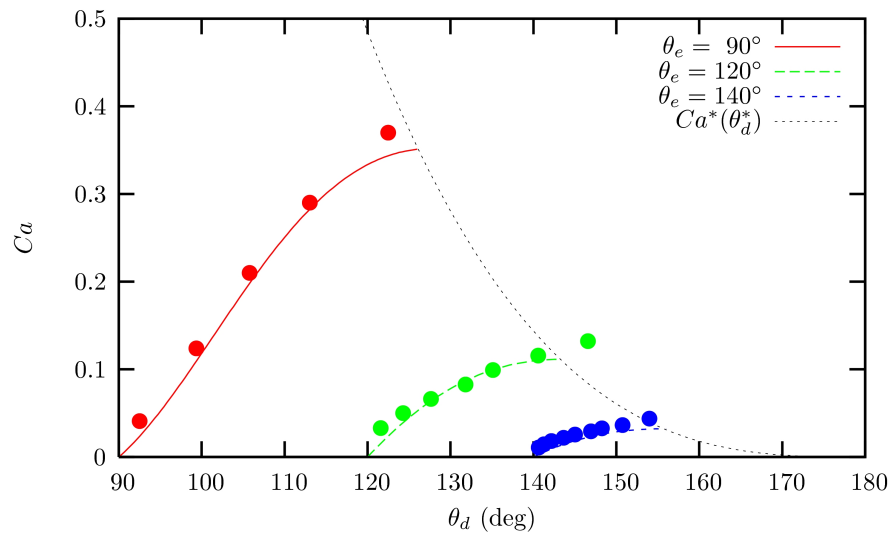
$$\theta_d = f(Ca)$$

$$Ca = \frac{\eta U}{\gamma}$$



Shape of contact line changes with attraction to solid substrate

6. Dynamic wetting instability



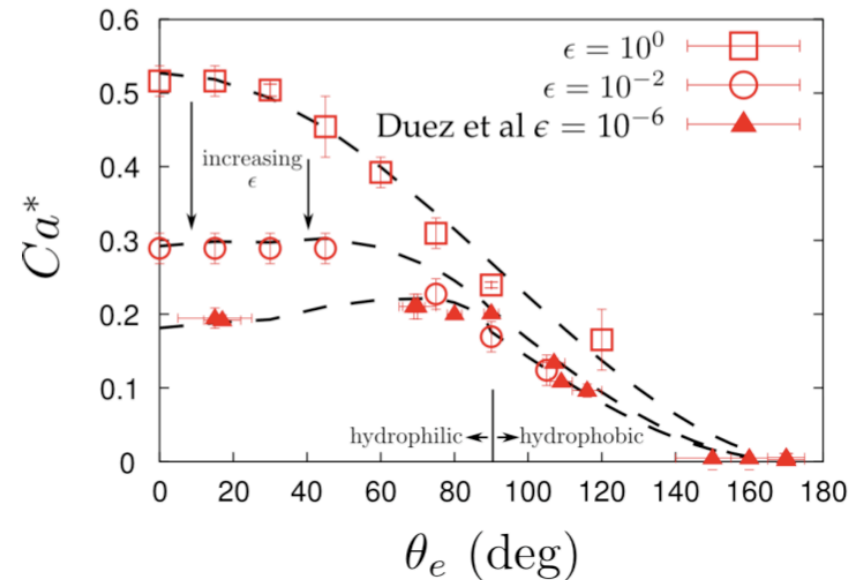
Identify critical Ca

Compares with lubrication theory

Larger angle requires less forcing

Drop emission

Periodic process



Hydrophobic substrate

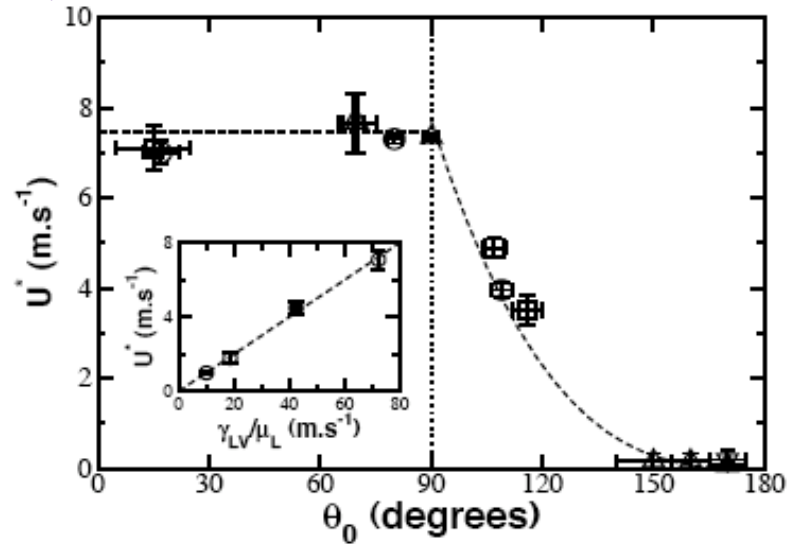
Controlled by gas wedge

Hydrophilic substrate

Foot explains plateau

6. Dynamic wetting instability

Bocquet et al. (2007)



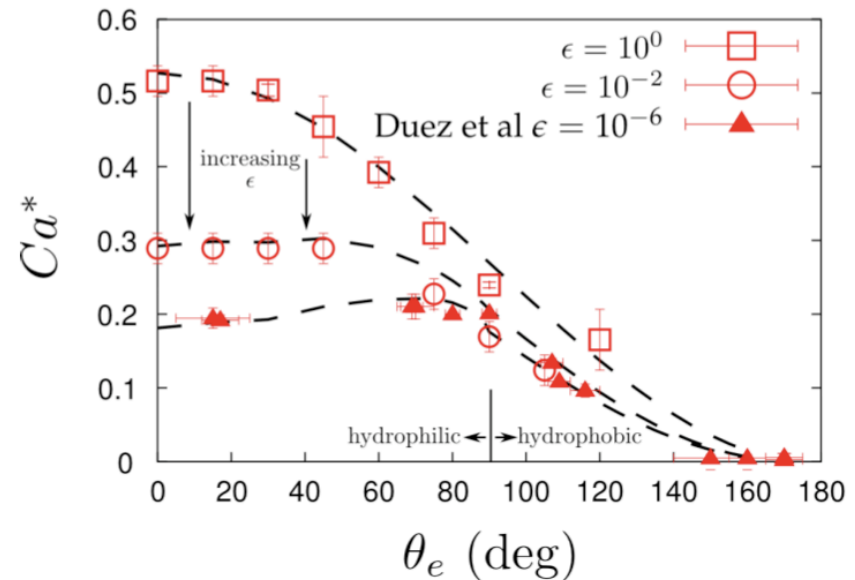
Identify critical Ca

Compares with lubrication theory

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Periodic process



Hydrophobic substrate

Controlled by gas wedge

Hydrophilic substrate

Foot explains plateau

7. Moment propagation

Principle

Probabilistic interpretation LB

In addition to $f_i(\mathbf{r}, \mathbf{c}_i, t)$ for the fluid (\sim solvent), scalar $P(\mathbf{r}, t)$ e.g. ρ_k

$$P(\mathbf{r}, t + 1) = \sum_i P(\mathbf{r} - \mathbf{c}_i \Delta t, t) p_i(\mathbf{r} - \mathbf{c}_i \Delta t, t) + P(\mathbf{r}, t) \left(1 - \sum_i p_i(\mathbf{r}, t) \right)$$

$p_i(\mathbf{r}, t)$ = probability to jump from \mathbf{r} to $\mathbf{r} + \mathbf{c}_i \Delta t$

$$p_i(\mathbf{r}, t) = \frac{f_i(\mathbf{r}, t)}{\rho_f(\mathbf{r}, t)} - w_i + \lambda w_i \left\{ \frac{1}{4} \beta q \mathbf{E} \cdot \mathbf{c}_i \Delta t + \frac{1}{1 + e^{-\beta[V(\mathbf{r}) - V(\mathbf{r} + \mathbf{c}_i \Delta t)]}} \right\}$$

Includes interactions + sensitivity to applied external fields

Ensures detailed balance

7. Moment propagation

Based on transition probabilities

Propagate other quantities using same scheme

Correlation functions

Propagate $P(\mathbf{r}, t)$ = probability to arrive in \mathbf{r} at t , weighted by $v(0)$

$$\langle v_\alpha(t)v_\alpha(0) \rangle = \sum_{\mathbf{r}} P(\mathbf{r}, t) \times \left(\sum_i p_i(\mathbf{r}, t) c_{i\alpha} \right)$$

Average over all possible trajectories done at once

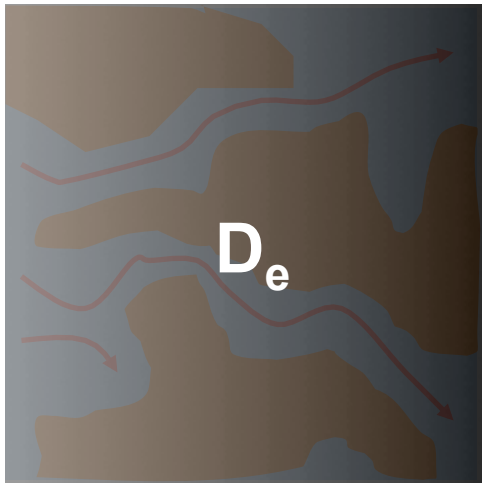
Choose equilibrium initial condition

$$P(\mathbf{r}, 1) = \sum_i \frac{e^{-\beta V(\mathbf{r} - \mathbf{c}_i \Delta t)}}{Q} p_i(\mathbf{r} - \mathbf{c}_i \Delta t) c_{i\alpha}$$

$$Z_\alpha(0) = \sum_{\mathbf{r}, i} \frac{e^{-\beta V(\mathbf{r})}}{Q} p_i(\mathbf{r}) c_{i\alpha} c_{i\alpha}$$

8. Diffusion in porous medium

Charged Tracer diffusion in a clay

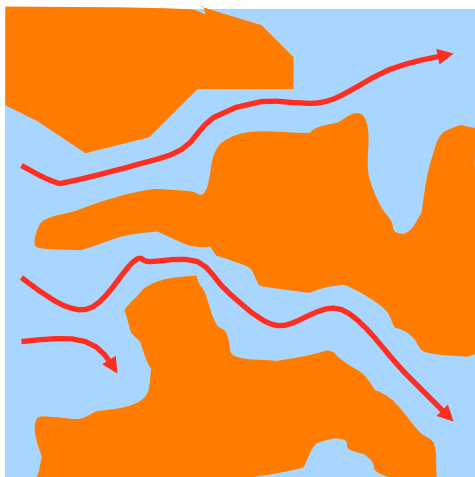


$$\frac{D_e^-}{D_0^-} < \frac{D_e^w}{D_0^w} < \frac{D_e^+}{D_0^+}$$

slower

faster

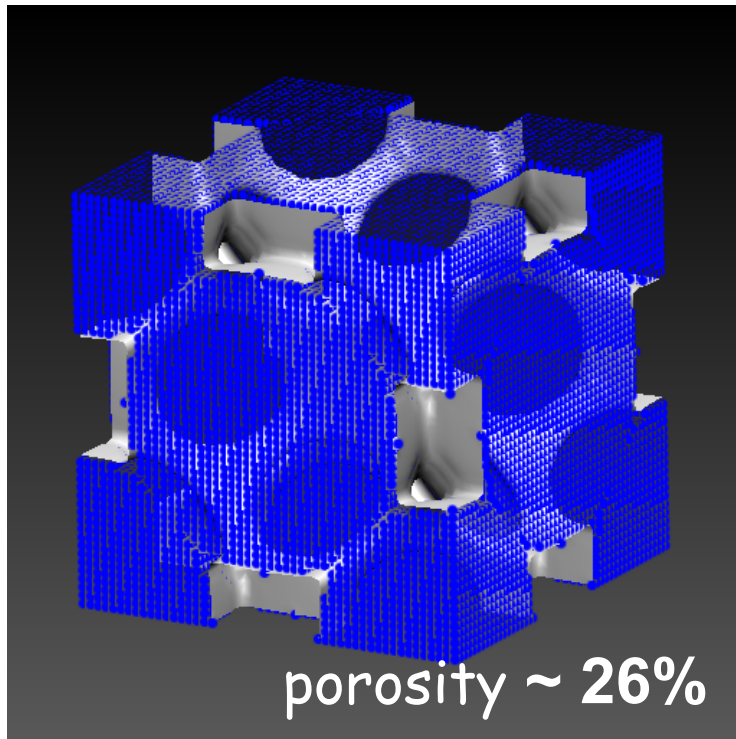
Effect decreases with ionic strength



Insight looking at pore scale?

8. Diffusion in porous medium

Model porous medium



FCC lattice. Spheres $Z < 0$

Counterions + salt (controls κ)

Large pores connected by small ones

Continuous double layer

Tracer of charge $z \in \{-1, 0, +1\}$
same D_0

Effective diffusion coefficient?

Relation to explored porosity?

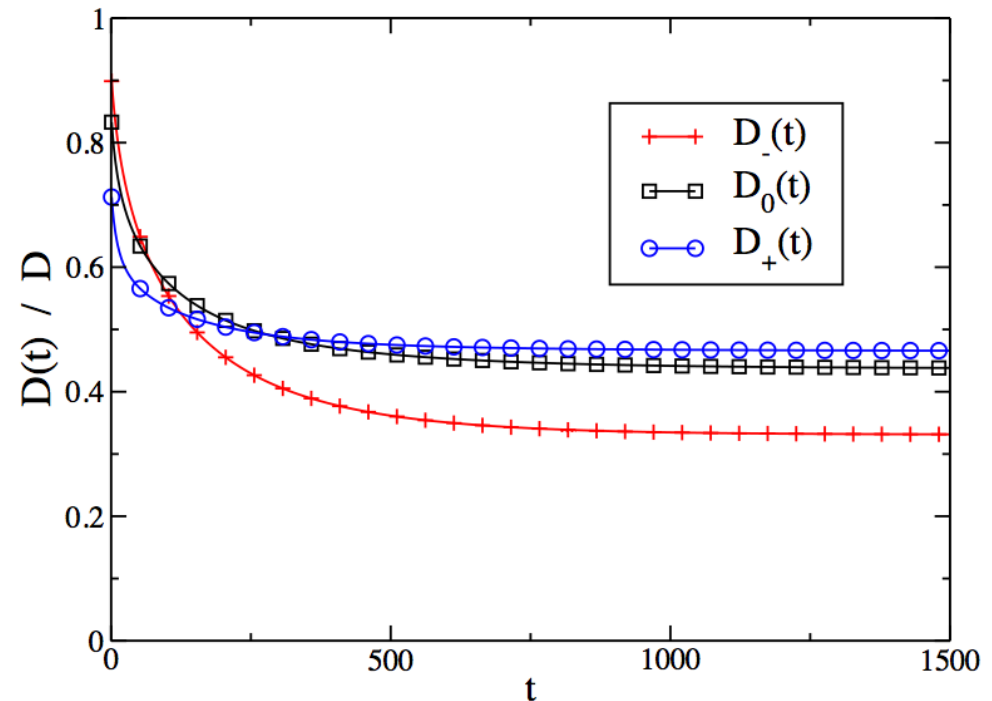
Time dependent diffusion, $D(t)$

8. Diffusion in porous medium

$D(t)$ and D_e

$$D_e = \lim_{t \rightarrow \infty} D(t)$$

$$D_+ > D_0 > D_-$$



Characteristic time

To explore the porosity accessible to each tracer

$$\tau = \int_0^{\infty} \frac{D(t) - D_e}{D - D_e} dt \quad \longrightarrow \quad \tau_+ < \tau_0 < \tau_-$$

4. Diffusion in porous medium

Observations

1) $D_+ > D_0 > D_-$

2) $\tau_+ < \tau_0 < \tau_-$

Cations

Close to surfaces

Double-layer = « highway »

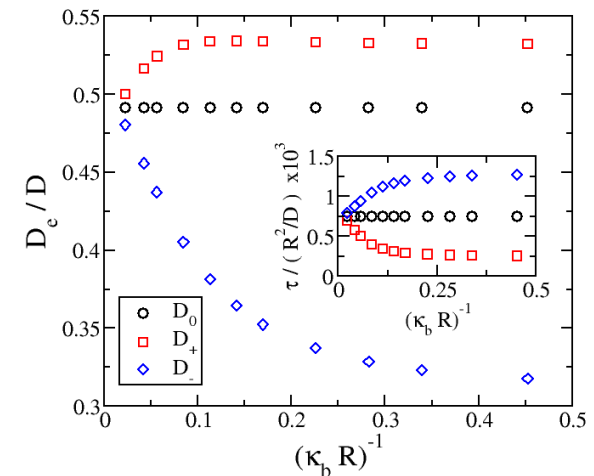
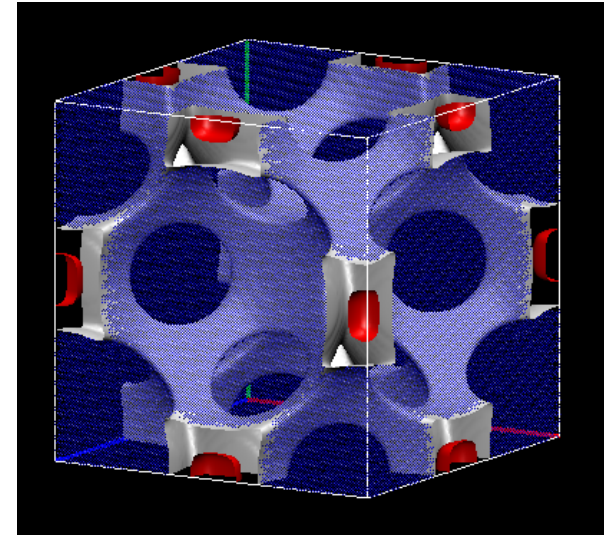
small, connected volume (small t , large D)

Anions

Far from surfaces

« Electrostatic bottlenecks » = barrier to jump

from one cavity to the next (larger t , smaller D)



8. Diffusion in porous medium

Couple localized species at liquid/solid boundary

Surface diffusion
+ Local Langmuir kinetics

$$\partial_t \Gamma(\mathbf{r}, t) = -k_d \Gamma(\mathbf{r}, t) + k_a \rho(\mathbf{r}, t)$$

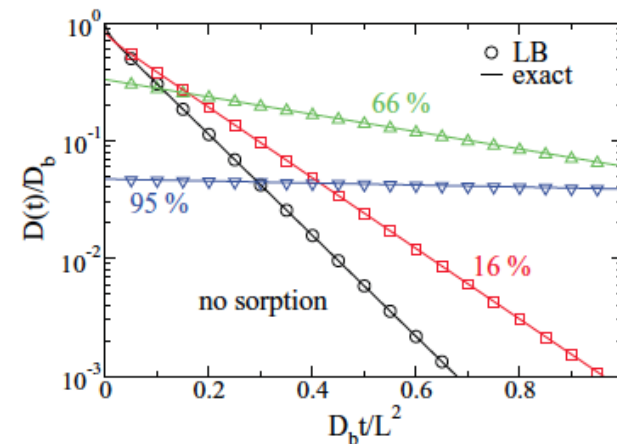
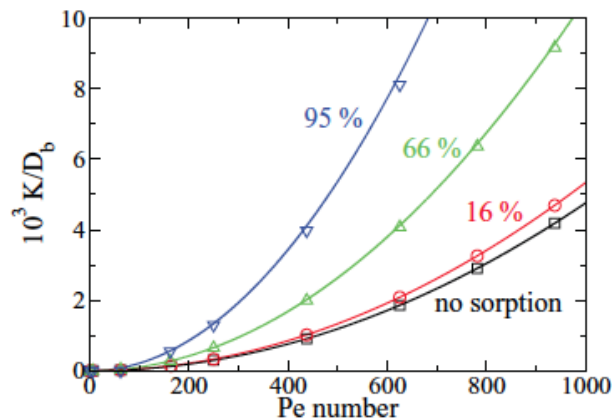
Local kinetic rules
efficient performance

$$P_{\text{ads}}(\mathbf{r}, t + \Delta t) = P(\mathbf{r}, t)p_a + P_{\text{ads}}(\mathbf{r}, t)(1 - p_d),$$

$$P(\mathbf{r}, t + \Delta t) = P^*(\mathbf{r}, t + \Delta t) + P_{\text{ads}}(\mathbf{r}, t)p_d.$$

Bulk vs Surface transport in porous medium?

New modes of transport?
Resonances in oscillatory driving?
Enhanced Taylor dispersion?



9. Fluctuating Lattice Boltzmann

Thermal fluctuations absent in LB

From Boltzmann equation recover Navier-Stokes

Need effect of other particles on single particle distribution function

Zwanzig-Bixon

recover fluctuating hydrodynamics

random distribution: conserves mass/momentum

stress satisfies FDT

Thermal noise added as an additional external force

$$f_i(\vec{r} + \vec{c}_i, t + 1) = f_i(\vec{r}, t) + \sum_j L_{ij} [f_j(\vec{r}, t) - f_j^{eq}(\vec{r}, t)] + \xi_i$$

9. Fluctuating hybrid schemes

Need to satisfy FDT at lattice

Diffusion model for order parameter

$$F(\psi) = \int [f(\psi) + \frac{K}{2} (\nabla\psi)^2] d\mathbf{r}$$

Equilibrium spectrum

$$f(\psi) = -\frac{A}{2}\psi^2 + \frac{B}{4}\psi^4$$

$$P[\psi(\mathbf{q})] = \frac{e^{\beta F_{\mathbf{q}}}}{Z}$$

Discretize on lattice?

$$\frac{\partial\psi}{\partial t} = M\nabla^2 \frac{\delta F}{\delta\psi} + \nabla \cdot \hat{\xi}$$

$$\langle \psi^2 \rangle = \frac{kT}{A}$$

$$\langle |\tilde{\psi}_{\mathbf{q}}|^2 \rangle = \frac{kT}{A + K\mathbf{q}^2}$$

9. Fluctuating hybrid schemes

How do we implement this scheme on previous link flux?

$$\partial_t \psi = \sum_i w_i \mathbf{c}_i \cdot \mathbf{j}^\mu(\mathbf{r} + \frac{1}{2} \mathbf{c}_i) + \sum_i w_i \mathbf{c}_i \cdot \hat{\xi}(\mathbf{r} + \frac{1}{2} \mathbf{c}_i)$$

Link fluxes (deterministic/random)

$$\mathbf{j}^\mu(\mathbf{r} + \frac{1}{2} \mathbf{c}_i) = M \frac{1}{2} [\nabla \mu(\mathbf{r}) + \nabla \mu(\mathbf{r} + \mathbf{c}_i)]$$

$$\nabla \cdot \mathbf{j}^\mu(\mathbf{r} + \frac{1}{2} \mathbf{c}_i) = \frac{1}{c_s^4} \sum_j w_j \mathbf{c}_j \left[M \sum_k w_k \mathbf{c}_k \mu(\mathbf{r} + \mathbf{c}_j + \mathbf{c}_k) \right]$$

$$\hat{\xi}(\mathbf{r} + \frac{1}{2} \mathbf{c}_i) = \frac{1}{2} [\hat{\xi}(\mathbf{r}) + \hat{\xi}(\mathbf{r} + \mathbf{c}_i)]$$

In Fourier they show proper symmetry on lattice

$$\partial_t \psi_{\mathbf{q}} = \Gamma_{\mathbf{q}} \cdot (M \Gamma_{\mathbf{q}} \mu_{\mathbf{q}}) + \Gamma_{\mathbf{q}} \cdot \xi_{\mathbf{q}}$$

$$\Gamma_{\mathbf{q}} \equiv \sum_i w_i \mathbf{c}_i \exp(i \mathbf{q} \cdot \mathbf{c}_i)$$

9. Fluctuating hybrid schemes

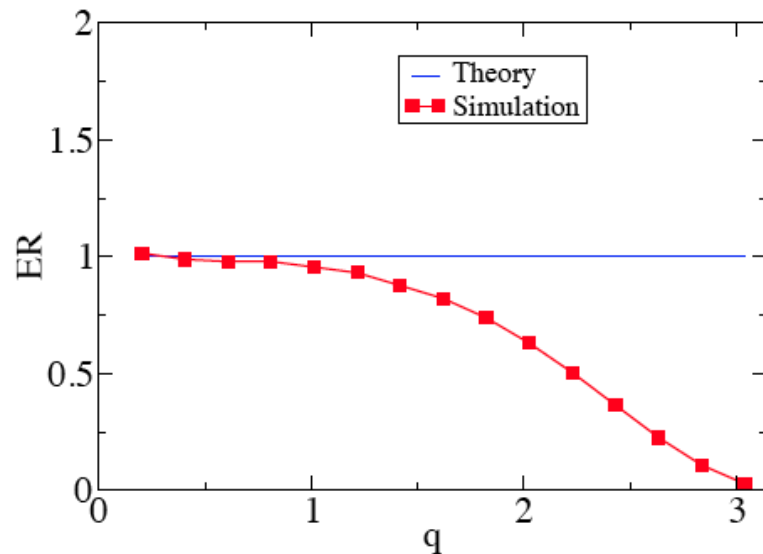
$$[\nabla^2 \psi](\mathbf{r}) = \sum_i \hat{\omega}_i \psi(\mathbf{r} + \mathbf{c}_i) \quad \longrightarrow \quad [\nabla^2 \psi](\mathbf{q}) = \sum_i \hat{\omega}_i e^{i\mathbf{q} \cdot \mathbf{c}_i} \tilde{\psi}(\mathbf{q}) = L(\mathbf{q}) \tilde{\psi}(\mathbf{q})$$

$$[\nabla \cdot \hat{\xi}](\mathbf{r}) = \sum_i \omega_i \mathbf{c}_i \cdot \hat{\xi}(\mathbf{r} + \mathbf{c}_i) \quad \longrightarrow \quad [\nabla \cdot \hat{\xi}](\mathbf{q}) = \sum_i \omega_i \mathbf{c}_i e^{i\mathbf{q} \cdot \mathbf{c}_i} \cdot \tilde{\xi}(\mathbf{q}) = \Gamma(\mathbf{q}) \cdot \tilde{\xi}(\mathbf{q})$$

$$\partial_t \tilde{\psi}(\mathbf{q}) = ML(\mathbf{q}) \tilde{\psi}(\mathbf{q}) + \Gamma(\mathbf{q}) \cdot \tilde{\xi}(\mathbf{q}).$$

FDT on lattice requires

$$L(\mathbf{q}) = \Gamma(\mathbf{q}) \cdot \Gamma(\mathbf{q}).$$

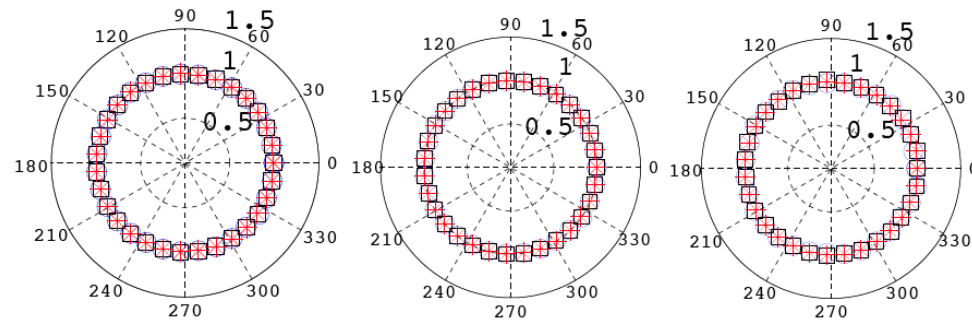
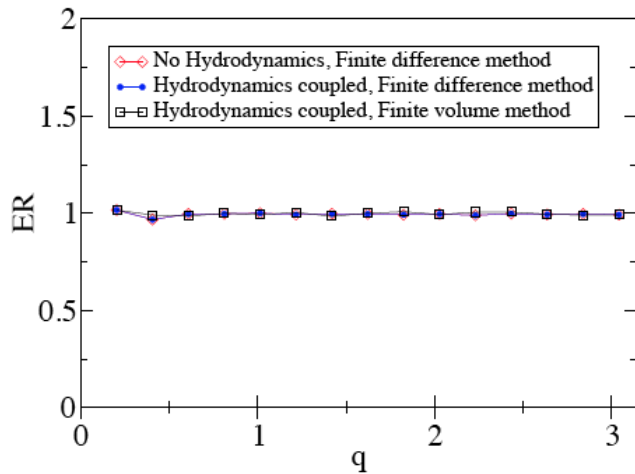


Standard approach to
Discretized Model B dynamics

9. Fluctuating hybrid schemes

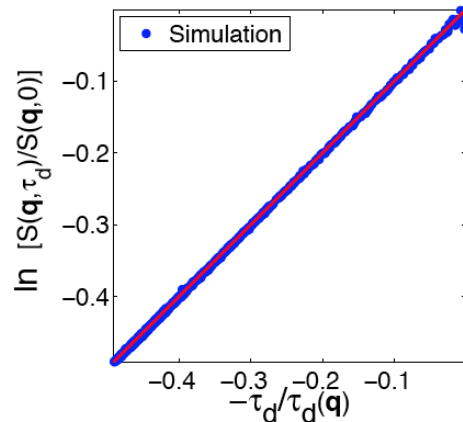
Proper fluctuation spectrum at all wave numbers

$$\langle \psi^2 \rangle = \frac{kT}{A}$$



No lattice anisotropies

Diffusive decay of correlations



$$\langle |\tilde{\psi}_q|^2 \rangle = \frac{kT}{A + Kq^2}$$

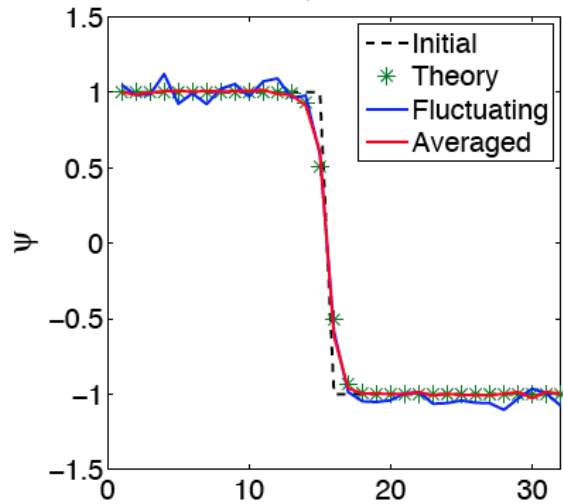
$$\langle \tilde{\psi}_q(t) \tilde{\psi}_q(t + \tau_d) \rangle = \frac{kT}{A + Kq^2} e^{-Mq^2(A + Kq^2)\tau_d}$$

9. Fluctuating hybrid schemes

Fluid/fluid coexistence

Capillary waves

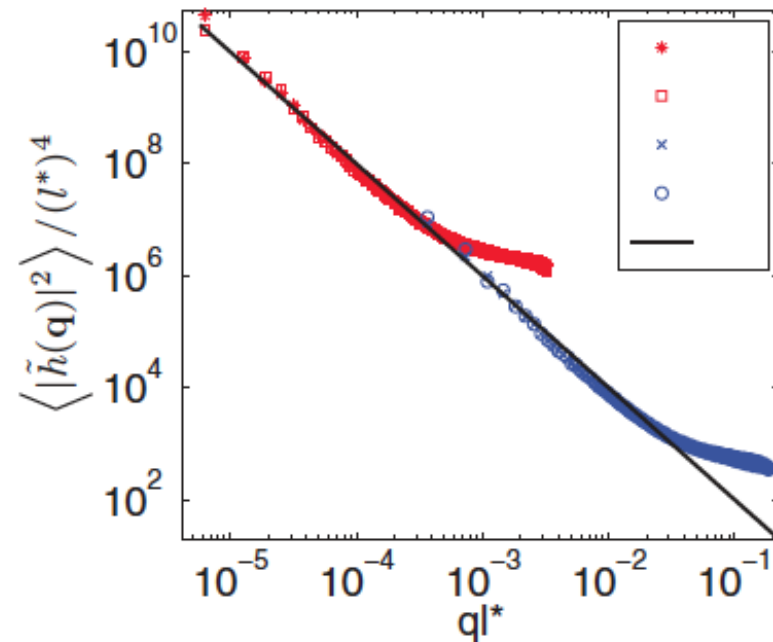
Average profile



Interface fluctuation spectrum

$$\Delta F_s = \frac{1}{2} \gamma \int dx dy (h_x^2 + h_y^2)$$

$$\Delta F_s = \frac{\gamma}{2} \sum_{\mathbf{q}} \mathbf{q}^2 |h(\mathbf{q})|^2 \quad \langle |h(\mathbf{q})|^2 \rangle = \frac{kT}{\gamma \mathbf{q}^2}$$



10. Conclusions

Hybrid lattice scheme

Lattice-Boltzmann for momentum conservation

Evolution of composition by link-flux

Consistent thermal fluctuations in general complex fluid mixtures

Moment propagation in general complex fluid mixtures

An efficient tool for charged, heterogeneous systems

Porous media (solid/liquid)

Immiscible solvents (liquid/liquid)

Charged oil/water mixture

Electrokinetic mobility of surfactant free emulsion

Emulsion electrokinetics

With solid surfaces

Capillary instabilities: role of dynamic wetting...